e+A(p) physics with an electron-ion collider at RHIC

Matthew A. C. Lamont Brookhaven National Lab

on behalf of the BNL EIC Science Task Force and friends



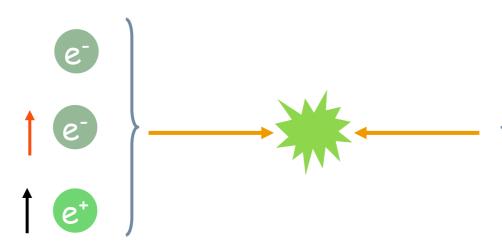


What is eRHIC?

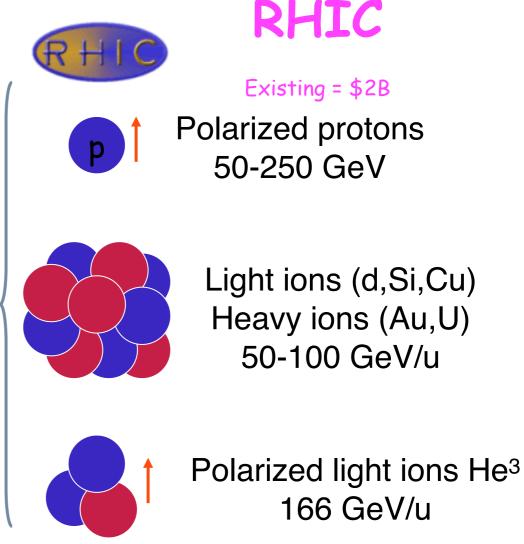
Electron accelerator

to be built

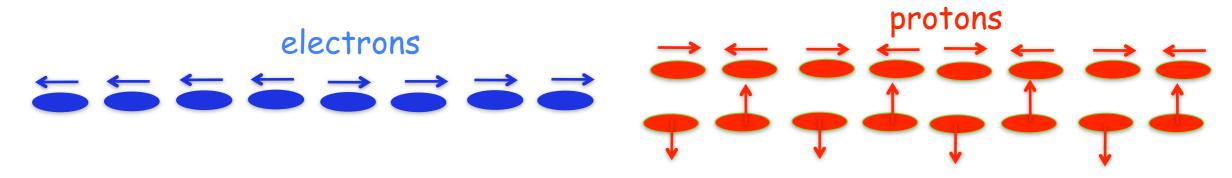
Unpolarized and polarized leptons 5-20 (30) GeV



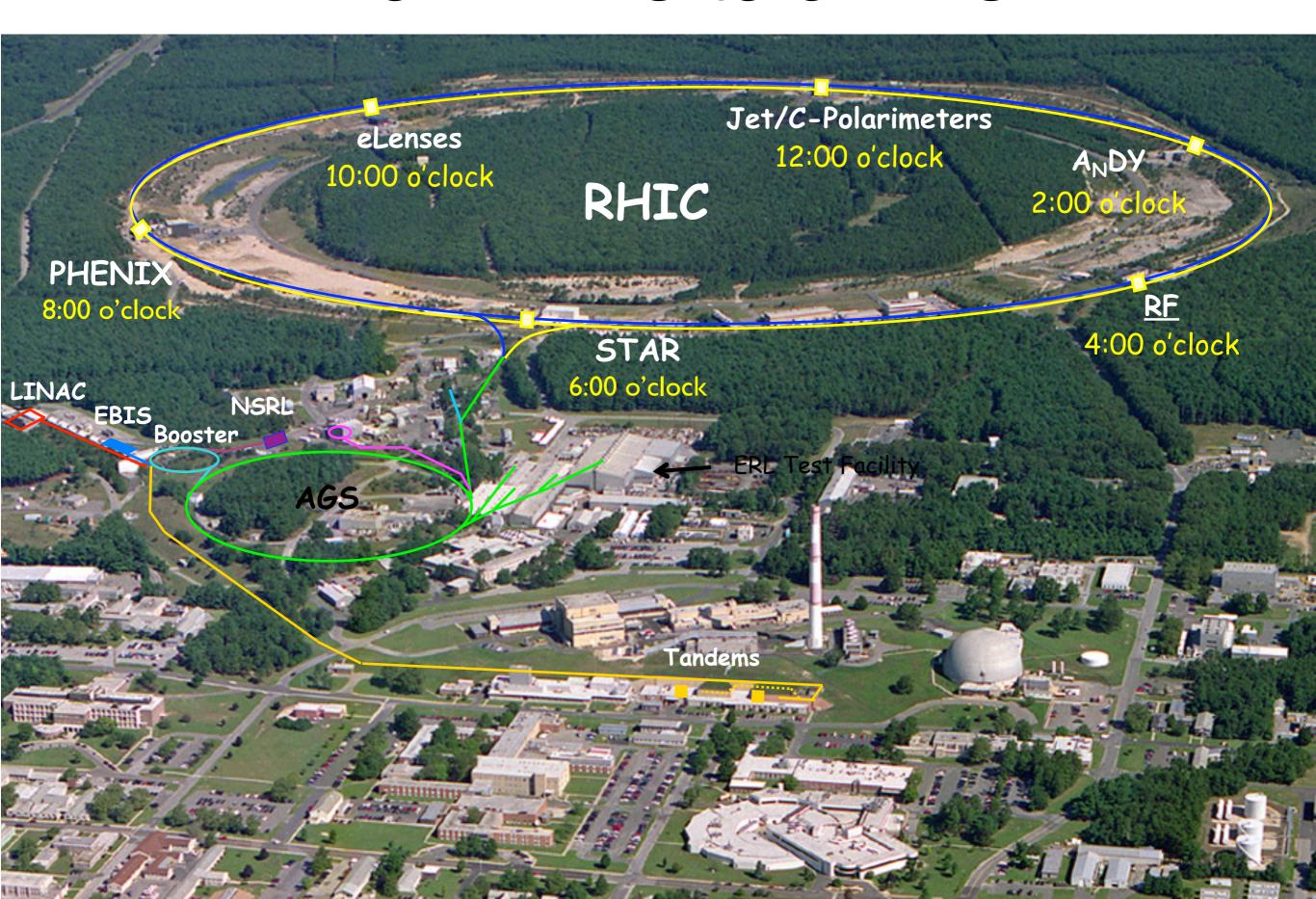
70% e- beam polarization goal polarized positrons?



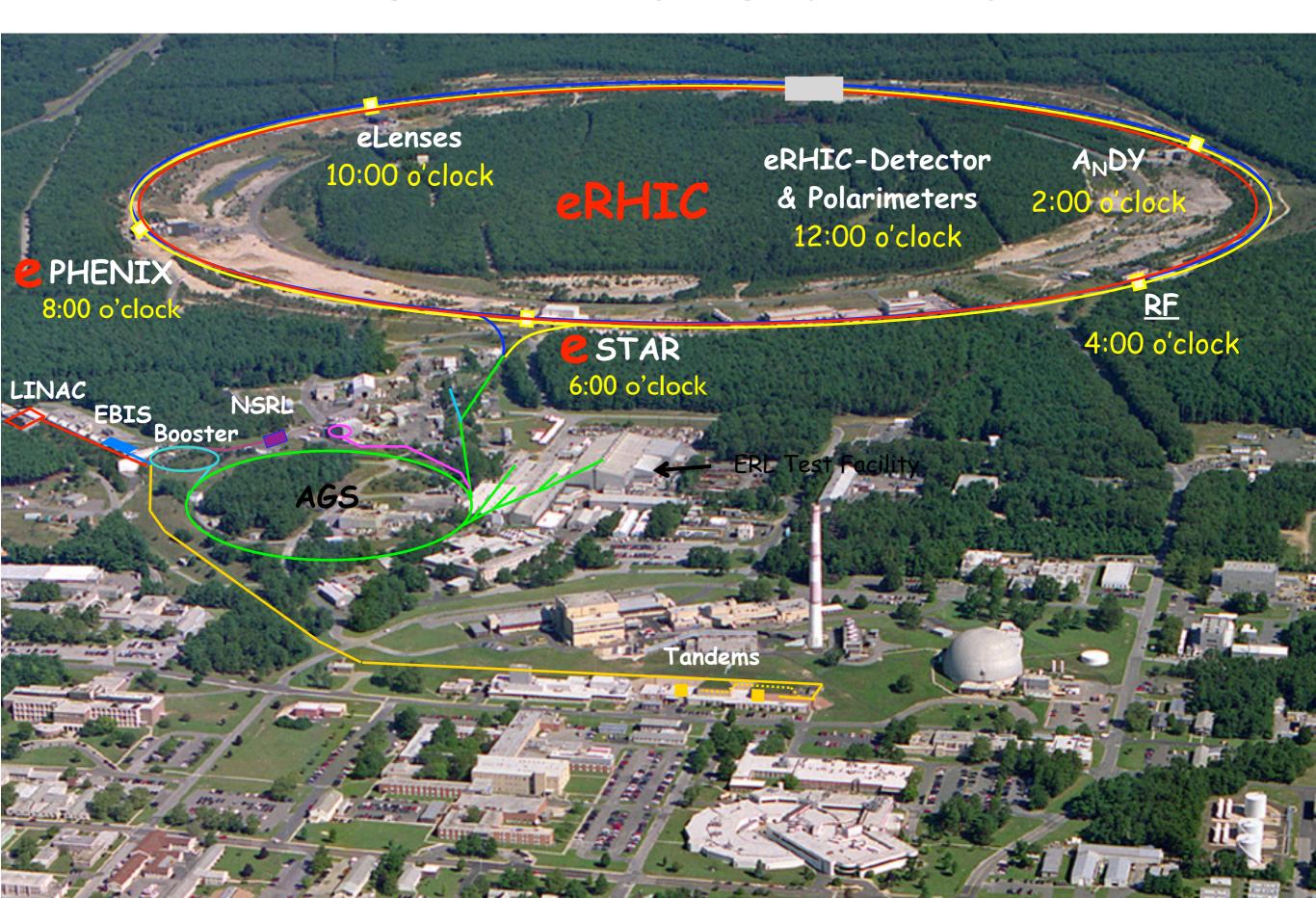
Centre-of-mass energy range: √s=30-200 GeV; L~100-1000xHera longitudinal and transverse polarization for p/He³ possible



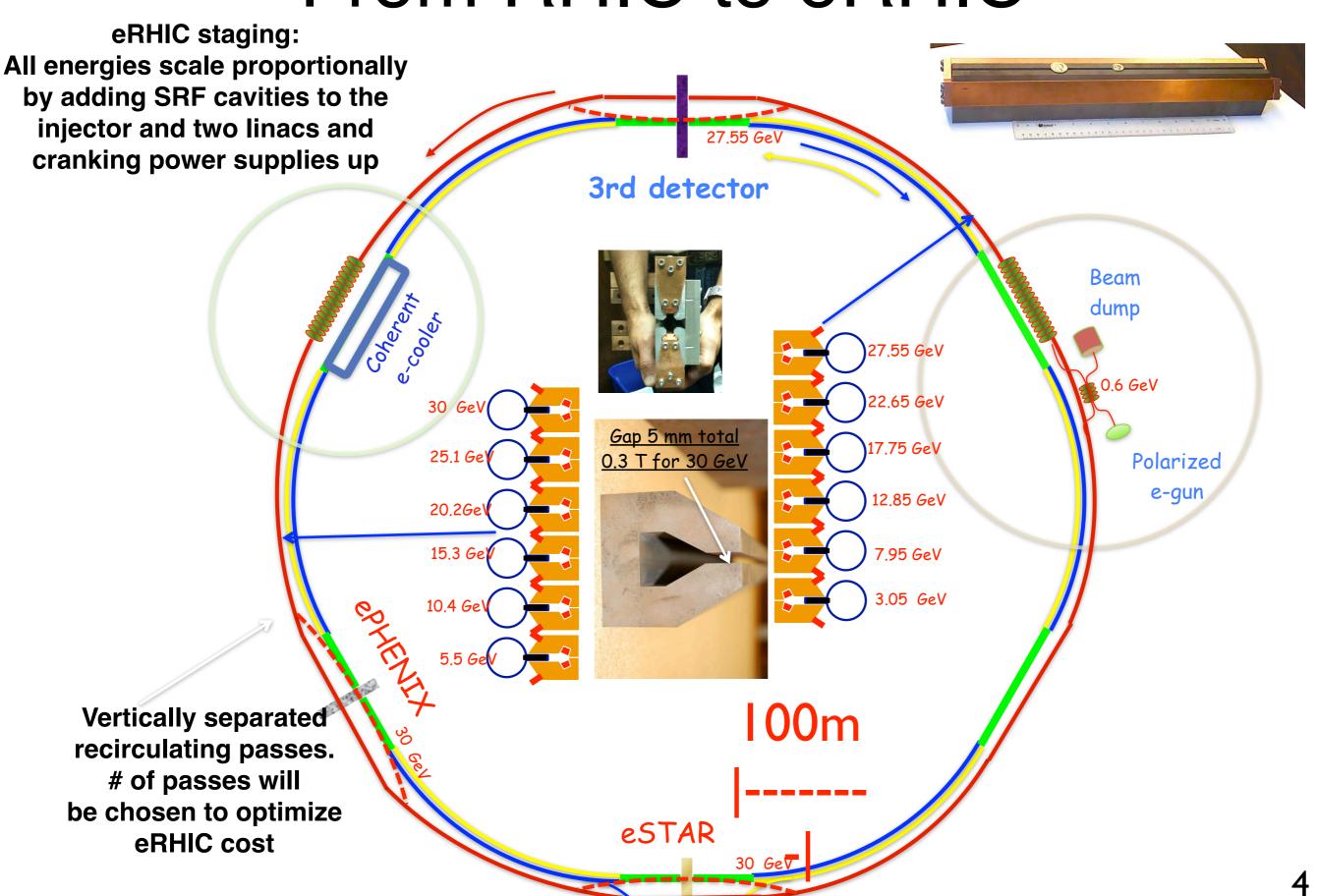
From RHIC to eRHIC



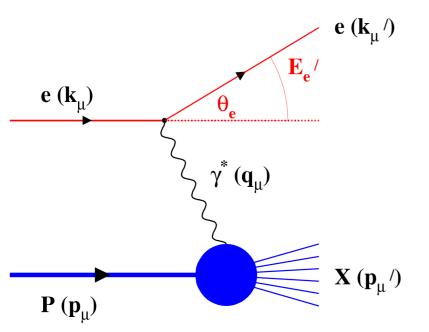
From RHIC to eRHIC



From RHIC to eRHIC



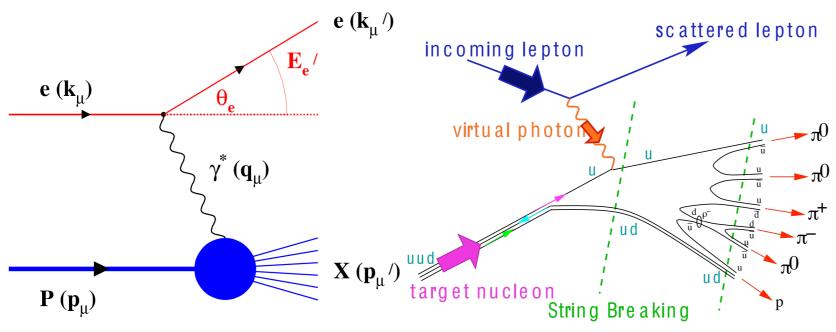
Detector requirements



Inclusive Reactions:

- Momentum/energy and angular resolution of e' critical
- Very good electron pid
- Moderate luminosity $> 10^{32}$ cm⁻¹ s⁻¹
- Need low x ~ 10^{-4} high \sqrt{s} (Saturation and spin physics)

Detector requirements



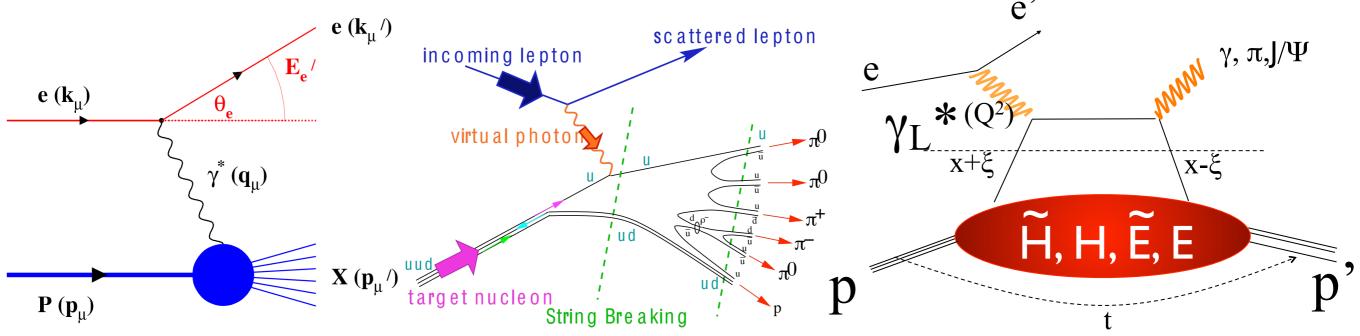
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Semi-inclusive Reactions:

- \bullet Excellent particle ID: π ,K,p separation over a wide range in η
- $_{ullet}$ full Φ -coverage around γ^*
- high luminosity > 10^{33} cm⁻¹ s⁻¹ (5d binning (x,Q²,z,p_t, Φ))
- Need low x ~ 10^{-4} high \sqrt{s}

Detector requirements



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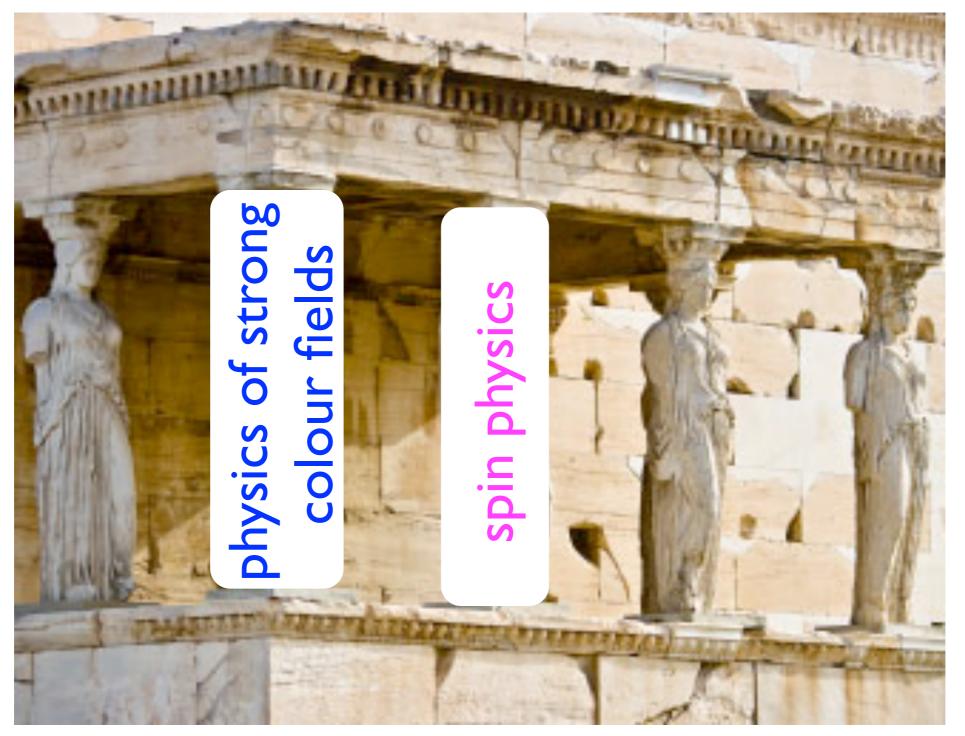
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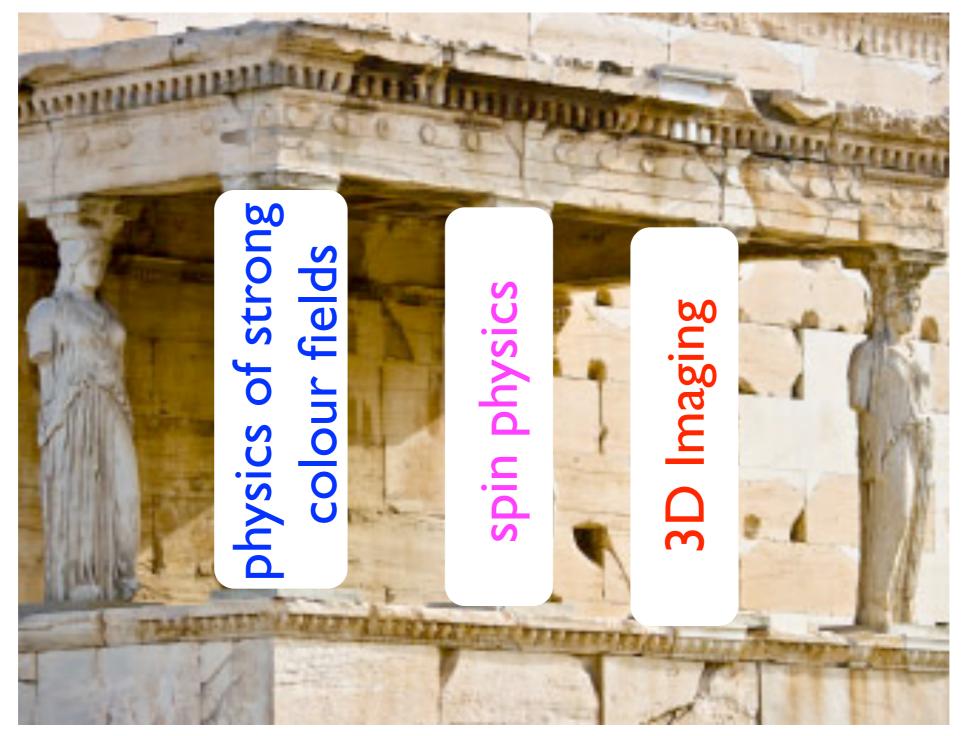
Exclusive Reactions:

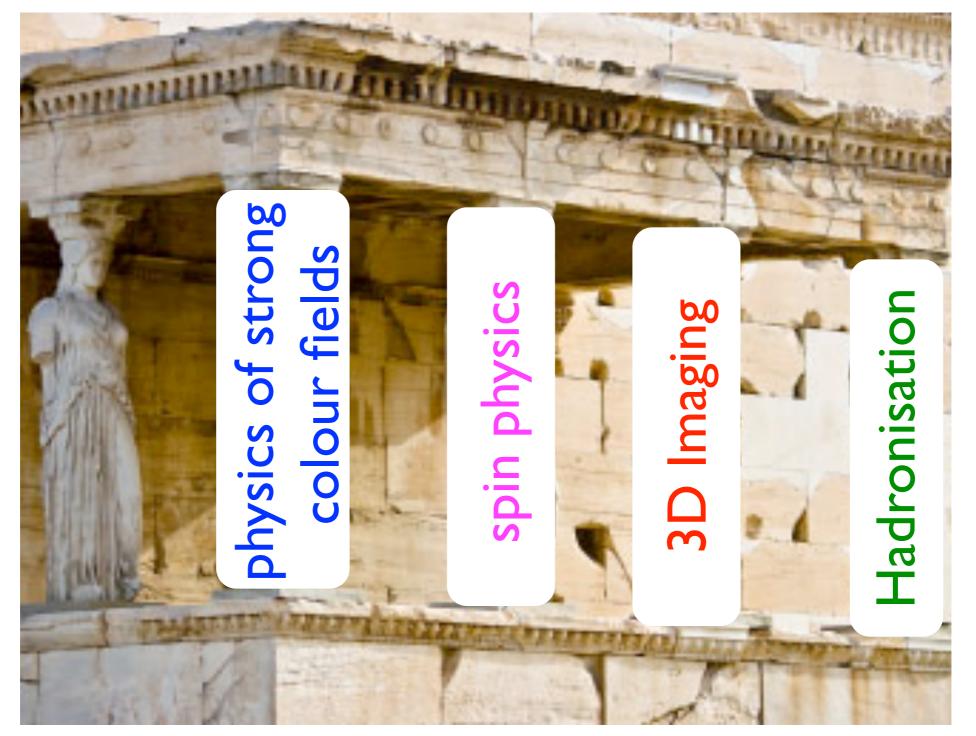
- Exclusivity → high rapidity coverage → rapidity gap events
- high resolution in t → Roman pots
- high luminosity > 10^{33} cm⁻¹ s⁻¹ (4d binning (x,Q²,t, Φ))

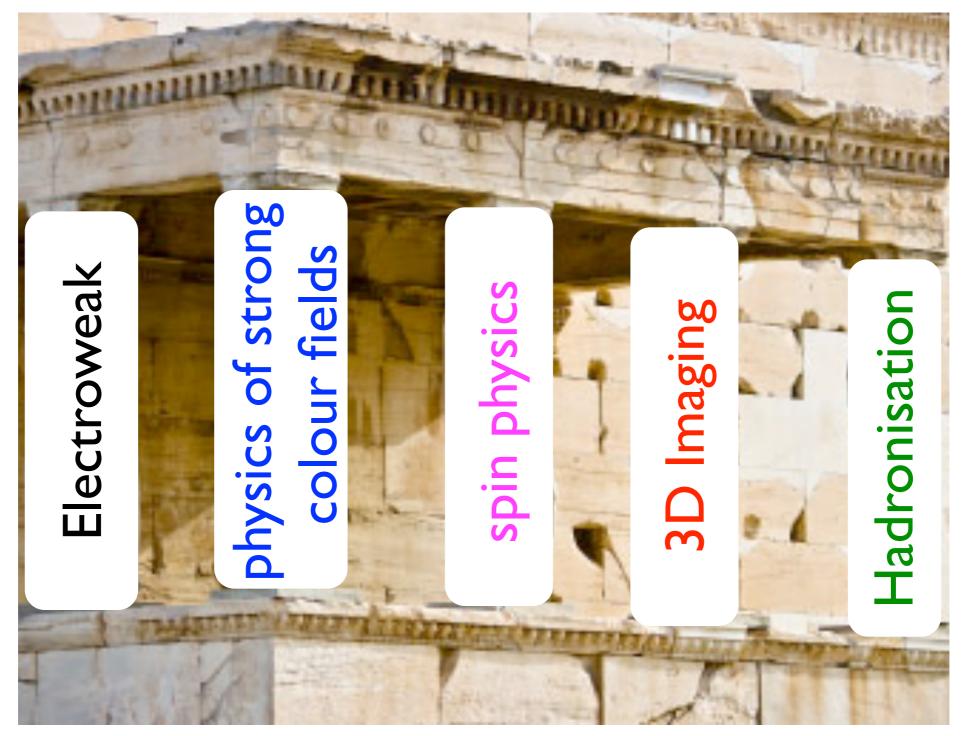


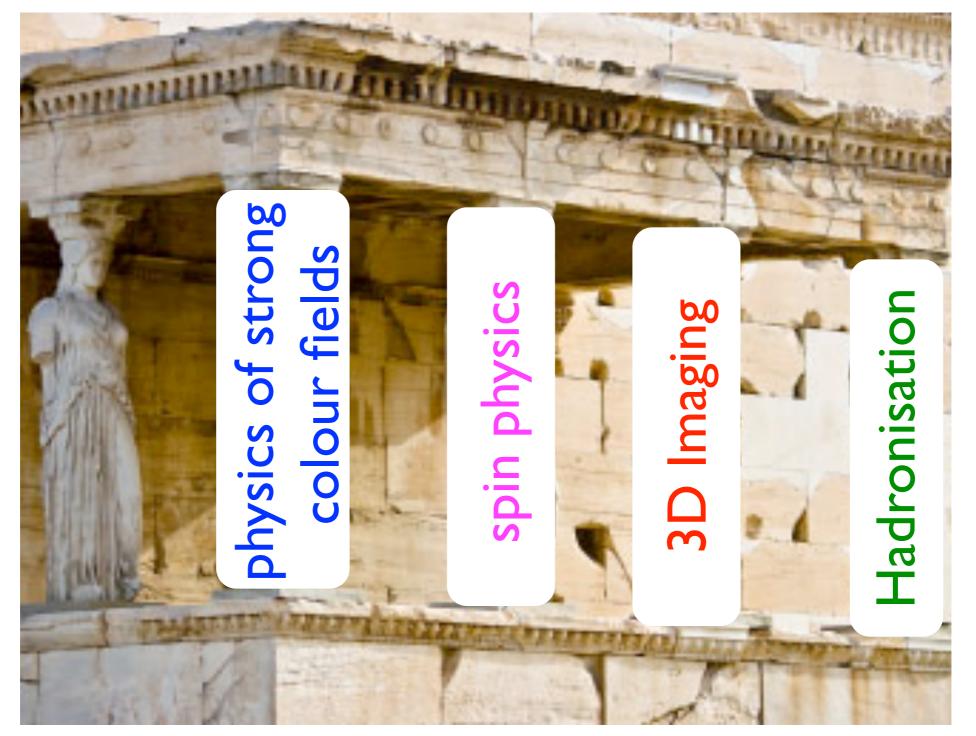


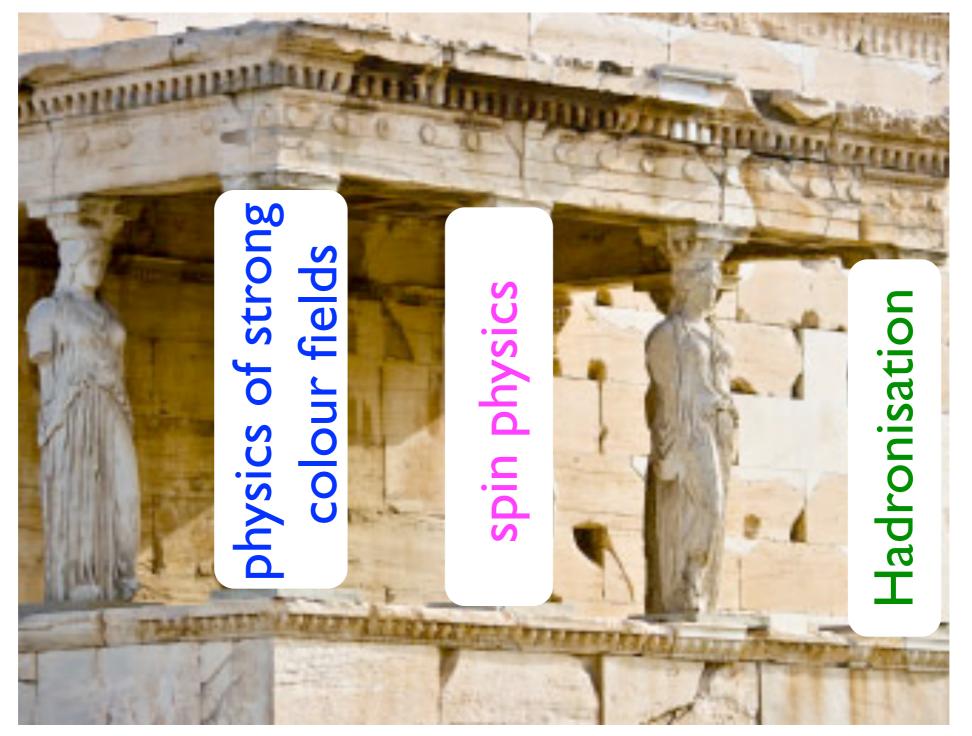










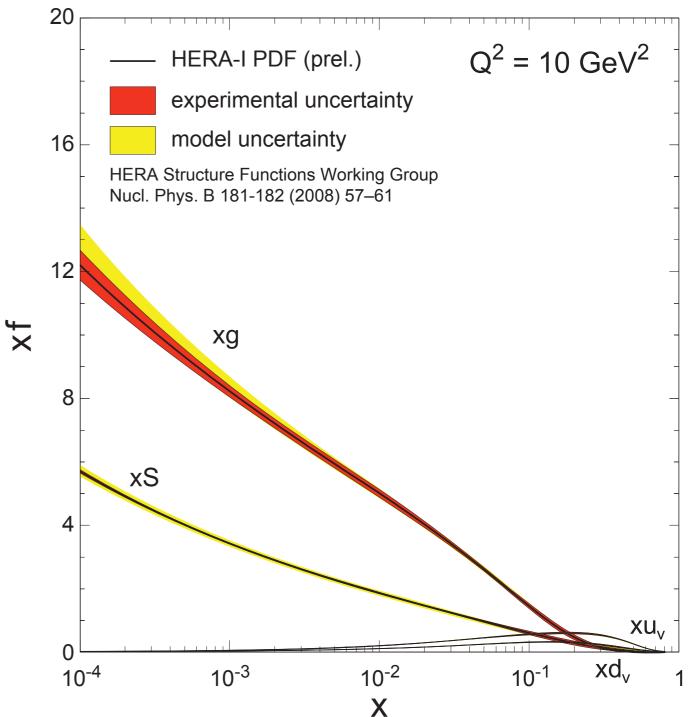


Most compelling physics questions

Spin physics

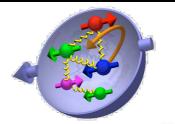
- What is the polarisation of gluons at small x where they dominate?
- What is the x-dependence and flavour decomposition of the polarised sea?

Determine quark and gluon contributions to the proton spin at last!!



Most compelling physics questions

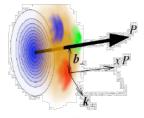
Spin physics



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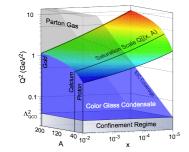
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Imaging

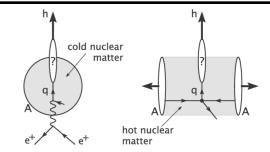


- What is the spatial distribution of quarks/ gluons in nucleons AND nuclei?
- Understand deep aspects of gauge theories revealed by k_T dependent distributions

Possible window to orbital angular momentum



Strong Colour Fields and Hadronisation



- Quantitatively probe the universality of strong colour fields in A+A, p+A and e+A
- Understand in detail the transition to the non-linear regime of strong gluon fields and the physics of saturation
- How do hard probes in e+A interact with the medium?

Currently have no experimental knowledge of gluons in nuclei at small x!!

spin physics

10 week INT programme - Fall 2010

Organizers:

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Talks online

Application form

Exit report

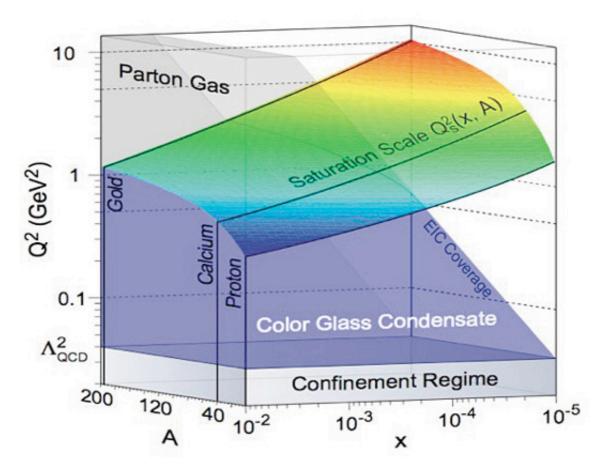
Friends of the INT

Obtain an INT preprint number

INT homepage

Gluons and the quark sea at high energies: distributions, polarization, tomography

September 13 to November 19, 2010



This INT program will address open questions about the dynamics of gluons and sea quarks in the nucleon and in nuclei. Answers to these questions are crucial for a deeper understanding of hadron and nuclear structure in QCD at high energies. Many of them are relevant for understanding QCD final states at the LHC, which often provide a background for physics beyond the standard model. The topics addressed in this program have important ramifications for understanding the matter produced in heavy-ion collisions at RHIC and the LHC.

http://www.int.washington.edu/PROGRAMS/10-3/

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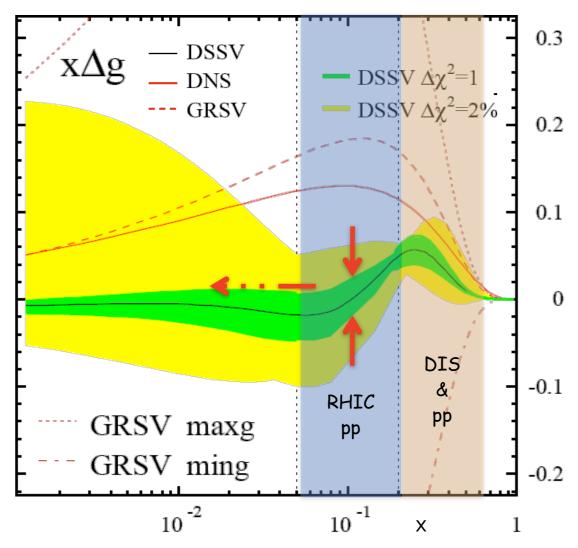
Golden measurements in spin

Deliverables	Observables	What we learn	Requirements
polarised gluon distribution Δ <i>g</i>	scaling violations in inclusive DIS	gluon contribution to the proton spin	coverage down to x ~ 10 ⁻⁴ ; £ of about 10 fb ⁻¹
polarised quark and antiquark densities	semi-incl. DIS for pions and kaons	quark contr. to proton spin; asym. like Δū-Δđ; Δs	similar to DIS; good particle ID
novel electroweak spin structure functions	inclusive DIS at high Q ²	flavour separation at medium-x and large Q ²	\sqrt{s} ≥ 100 GeV; L ≥ 10 fb ⁻¹ ; positrons; polarised ³ He beam

The quest for Δg - where do we stand?

- inclusive pions and jets remain the main probes
- jet/hadron correlations are essential to cover smaller x

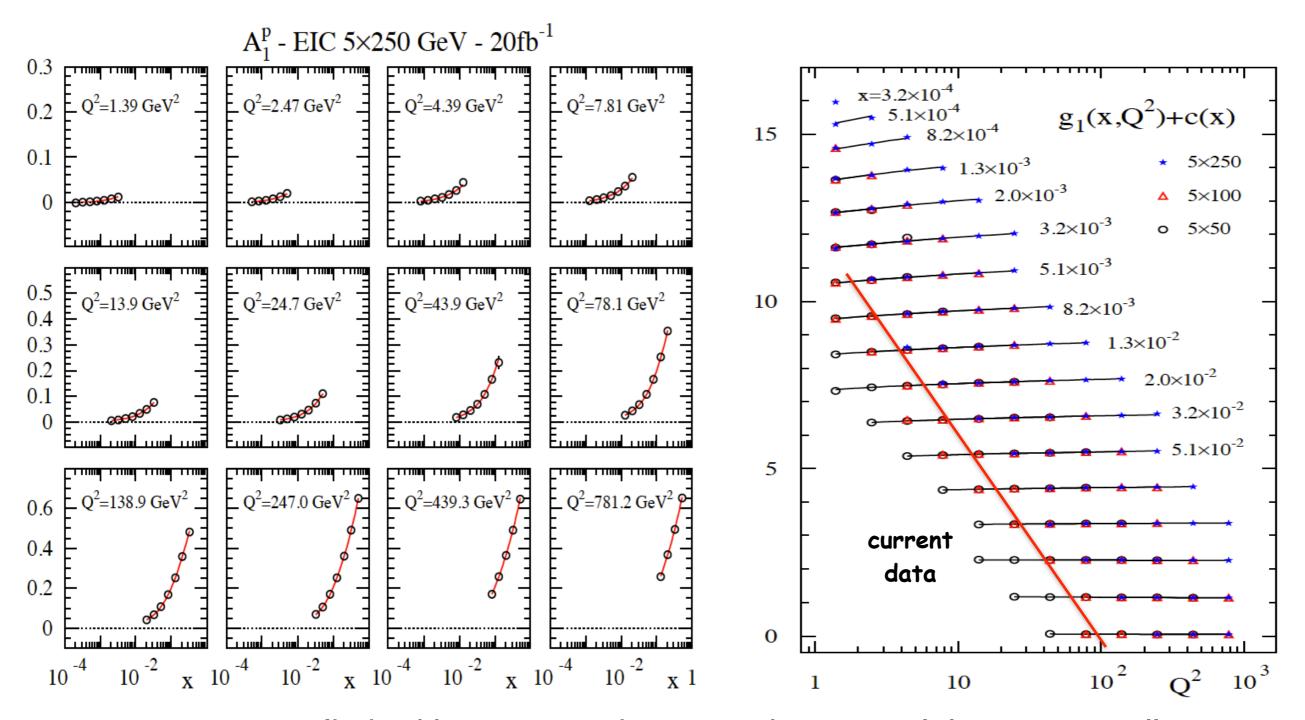




- low-x behaviour is unconstrained
 - significant polarisation still possible
 - → no reliable error estimate for 1st moment
- By 2015 expect to have:
 - DSSV 2.0 global analysis on new world data
 - → reduced uncertainties in Δg in current x range
 - evidence of a node further scrutinised
 - extend x-range towards lower x
 - 500 GeV running and particle correlations

DSSV includes "only" RHIC run6 data

strategy to quantify impact: global QCD fits with realistic pseudo-data



measurements limited by systematics – need to control them very well

issues: bunch-by-bunch polarimetry, relative luminosity, detector performance, ...

how effective are scaling violations ? quantitative studies based on simulated data for eRHIC stage-1: $5 \times (50, 100, 250, 325)$ GeV

$$\chi^{2} \text{ profile for } \int_{10^{-4}}^{1} \Delta g(x,Q^{2}) \, dx$$

$$\Delta \chi^{2}_{i}$$
 from current ep & RHIC data
$$\Delta g^{1,[0.0001-1]}$$

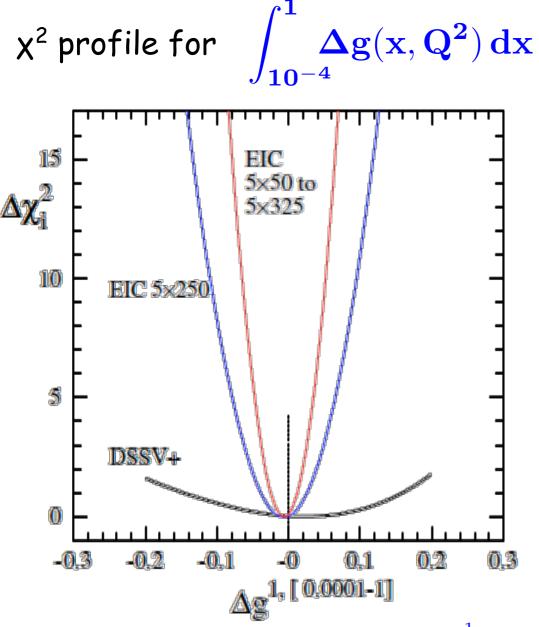
expect to determine $\int_0^1 dx \, \Delta g(x,Q^2)$ at about 10% level (more studies needed) kinematic reach down to $x = 10^{-4}$ essential to determine integral

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$$\chi^2$$
 profile for $\int_{10^{-4}}^{1} \Delta g(x, Q^2) dx$
 $\Delta \chi^2_i$
 $\int_{-0.3}^{15} -0.2 -0.1 -0 0.1 0.2 0.3$
 Δg^{1} , [0.00001-1]

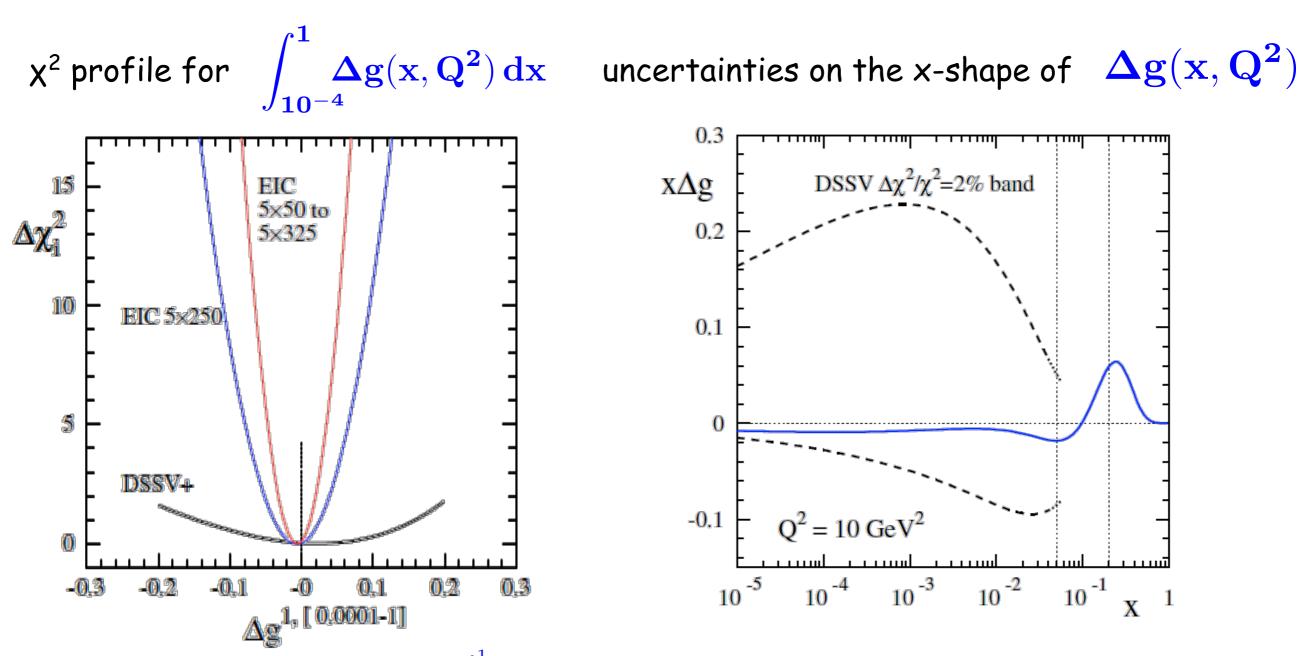
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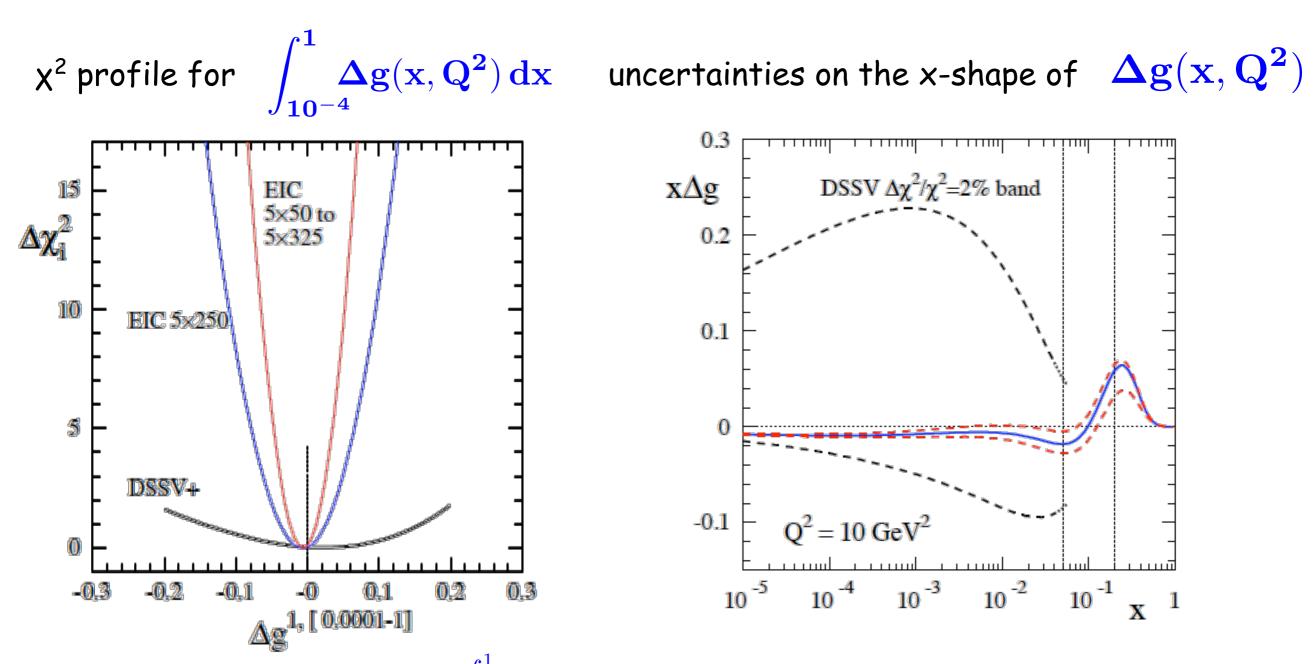
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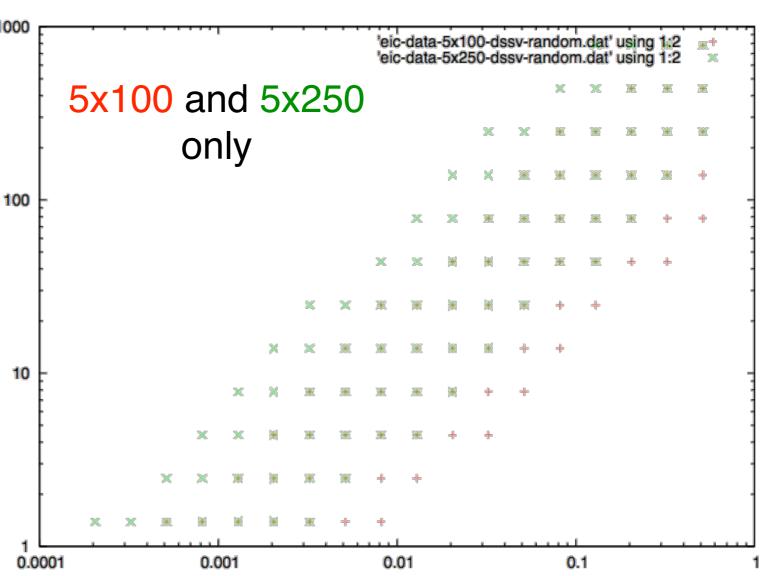
New pseudo-data

data for DIS and SIDIS (π±, K±)

10 fb⁻¹ each, 70% beam pol.

Cuts:

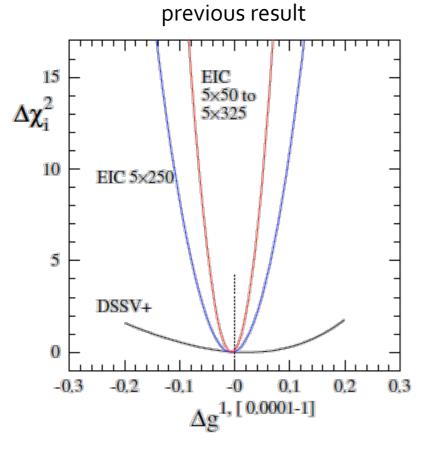
- $W^2 > 10 \text{ GeV}^2$
- depol. factor > 0.1
- $\odot 0.001 < y < 0.95$
- \bullet 1° < θ < 179°
- p_e > 0.5 GeV
- phadr > 1 GeV

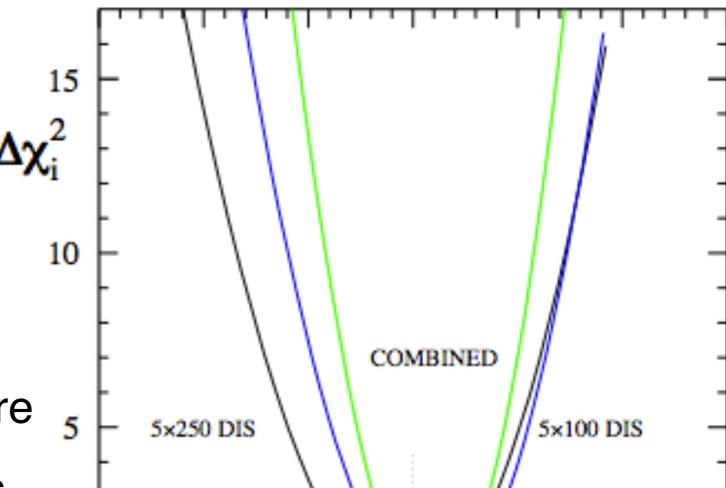


• Global analysis:

- use relative uncertainty of each point to produce mock data (based on DSSV)
- randomise data within 1σ
- → for SIDIS: incl. 5%(10%) uncertainty from pion (kaon) frag. functions
- \rightarrow map out X^2 profiles with Lagrange multiplier method (Hessian is work in progress)

Update on Δg





DSSV+

-0.2

-0.1

 $\Delta g^{1,[0.001-1]}$

-0.3

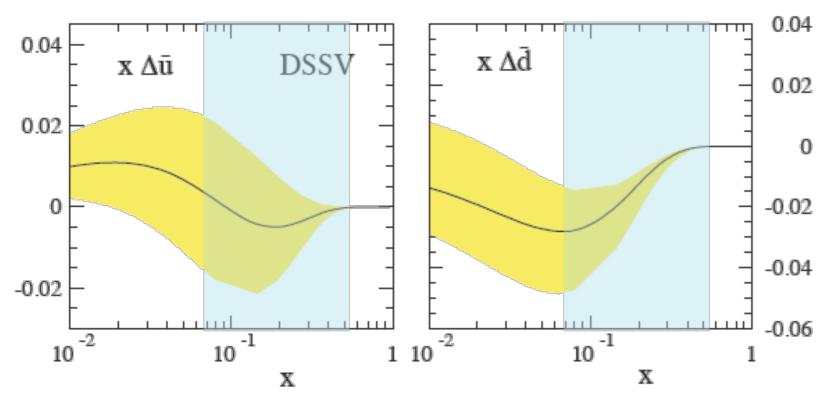
- Very similar results to before
- Slightly larger uncertainties
- need to study 10⁻⁴→1 range
- need to translate into error on x-shape of Δg

0.3

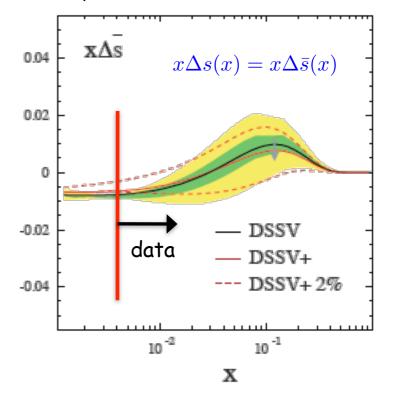
0.2

What about Δq ?



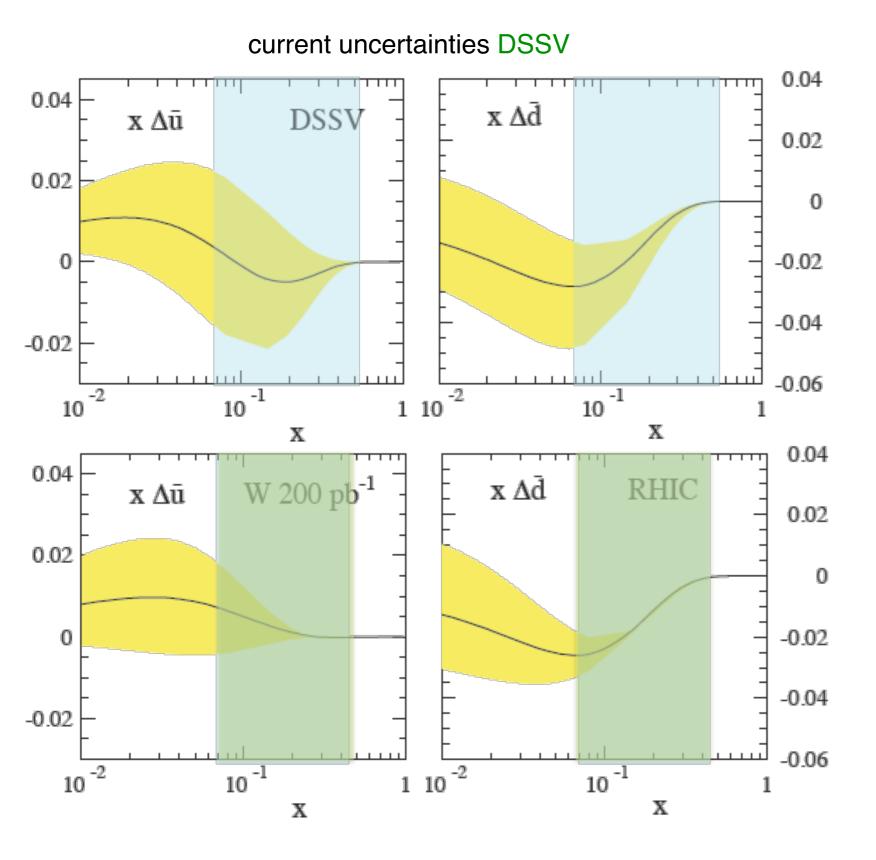


DSSV (incl. latest COMPASS data)

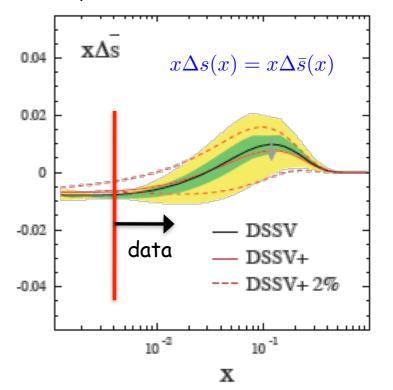


- surprise: ∆s small & positive from SIDIS data
- but 1st moment is negative and sizable due to "constraint" from hyperon decays (F,D) (assumed SU(3) symmetry - debatable M. Savage)
- \odot drives uncertainties on $\Delta\Sigma$ (spin sum)

What about Δq ?



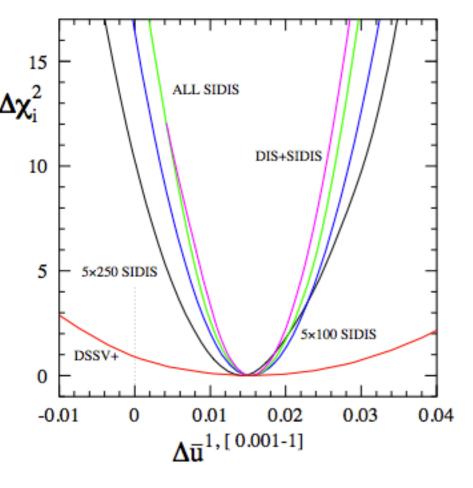
DSSV (incl. latest COMPASS data)



simulated impact of RHIC

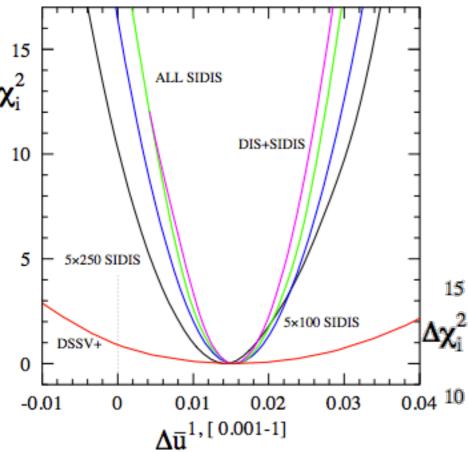
- W boson data on global fit
- reduction of uncertainties for 0.07 < x < 0.4 can test consistency of low Q² SIDIS data in that x regime

First results on the quark sea

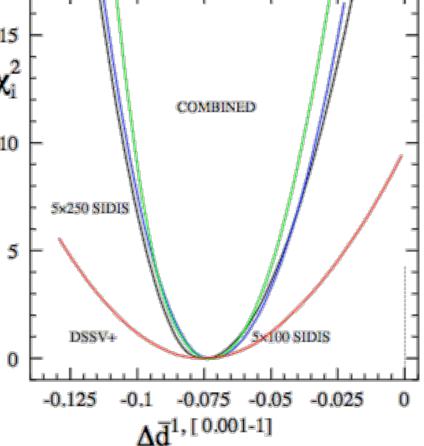


- very encouraging results
- as expected, DIS has no impact
- need to study 0.0001-1 range
- need to translate into error on x-shape

First results on the quark sea

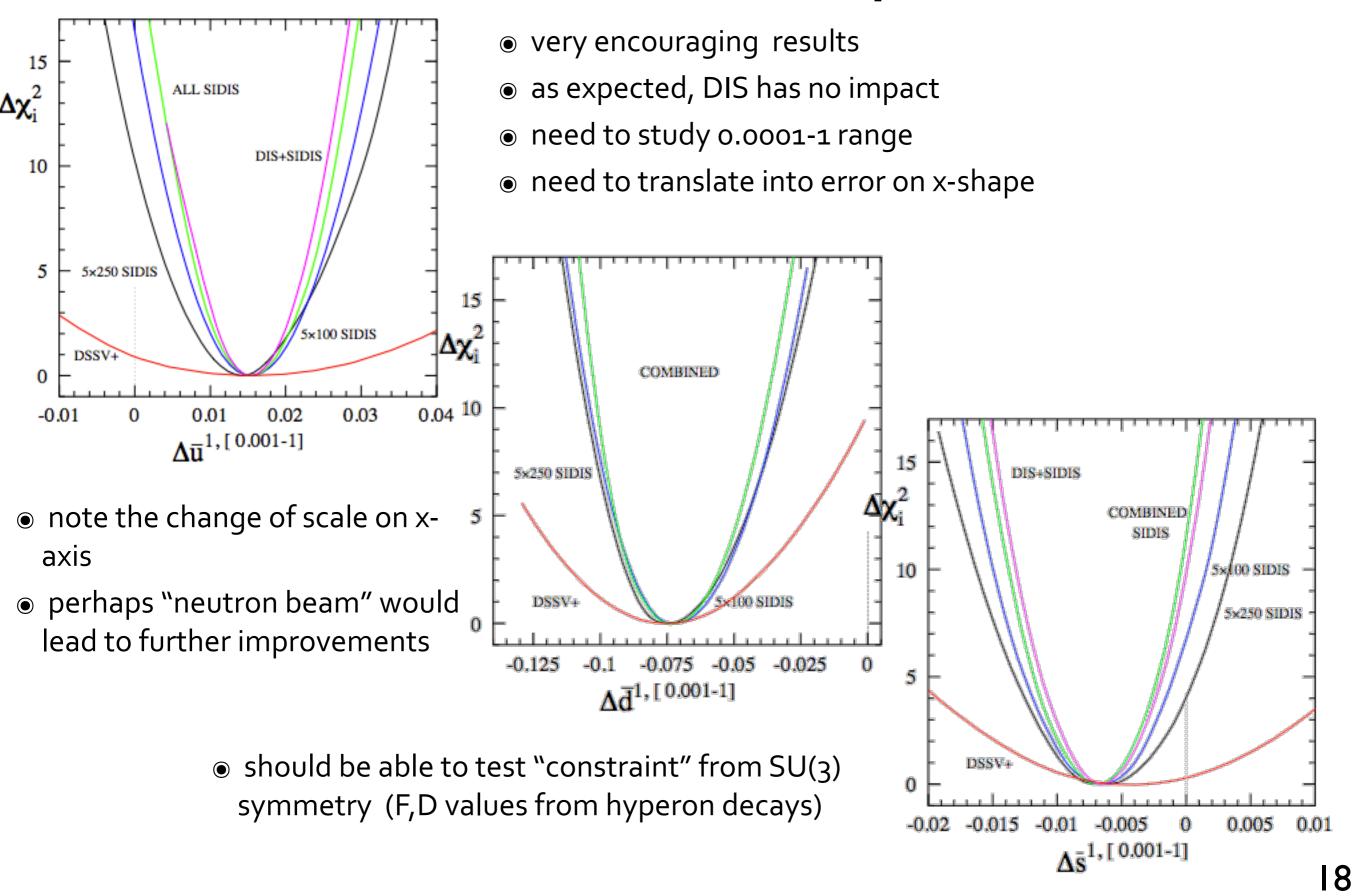


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- note the change of scale on xaxis
- perhaps "neutron beam" would lead to further improvements

First results on the quark sea



e+A physics

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Golden Measurements

Deliverables	Observables	What we learn	Stage-1	Stage-II
integrated gluon distributions	F _{2,L}	nuclear wave function; saturation, Qs	gluons at 10 ⁻³ < x < 1	saturation regime
k _T dependent gluons; gluon correlations	di-hadron correlations	non-linear QCD evolution / universality	onset of saturation	measure Q _s
transport coefficients in cold matter	large-x SIDIS; jets	parton energy loss, shower evolution; energy loss mechanisms	light flavours and charm; jets	rare probes and bottom; large-x gluons

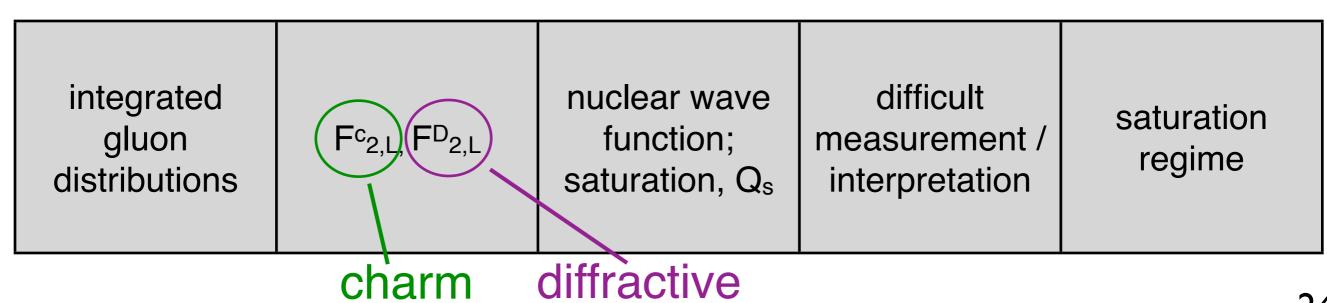
Silver Measurements

Deliverables	Observables	What we learn	Stage-I	Stage-II
integrated gluon distributions	F ^c _{2,L,} F ^D _{2,L}	nuclear wave function; saturation, Q _s	difficult measurement / interpretation	saturation regime
flavour separated nuclear PDFs	charged current and γZ structure functions	EMC effect origin	full flavour separation for $10^{-2} < x < 1$	measure Q _s
k⊤ dependent gluons	SIDIS at small x	non-linear QCD evolution / universality	onset of saturation	rare probes and bottom; large-x gluons
b-dependent gluons; gluon correlations	DVCS; diffractive vector mesons	interplay between small- x evolution and confinement	moderate x with light, heavy nuclei	smaller x, saturation

Integrated gluon distributions from inclusive structure functions

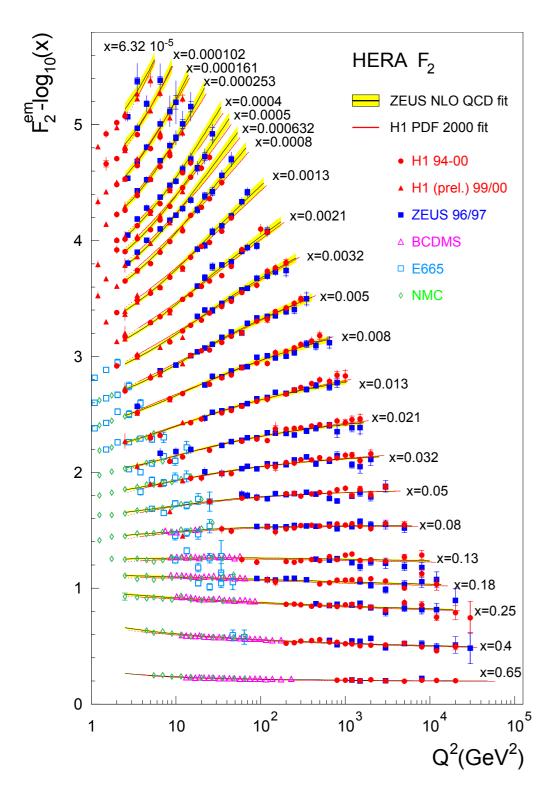
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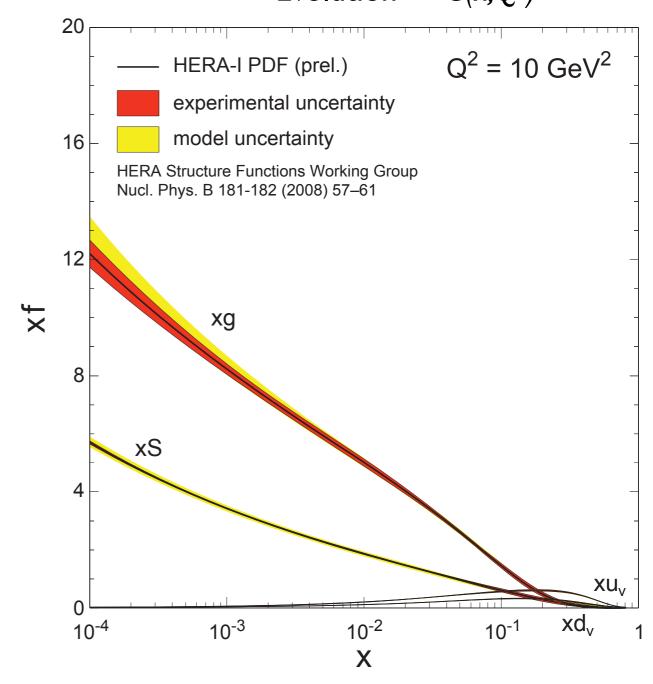


Measuring the glue via Structure Functions

$$\sigma_r(x, Q^2) = F_2^A(x, Q^2) - \frac{y^2}{Y^+} F_L^A(x, Q^2)$$

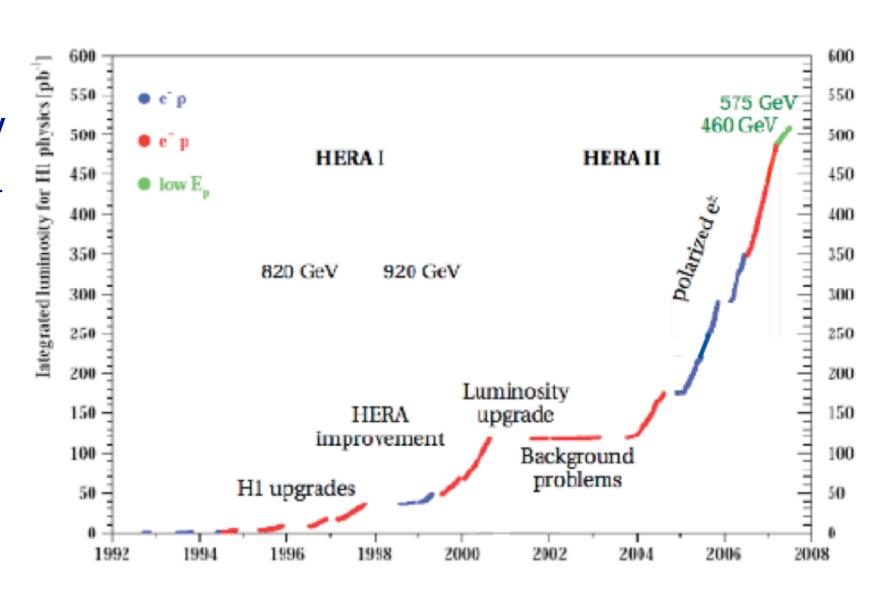


Scaling violation: $dF_2/dlnQ^2$ and linear DGLAP Evolution $\Rightarrow G(x,Q^2)$



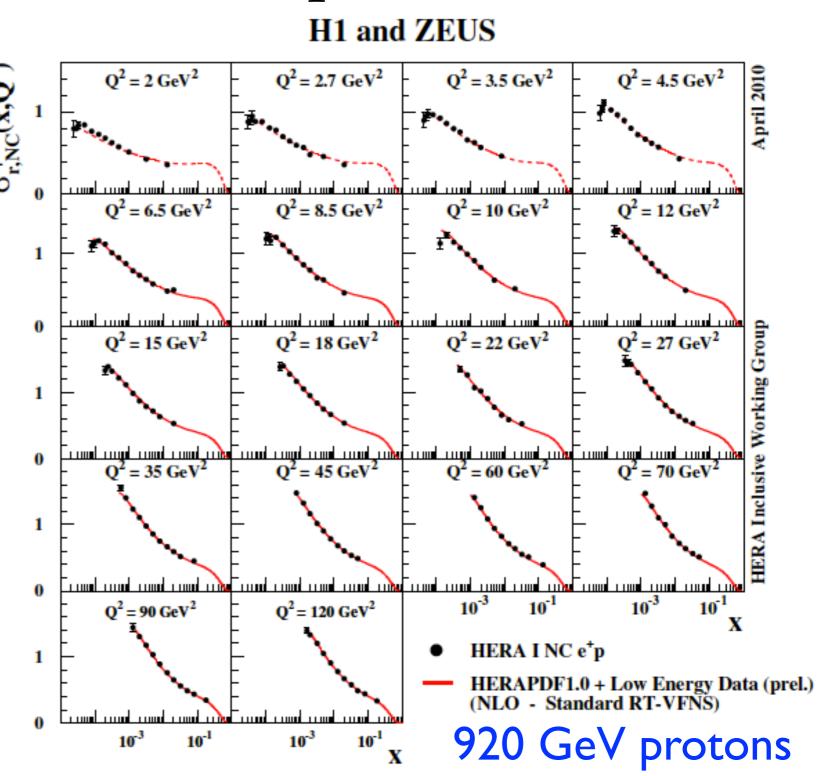
$$\sigma_r(x, Q^2) = F_2^A(x, Q^2) - \frac{y^2}{Y^+} F_L^A(x, Q^2)$$

- \bullet F_L ~ α_s xG(x,Q²)
 - \rightarrow y = Q²/xs
 - → require an energy scan to extract F_L
- 3 different proton energies run at HERA
 - 2 low-statistics runs
 - → bad for F_L extraction



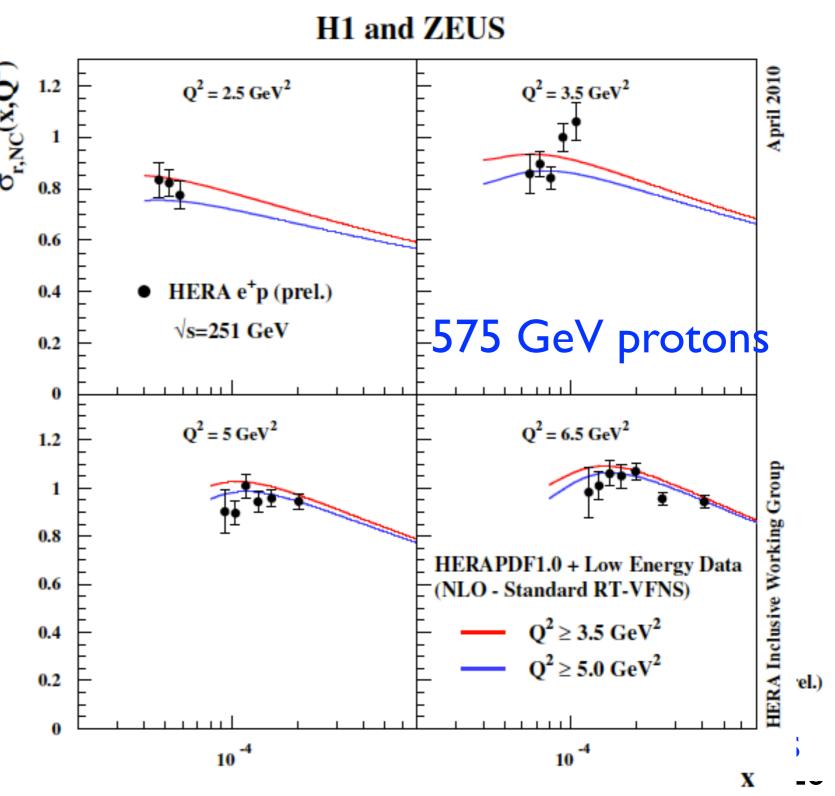
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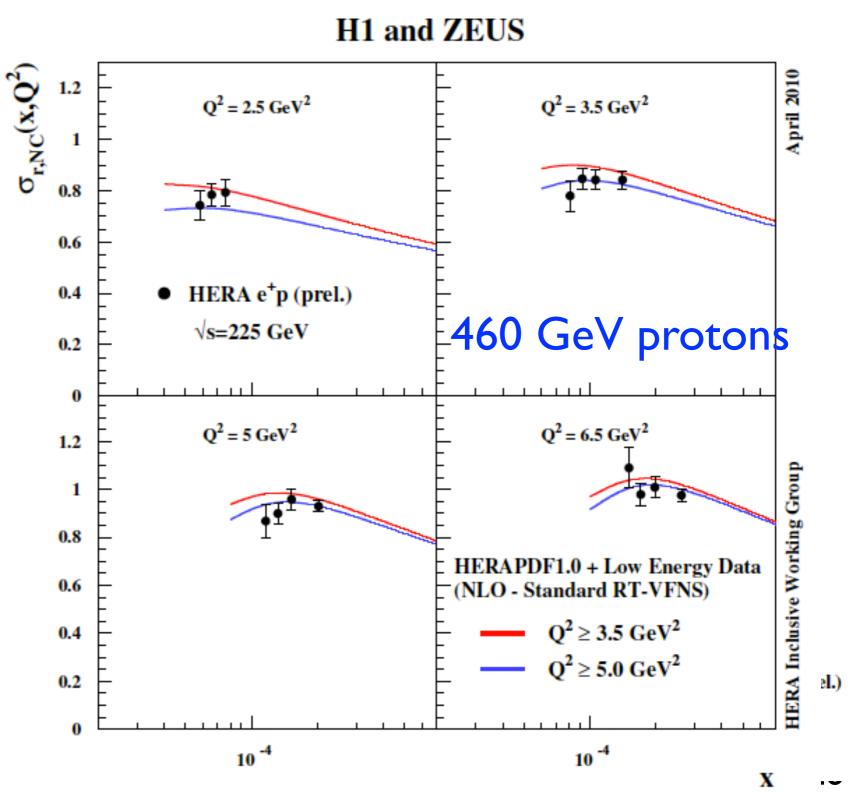
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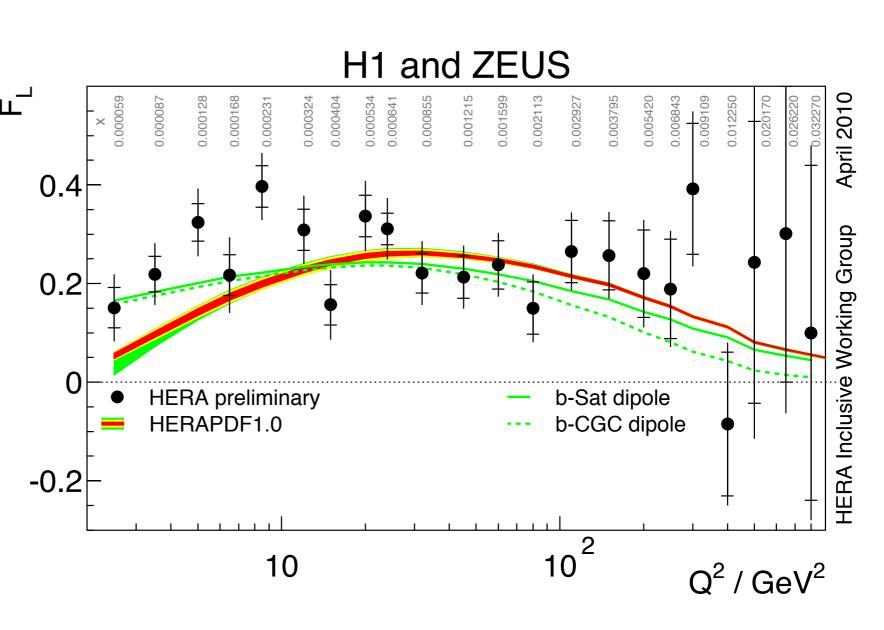
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- $F_L \sim \alpha_s xG(x,Q^2)$
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 - 2 low-statistics runs
 - → bad for F_L extraction



Are F₂/F_L good differentiators of models?

 In order to see if measuring F₂/F_L at an EIC is "worthwhile", we need to see if a measurement could differentiate between models

Models:

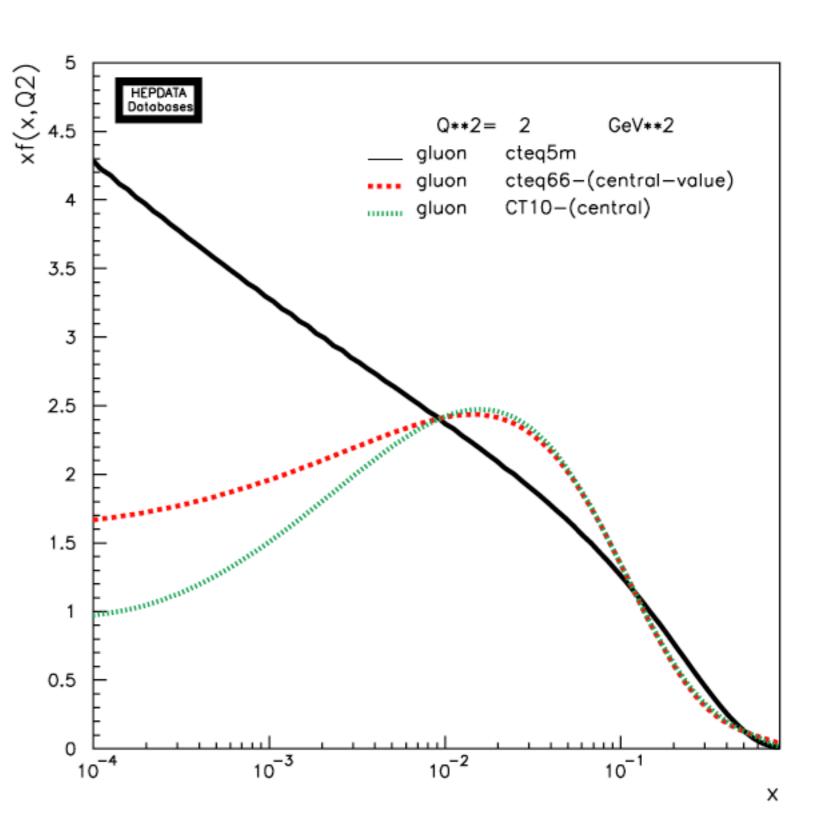
- MSTW08: code downloadable from HEPFORGE. Code to extract F₂/F_L obtained privately from Graeme Watt
 - Global fit, using total cross-section from HERA
 - DGLAP evolution
- → IPSat: data kindly provided by T. Lappi
 - Fit to ZEUS'96 data χ^2 /d.o.f. ~ 1.2
- bCGC: data kindly provided by T. Lappi
 - Fit to Zeus'96 data χ²/d.o.f. ~ 1.62
- rcBK: AAQMS data kindly provided by J. Albacete
 - Evolution along x with BK equation
 - ► Fit to H1+ZEUS combined 2006 data
- → Leading-Twist Shadowing: FGS10 data kindly provided by V. Guzey
 - Evolved with DGLAP

• Data:

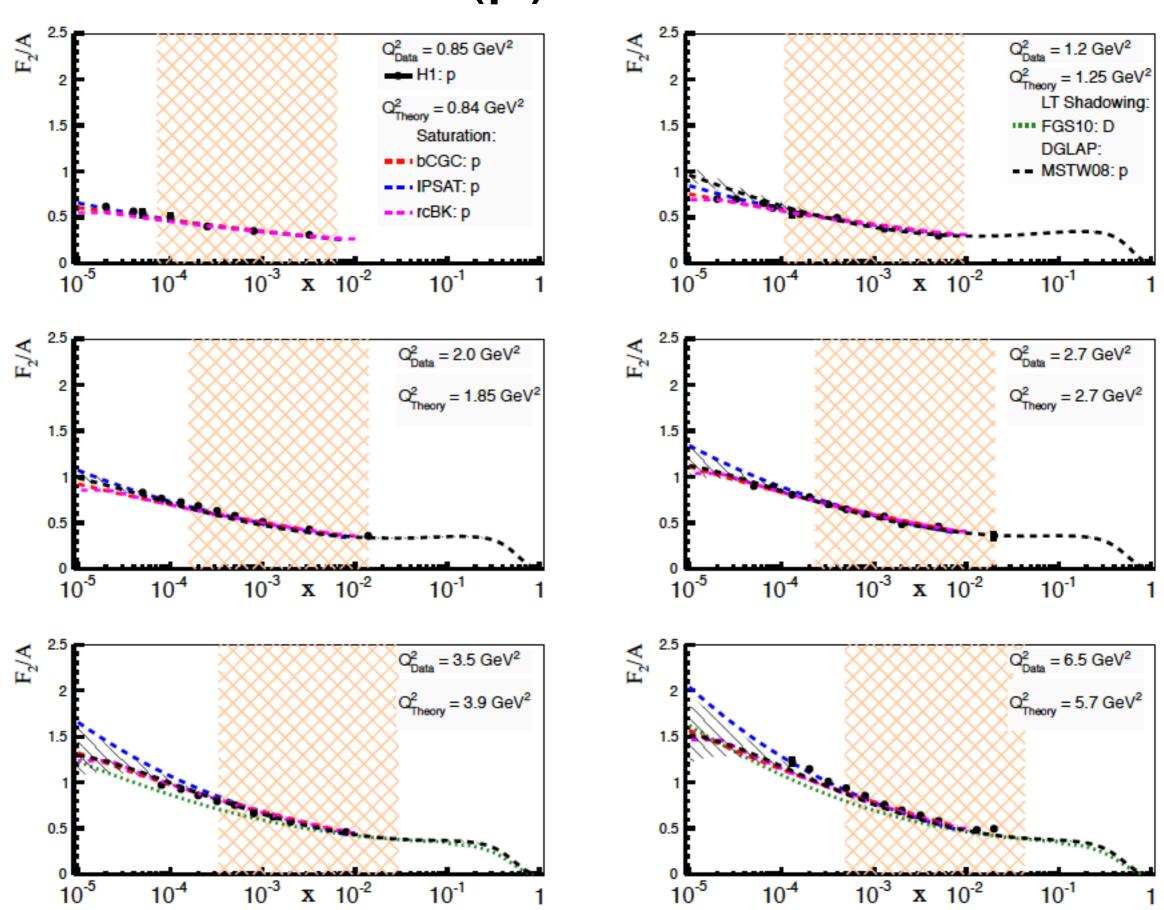
- → F₂: H1&Zeus combined data from: http://www-h1.desy.de/psfiles/papers/desy09-158.pdf
- → F_L: H1 data from: http://www-h1.desy.de/psfiles/papers/desy10-228.pdf

Using the correct PDF

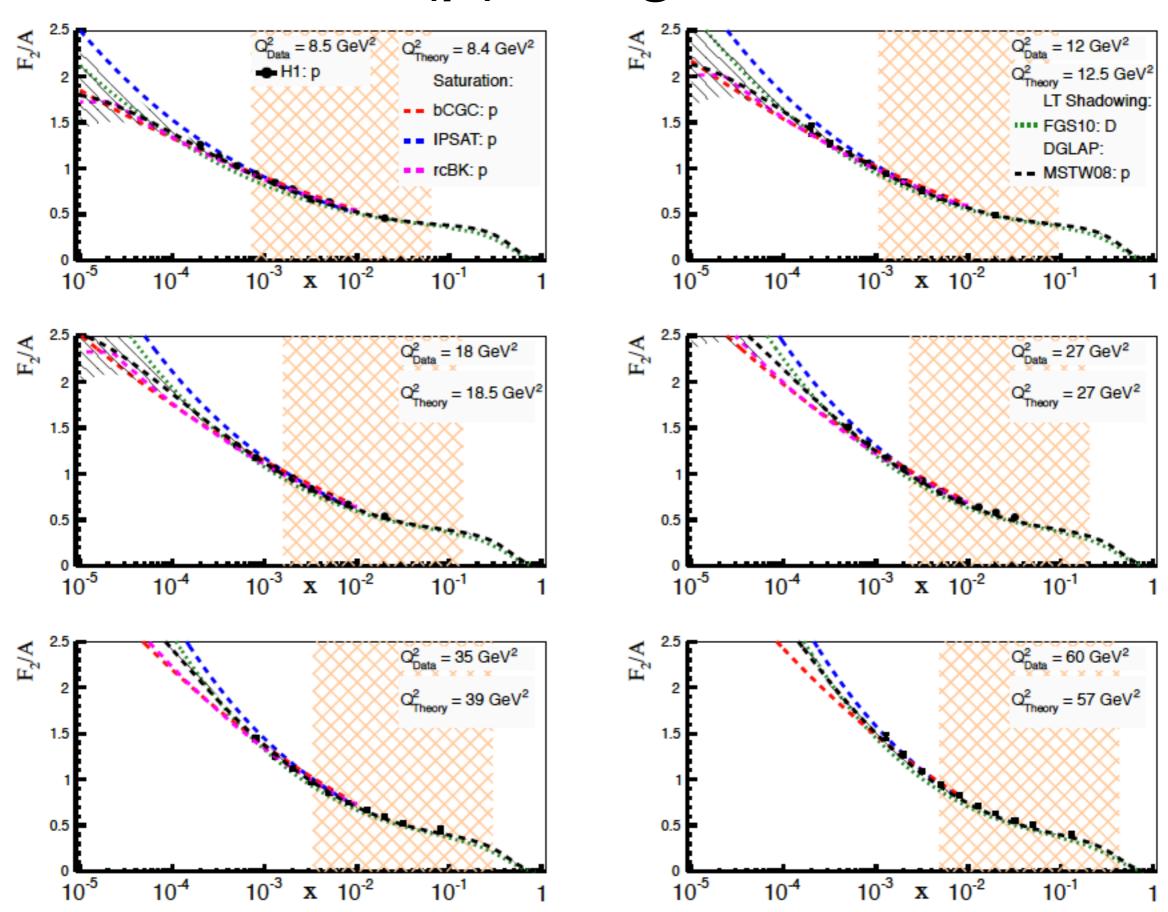
- This overestimates the gluon contribution quite drastically compared to more modern calculations
 - New curves are on their way from FGS10 with CTEQ6
 - Not ready for this meeting
 - → Following F_L data therefore still uses CTEQ5m for FGS10



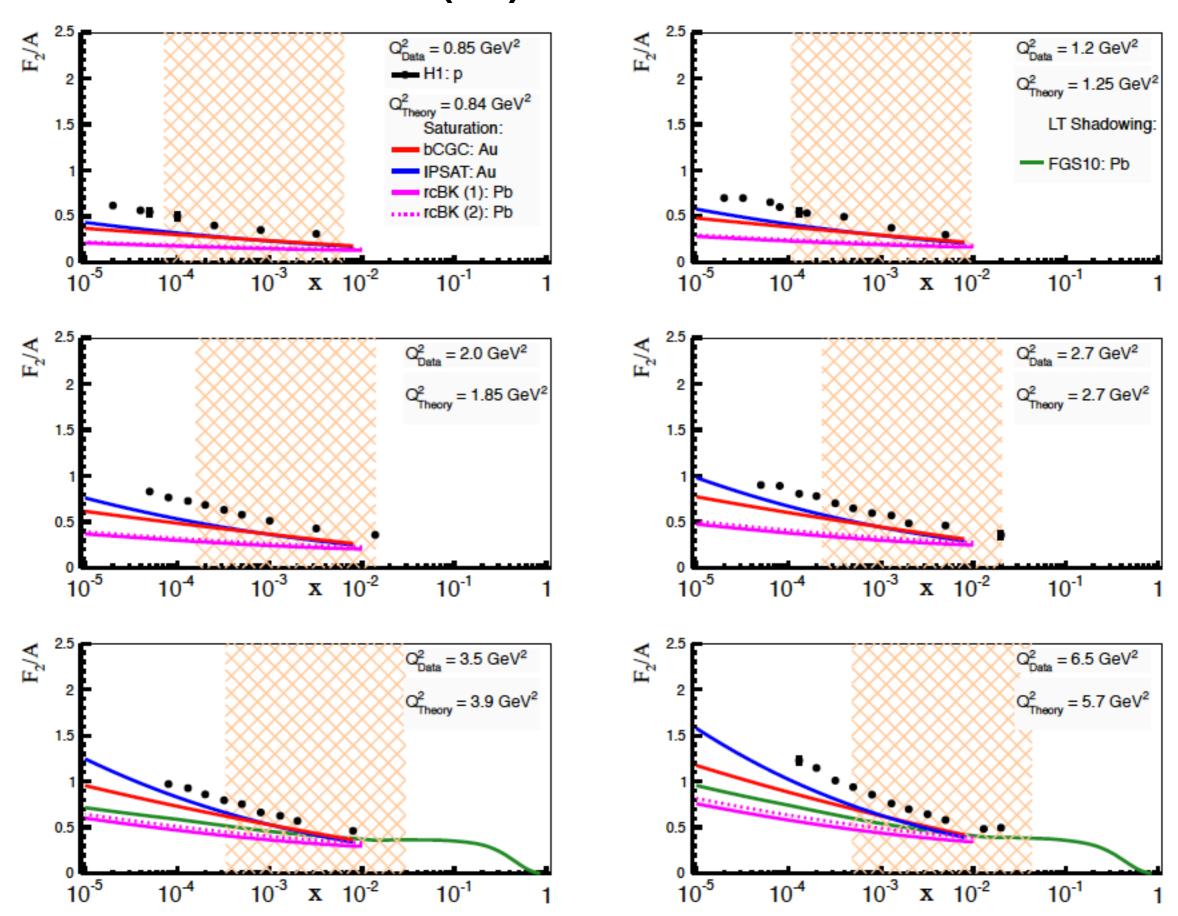
$F_2(p)$ - low Q^2



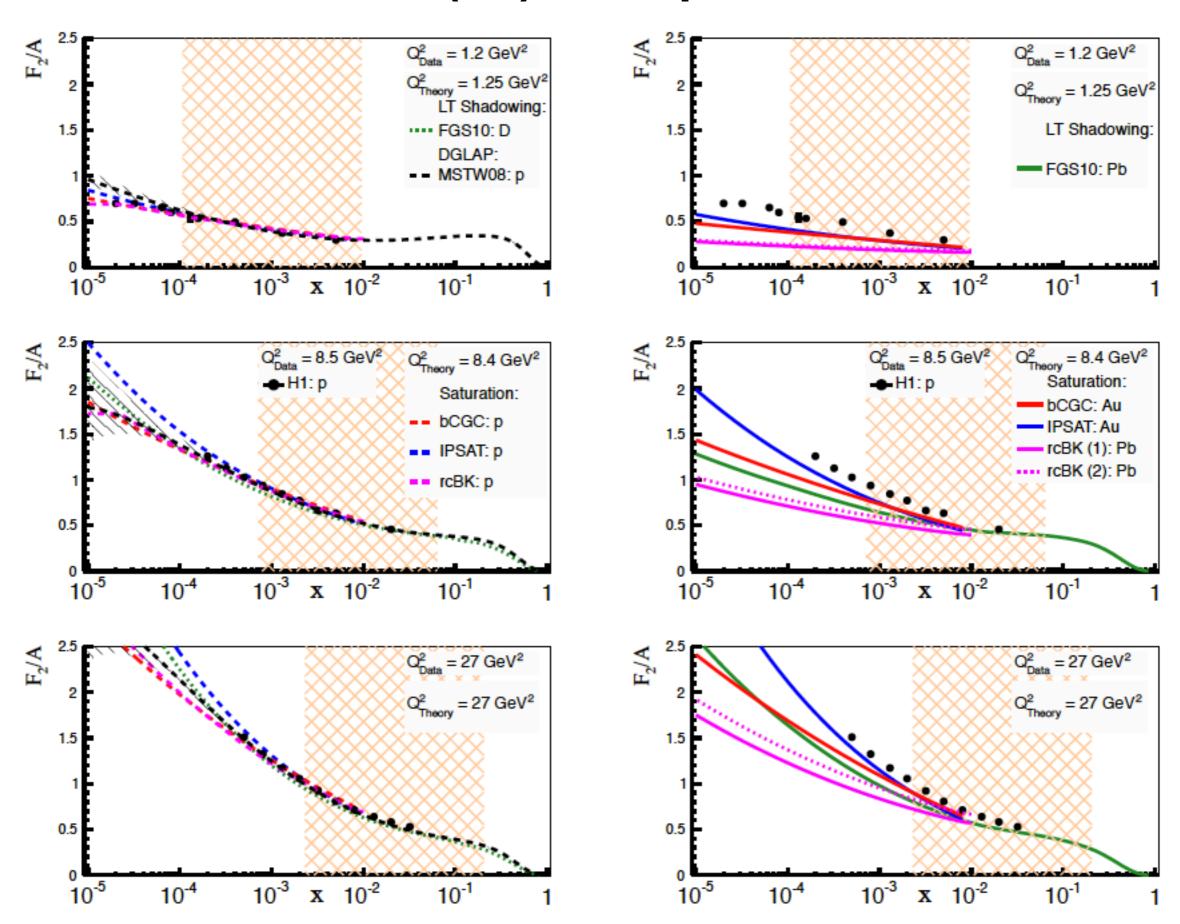
$F_2(p)$ - higher Q^2



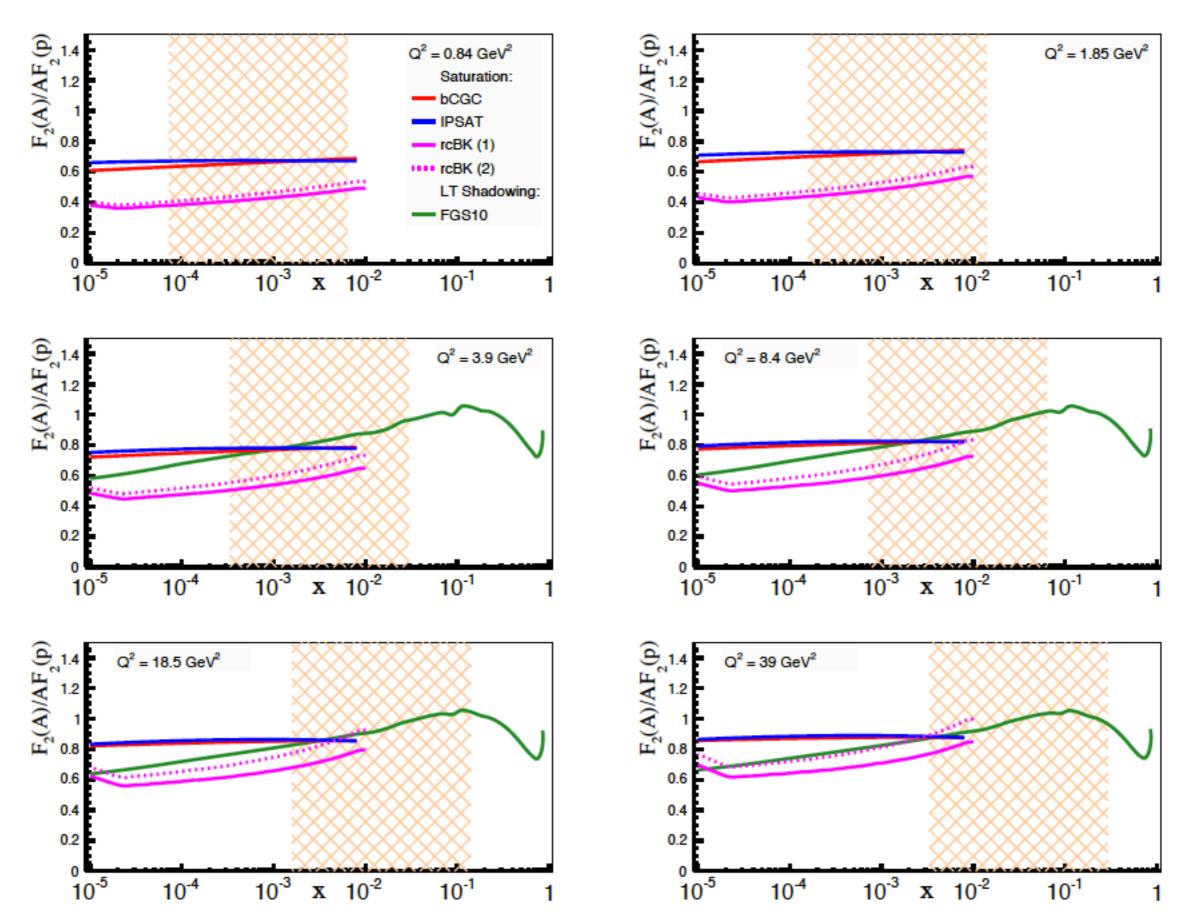
$F_2(A)/A - low Q^2$



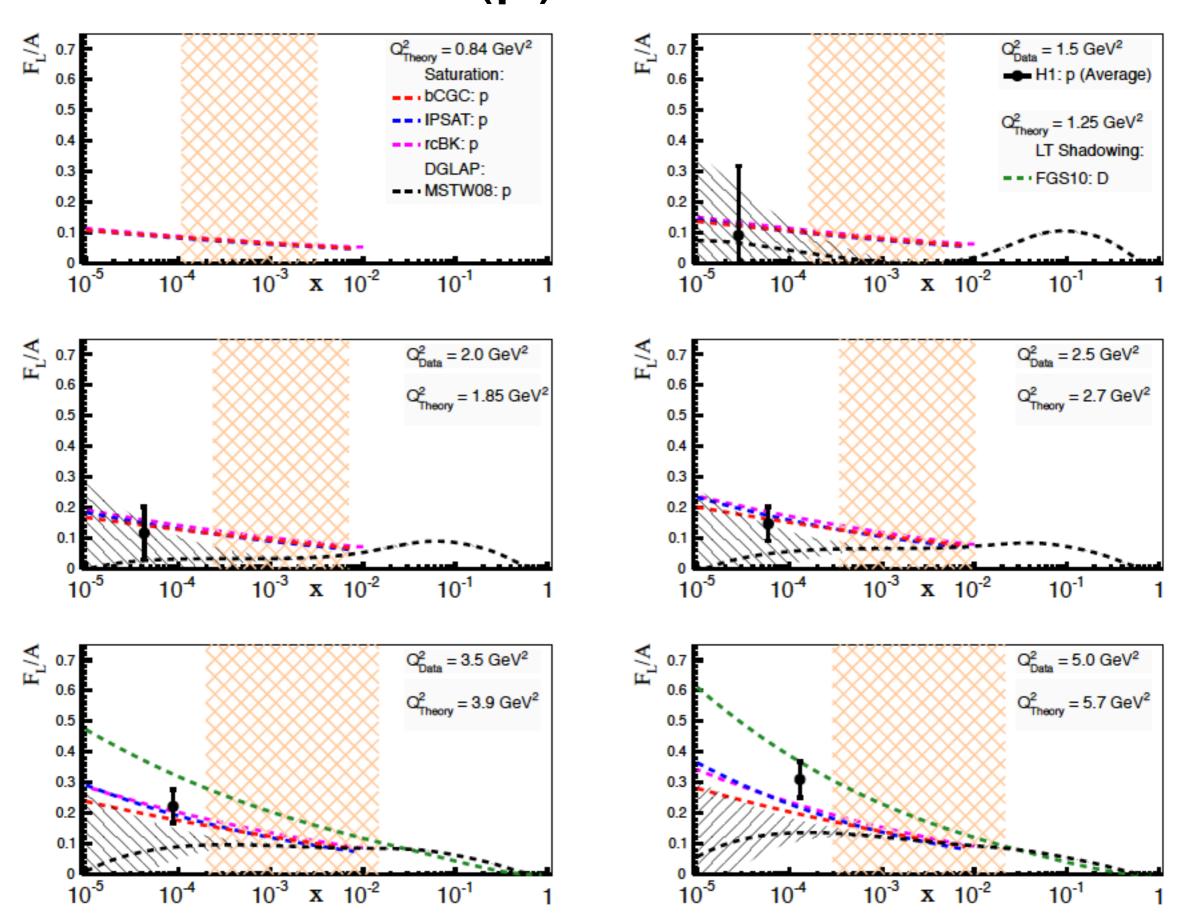
$F_2(A)/A - p vs A$



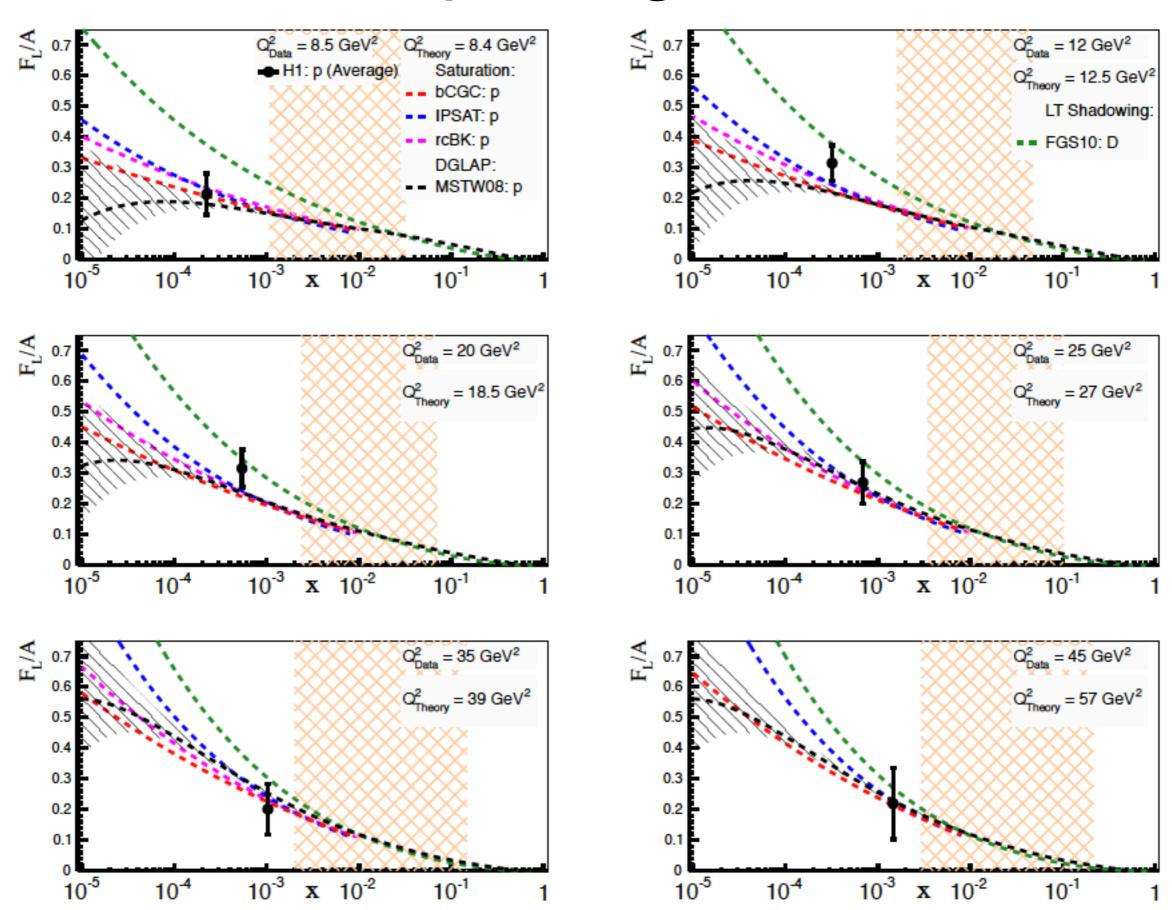
F_2 ratios: $F_2(A)/AF_2(p)$



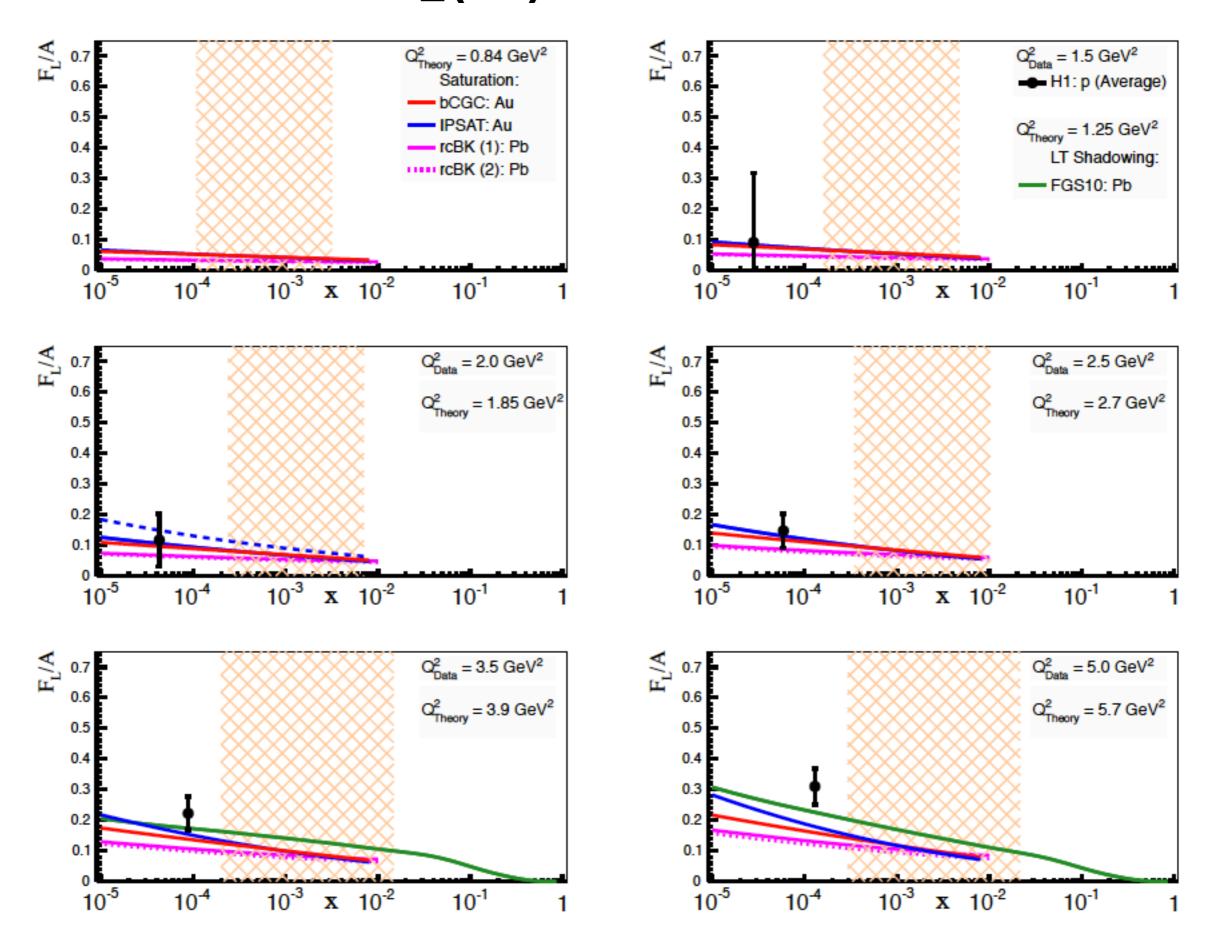
$F_L(p)$ - low Q^2



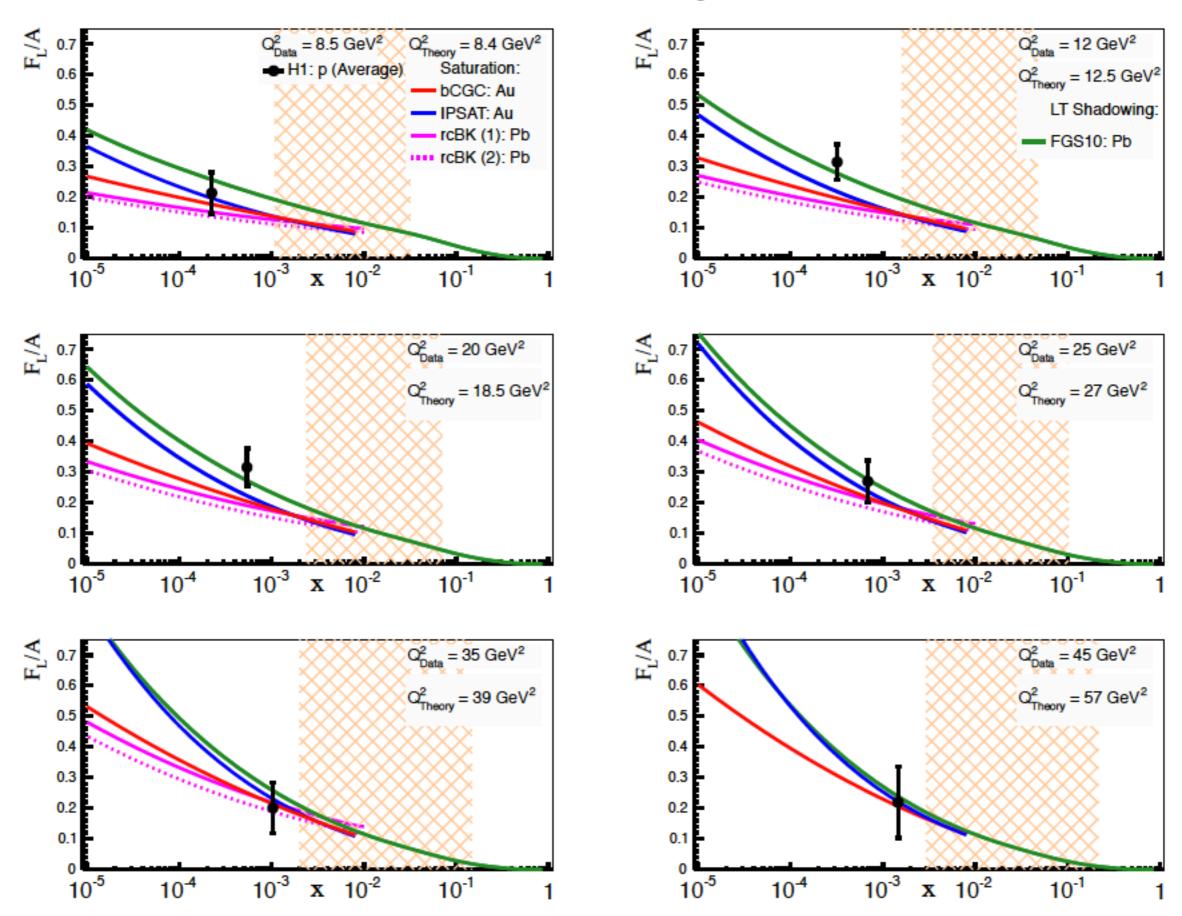
F_L(p) - higher Q²



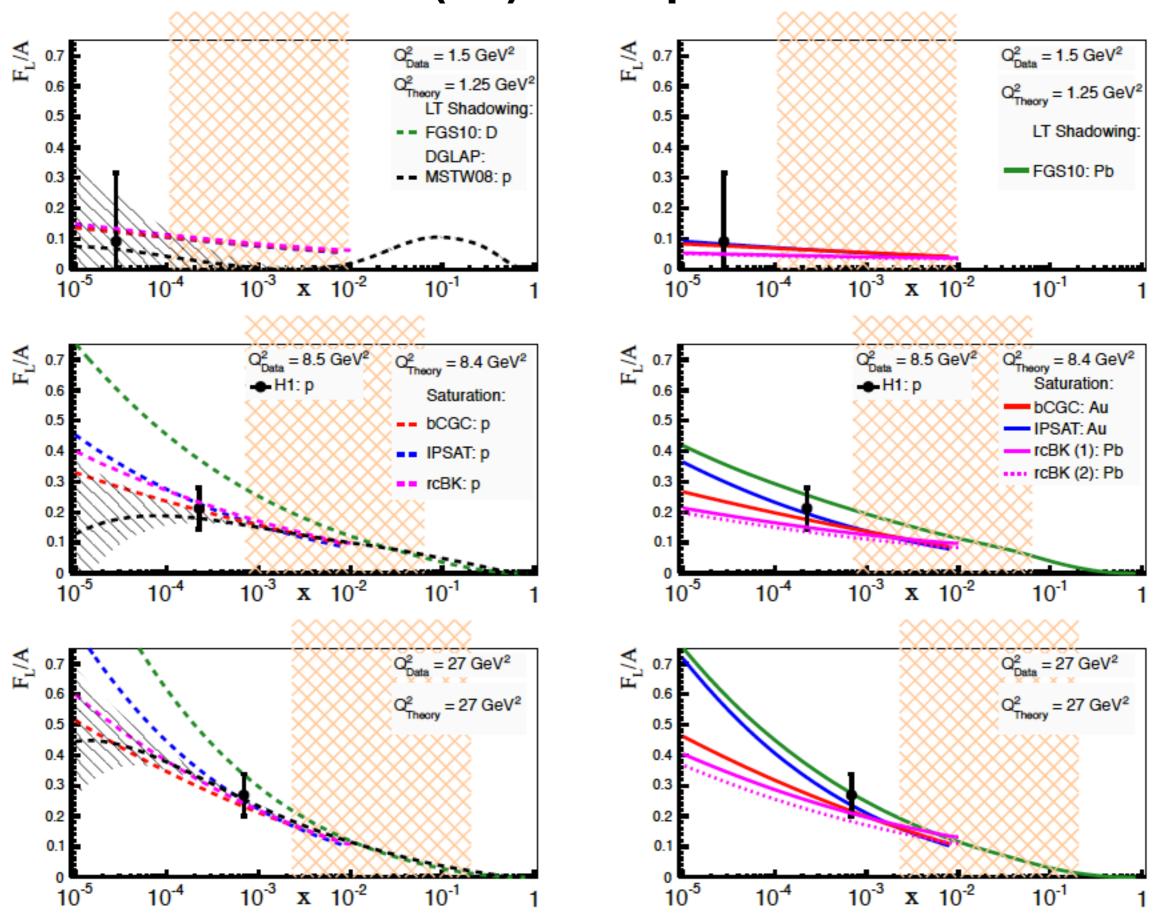
$F_L(A)/A - low Q^2$



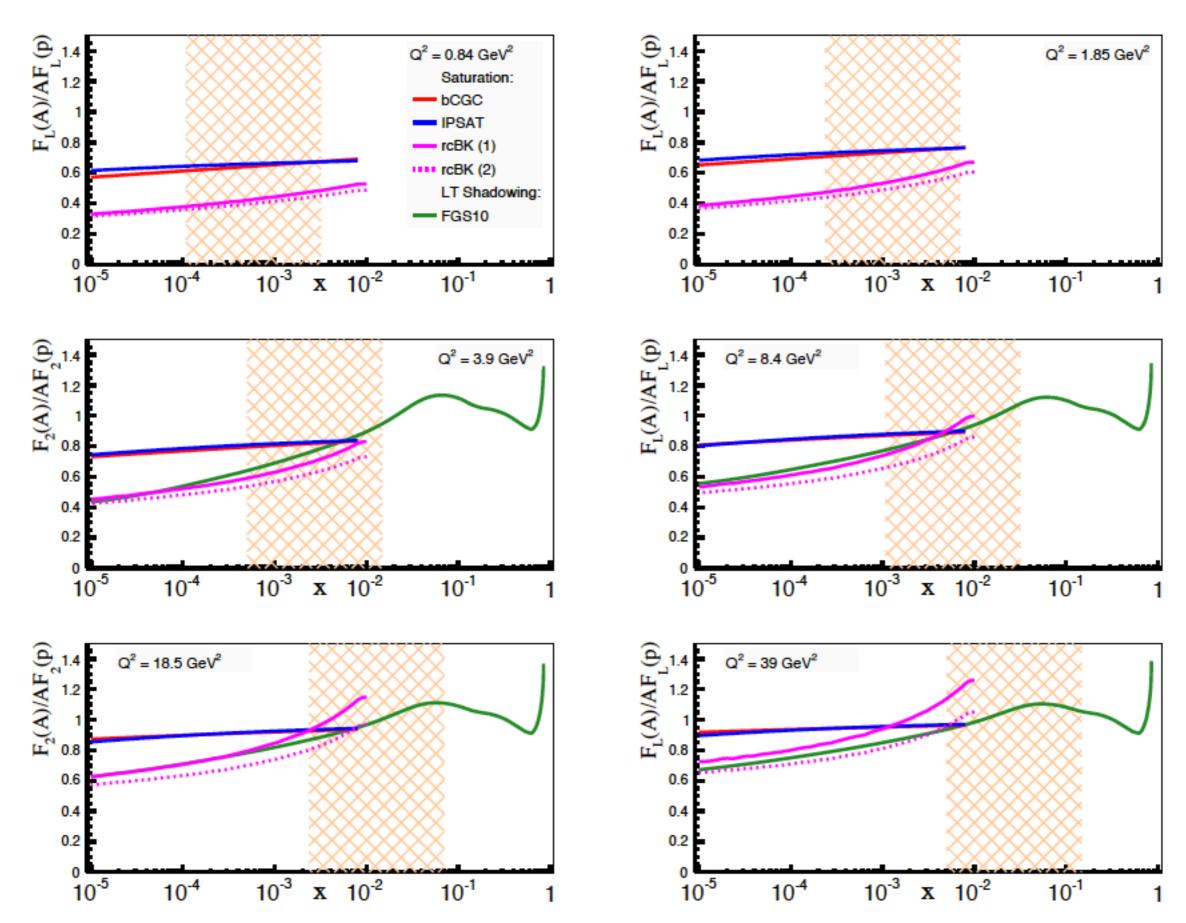
F_L(A)/A - higher Q²



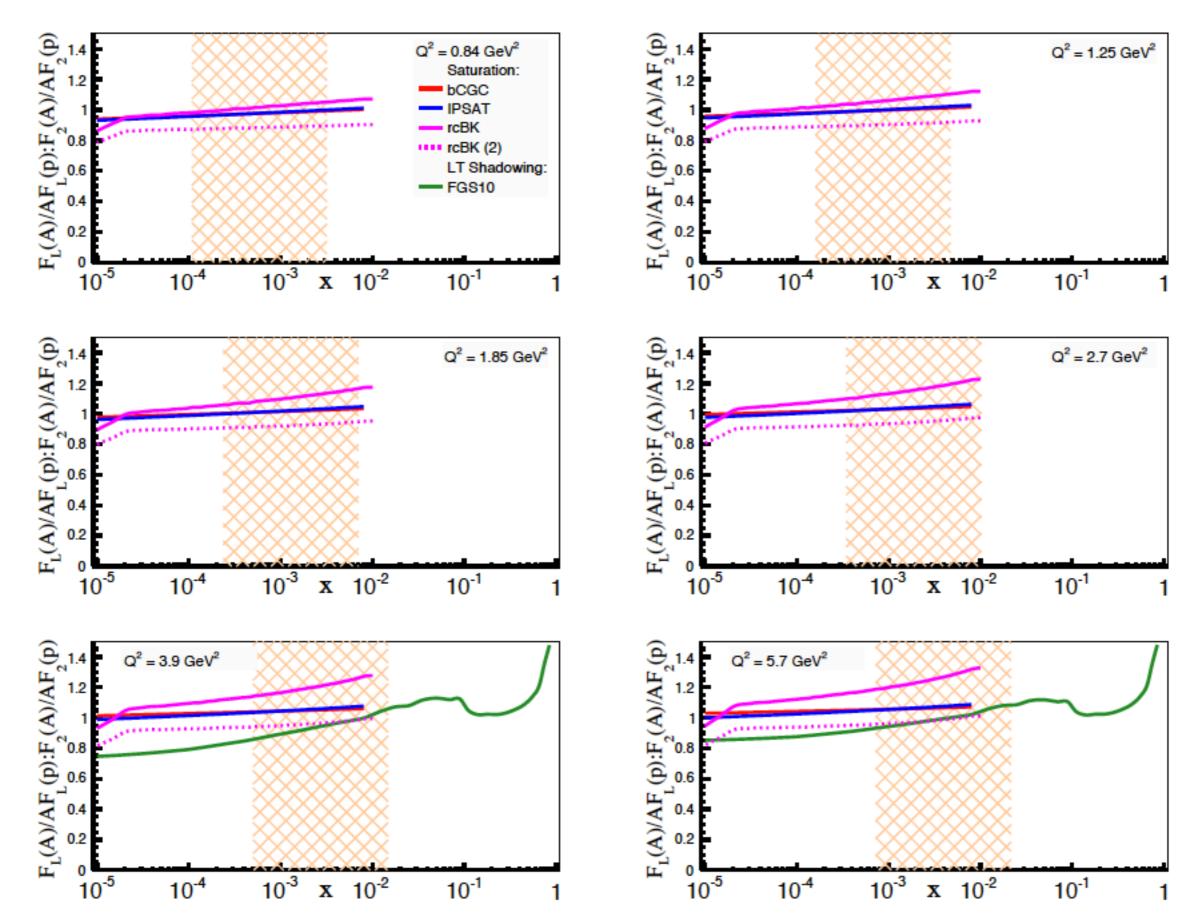
$F_L(A)/A - p vs A$



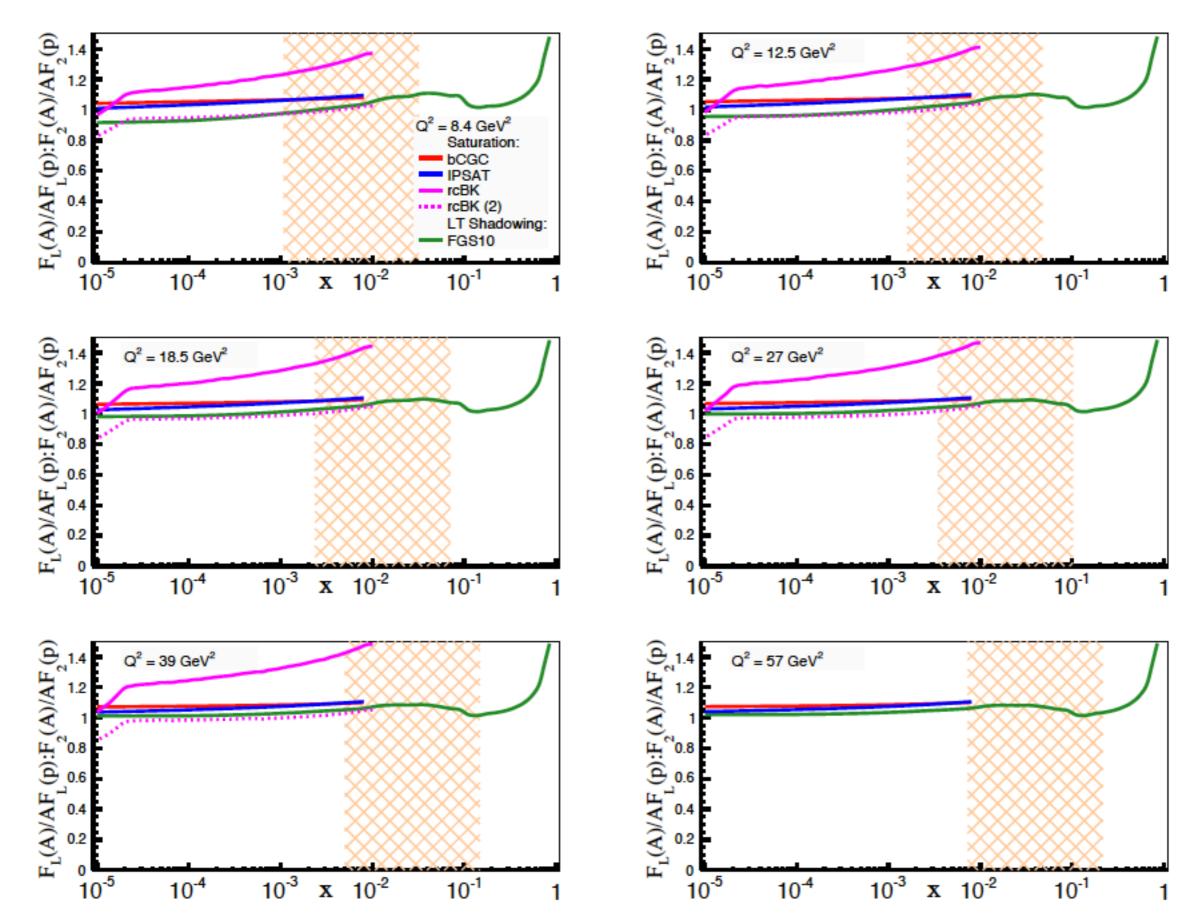
F_L ratios: $F_L(A)/AF_L(p)$



Double ratios - $F_L(A)/AF_L(p)$: $F_2(A)/AF_2(p)$



Double ratios - $F_L(A)/AF_L(p)$: $F_2(A)/AF_2(p)$

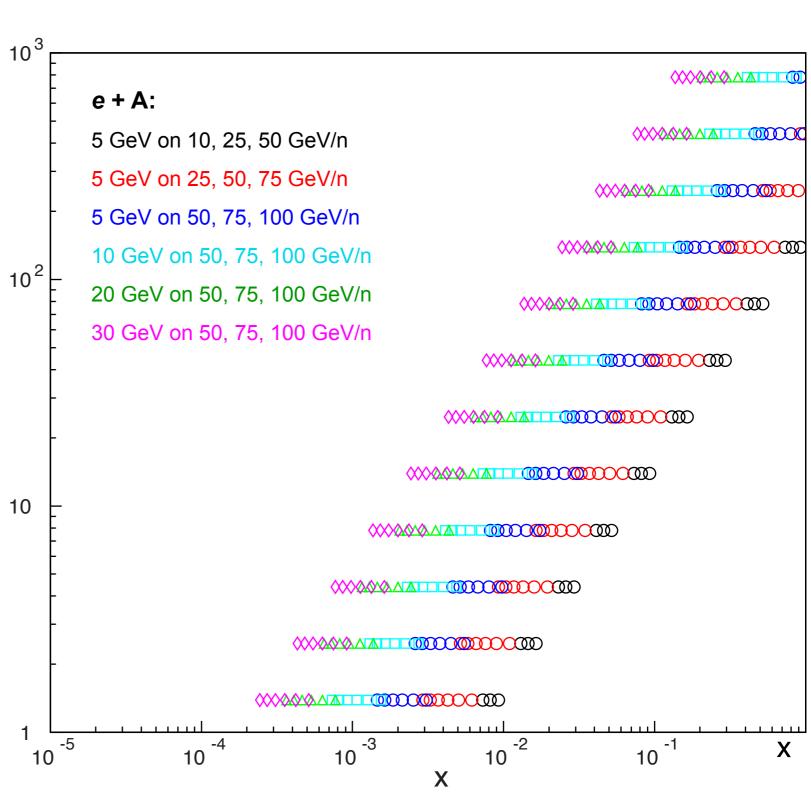


Feasibility study:
$$\sigma_r(x,Q^2) = F_2^A(x,Q^2) - \frac{y^2}{Y^+} F_L^A(x,Q^2)$$

- Simulated data for e+A coverage in x-Q² space
 - → 3 energies is the minimum requirement in the F_L capability study
 - → 1st stage only gets to medium x
 - Need high electron energy to get to "small" X

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Feasibility study: $\sigma_r(x,Q^2) = F_2^A(x,Q^2) - \frac{y^2}{Y^+} F_L^A(x,Q^2)$

Strategies:

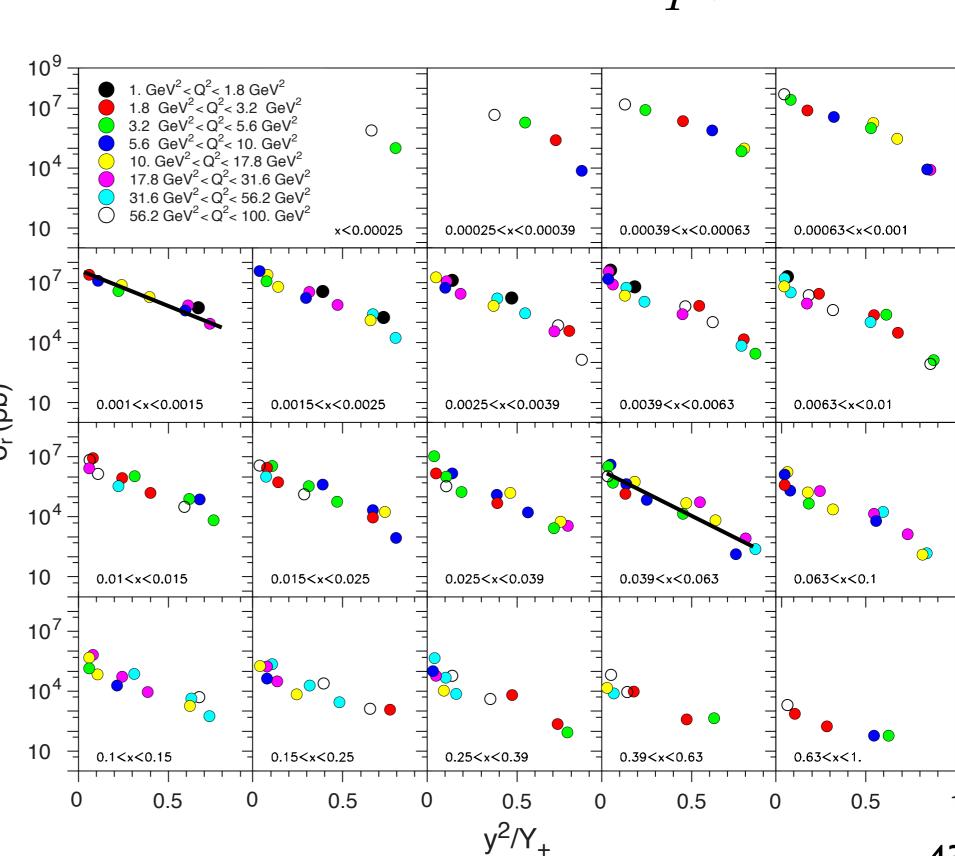
slope of y^2/Y_+ for different s at fixed x & \mathbb{Q}^2

e+p: 1st stage 5x50 - 5x325 running combined 4 weeks/each (50% eff)

stat. error shown and negligible

To Do:

refine method & test how well we can extract F_L in e+A collisions

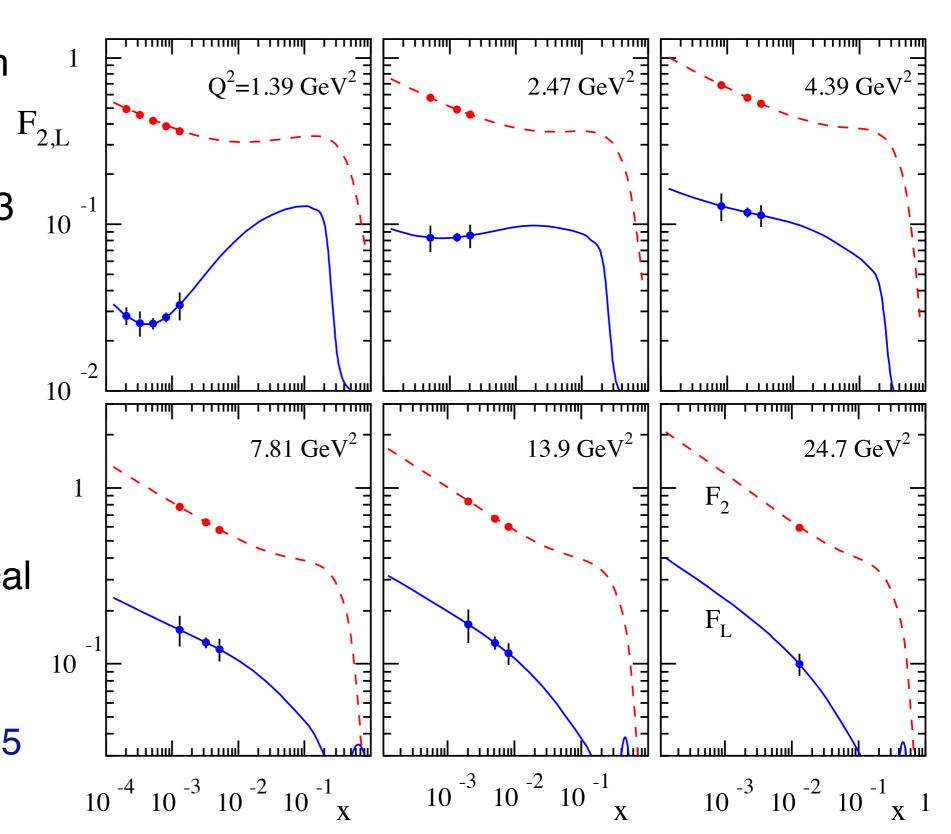


Extracting F₂ and F_L at the EIC

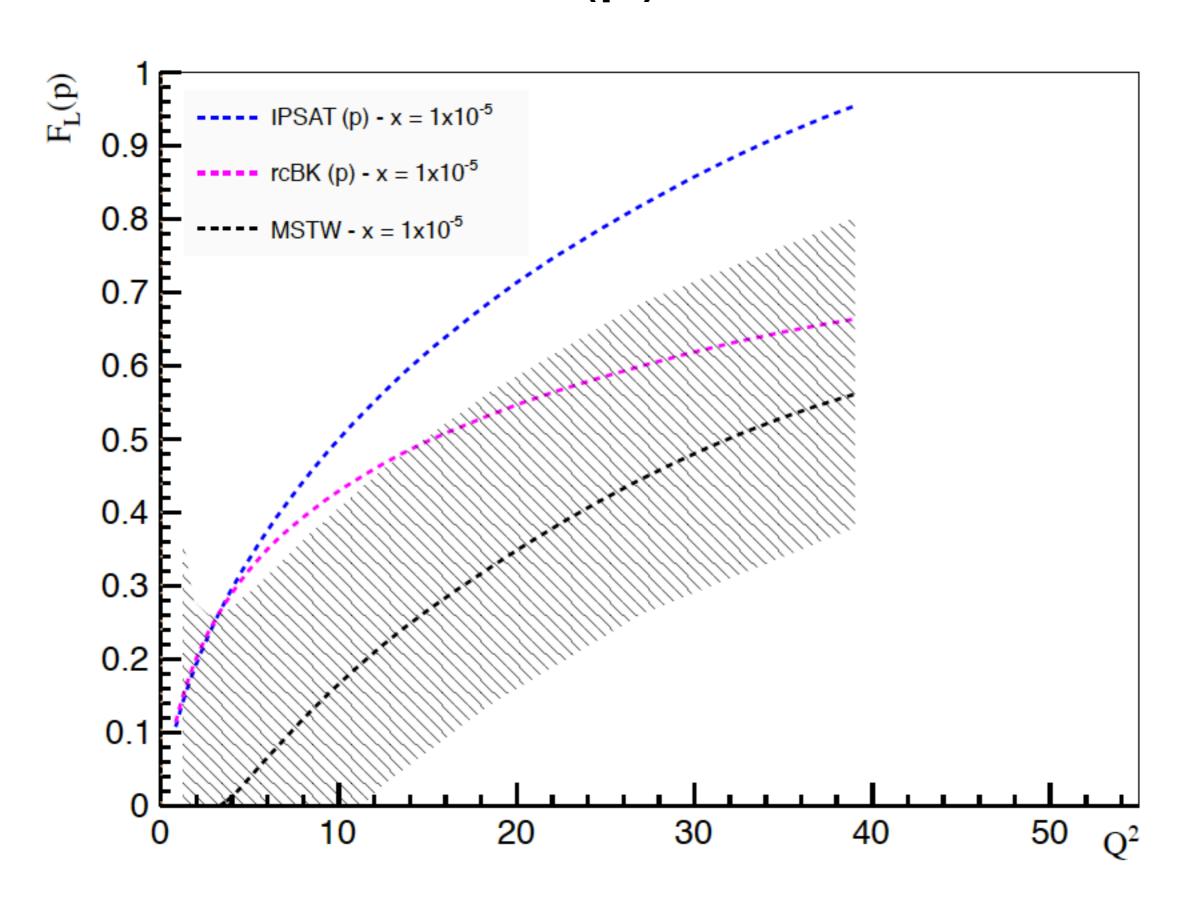
F_{2,L} extracted from pseudo-data
 generated for 1
 month running at 3
 eRHIC energies



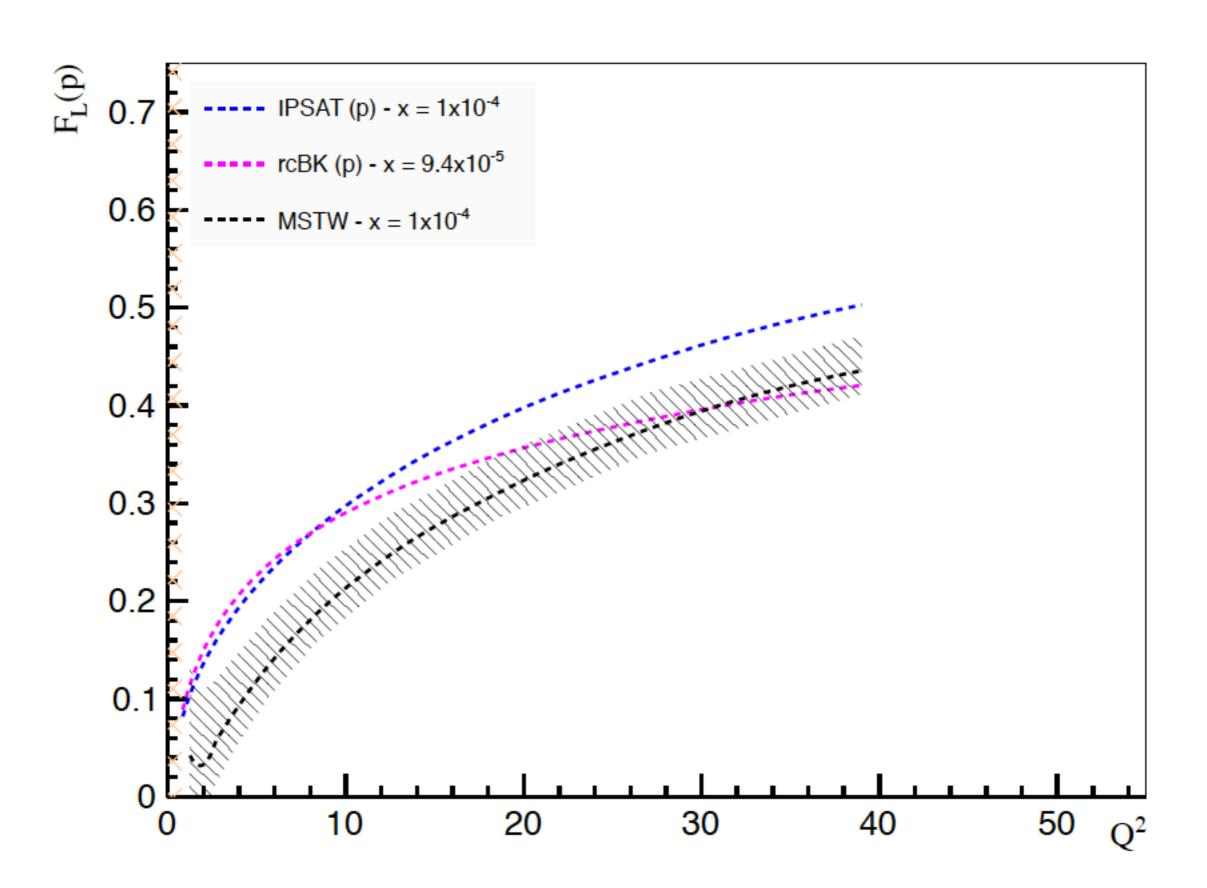
- → 5+250 GeV
- → 5+325 GeV
- Data, with errors, added to theoretical expectations from ABKM09 PDF set
 - → valid for Q² > 2.5 GeV²



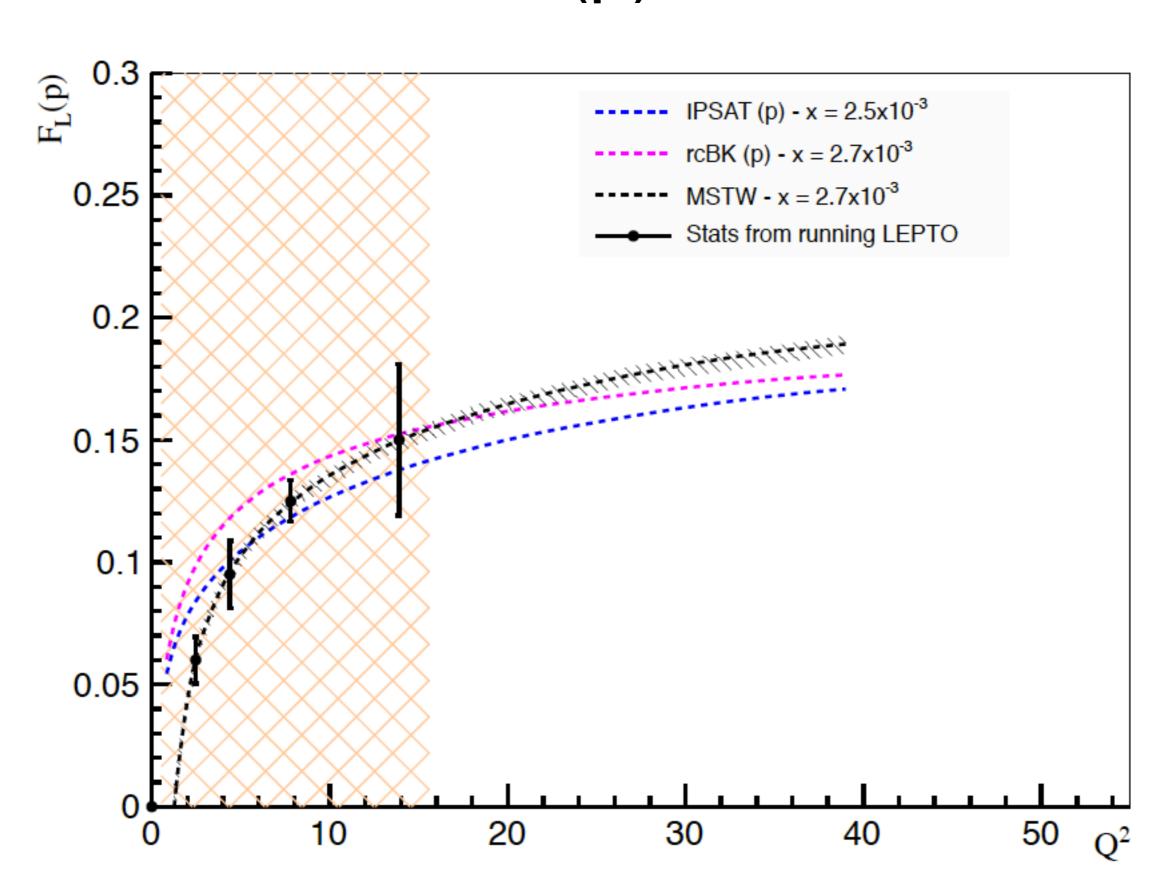
Evolution of $F_L(p)$ with Q^2 - fixed x



Evolution of $F_L(p)$ with Q^2 - fixed x

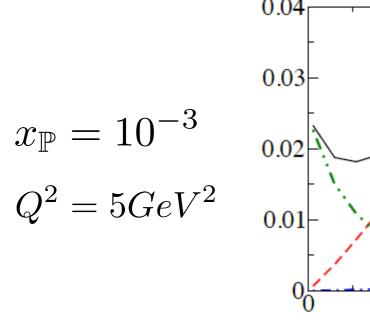


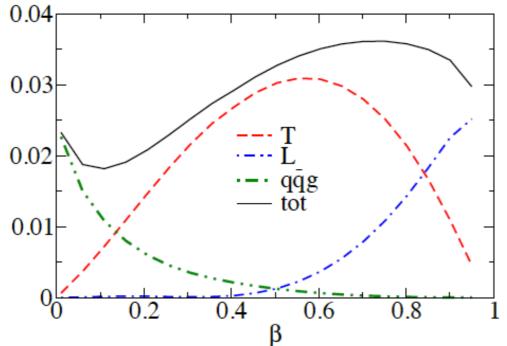
Evolution of $F_L(p)$ with Q^2 - fixed x

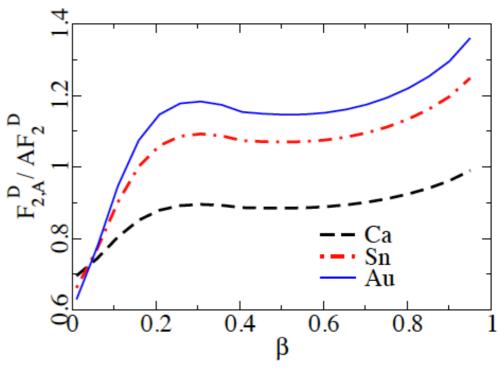


Charm and diffractive structure functions, FD2,L, FC2,L

- Fc_{2,L} give more direct access to the gluon distribution than the inclusive F₂ structure function
 - → QCD calculations with non-zero m_c are scheme dependent and can absorb saturation signals if not handled correctly
- FD_{2,L} is also sensitive to the gluon distribution
 - → Differences between linear and non-linear models appear at higher Q² than for F₂ (8 GeV² vs 2 GeV²)
 - ► More experimentally challenging measurement than F₂





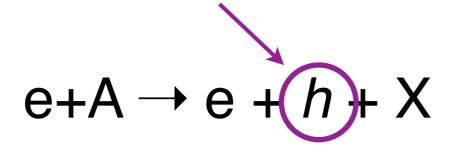


k_T dependent gluons, gluon correlations from di-hadron correlations, SIDIS (semi-inclusive DIS)

k_T dependent gluons, gluon correlations from di-hadron correlations, SIDIS (semi-inclusive DIS)

Deliverables	Observables	What we learn	Stage-I	Stage-II
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Direct link between p_T of produced hadron and that of the small-x gluon



k _⊤ dependent gluons	SIDIS at small x	non-linear QCD evolution / universality	onset of saturation	rare probes and bottom; large-x gluons
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k_T dependent gluons, gluon correlations from di-hadron correlations, SIDIS (semi-inclusive DIS)

$$e+A \rightarrow e + h_1 + h_2 + X$$

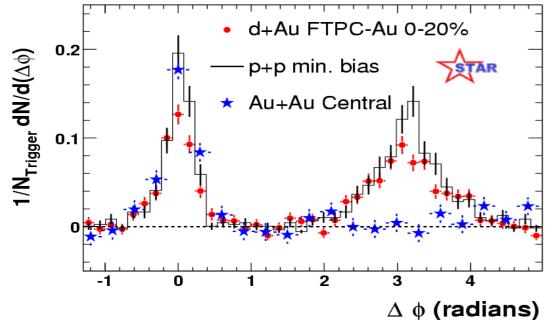
Deliverables	Observables	What we learn	Stage-I	Stage-II
k _T dependent gluons; gluon correlations	di-hadron correlations	non-linear QCD evolution / universality	onset of saturation	measure Q _s

$$e+A \rightarrow e + h + X$$

k⊤ dependent gluons	SIDIS at small x	non-linear QCD evolution / universality	onset of saturation	rare probes and bottom; large-x gluons
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di-hadron angular correlations in d+A

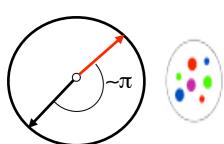
comparisons between d+Au \rightarrow h₁ h₂ X (or p $+Au \rightarrow h_1 h_2 X$) and $p+p \rightarrow h_1 h_2 X$



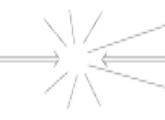
- At y=0, suppression of awayside jet is observed in A+A collisions
- No suppression in p+p or d+A

→
$$x \sim 10^{-2}$$

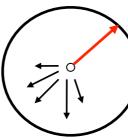
$$x_A = \frac{k_1 e^{-y_1} + k_2 e^{-y_2}}{\sqrt{s}} << 1$$





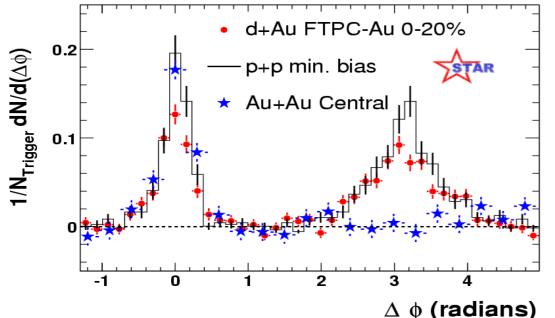






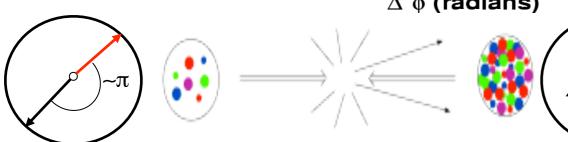
di-hadron angular correlations in d+A

comparisons between d+Au $\to h_1 \; h_2 \; X$ (or p +Au $\to h_1 \; h_2 \; X$) and p+p $\to h_1 \; h_2 \; X$



- At y=0, suppression of awayside jet is observed in A+A collisions
- No suppression in p+p or d+A

$$x_A = \frac{k_1 e^{-y_1} + k_2 e^{-y_2}}{\sqrt{s}} << 1$$



 $<\eta_L>=3.2, <\eta_S>=3.1$

 $\frac{p+p \rightarrow \pi^0 \pi^0 + X, \sqrt{s} = 200 \text{ GeV}}{p_{TL} > 2 \text{ GeV/c}, 1 \text{ GeV/c} < p_{TS} < p_{TL}}$

0.015

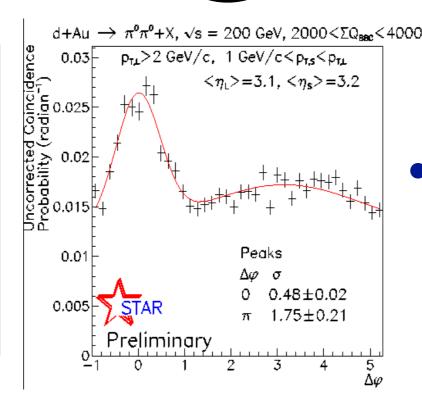
Ā0.0125

0.01

0.005

0.0075

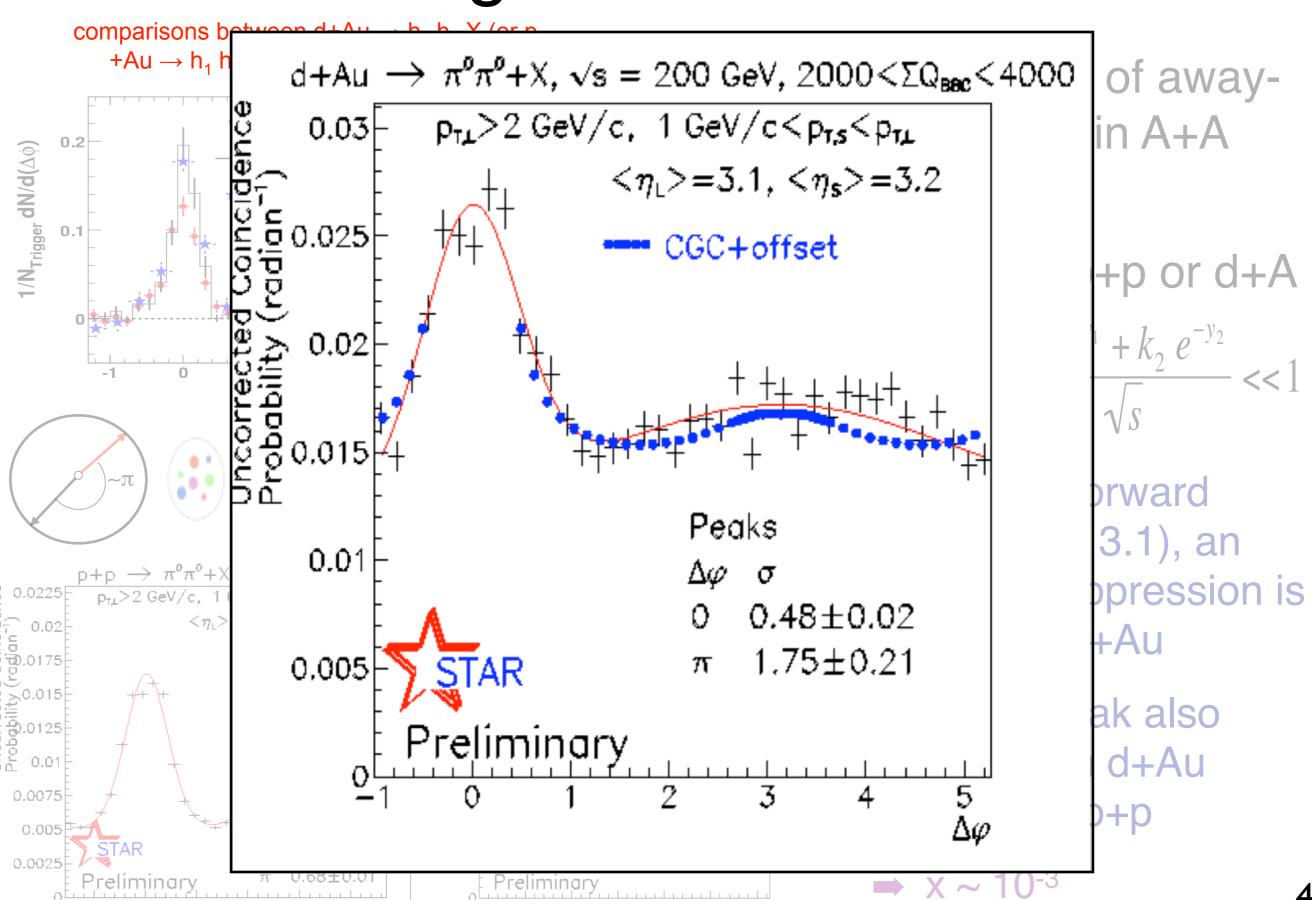
0.0025



- However, at forward rapidities (y ~ 3.1), an away-side suppression is observed in d+Au
- Away-side peak also much wider in d+Au compared to p+p

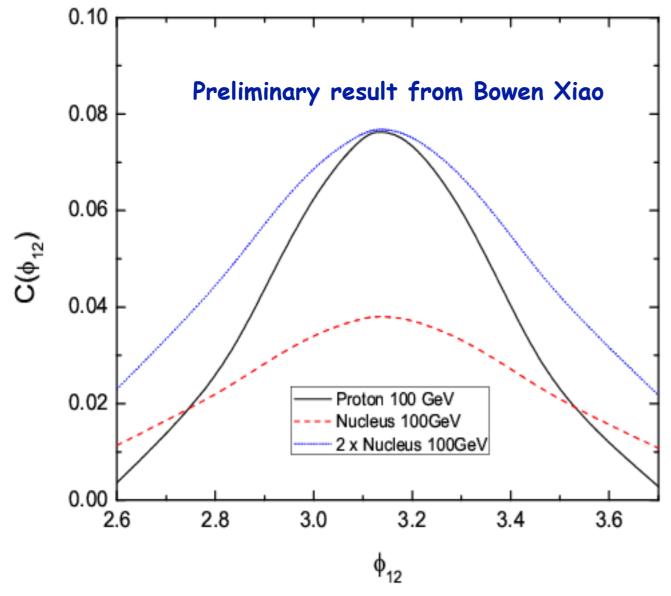
→
$$X \sim 10^{-3}$$

di-hadron angular correlations in d+A



di-hadron correlations in e+A

Never been measured - we expect to see the same effect in e+A as in d+A



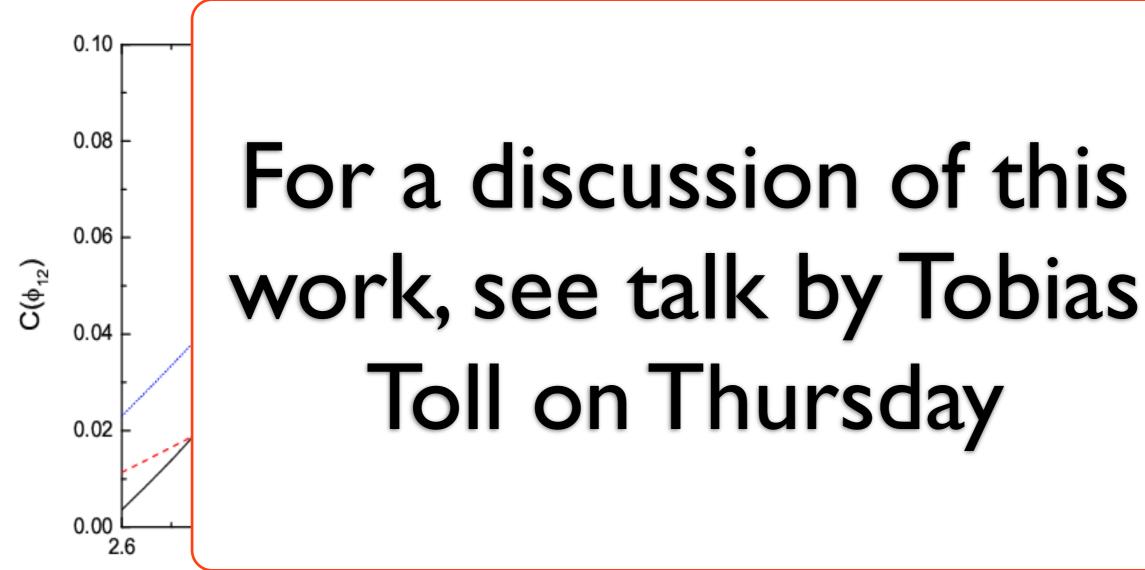
$$Q^2 = 4GeV^2$$
; $z_{h1} = z_{h2} = 0.3$
2 GeV < p_T^T < 3GeV
1GeV < p_T^A < 2GeV

- At small-x, multi-gluon distributions are as important as single-gluon distributions and they contribute to dihadron correlations
 - → The non-linear evolution of multi-gluon distributions is different from that of singlegluon distributions and it is equally important that we understand it
- The d+Au RHIC data is therefore subject to many uncertainties
 - these correlations in e+A can help to constrain them better

di-hadron correlations in e+A

Never been measured - we expect to see the same effect in e+A as in d+A

 At small-x, multi-gluon distributions are as important



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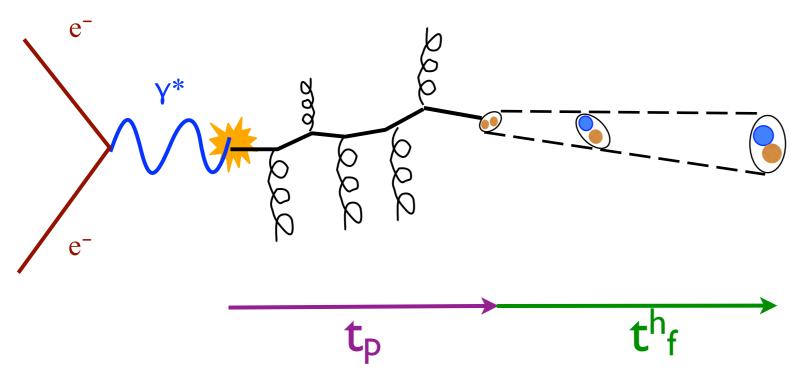
ny

 $Q^2 = 4GeV^2$; $z_{h1} = z_{h2} = 0.3$ 2 GeV < p_T^T < 3GeV 1GeV < p_T^A < 2GeV these correlations in e+A can help to constrain them better transport coefficients in cold nuclear matter from large-x semi-inclusive DIS and jets

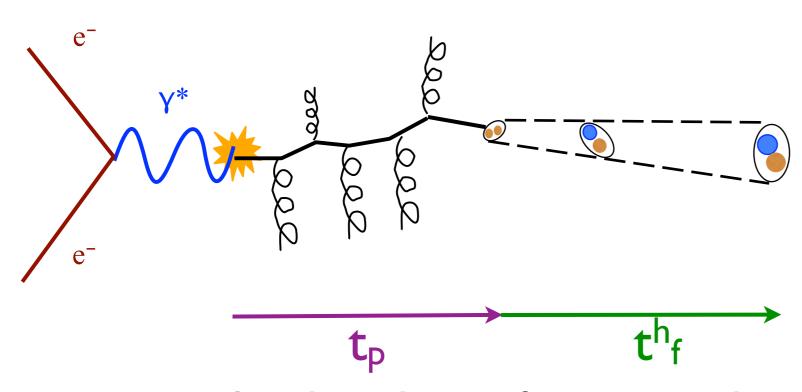
Transport coefficients in cold nuclear matter

Deliverables	Observables	What we learn	Stage-I	Stage-II
transport coefficients in cold matter	large-x SIDIS; jets	parton energy loss, shower evolution; energy loss mechanisms	light flavours and charm; jets	rare probes and bottom; large-x gluons

A. Accardi R. Dupre



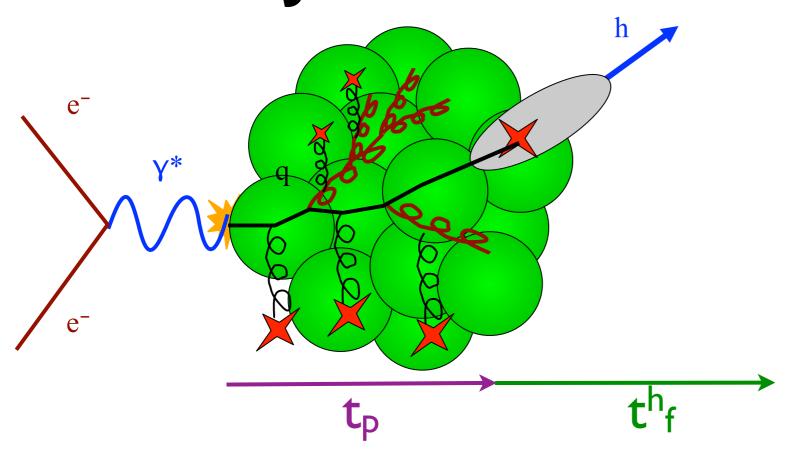
- tp production time of propagating quark
- th_f hadron formation time



A.Accardi R. Dupre What happens if we add a nuclear

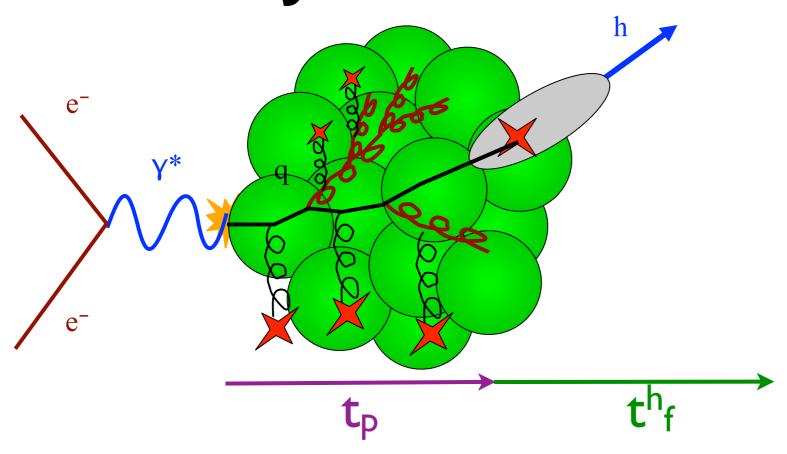
medium?

- tp production time of propagating quark
- th_f hadron formation time



A.Accardi R. Dupre What happens if we add a nuclear medium?

- tp production time of propagating quark
- th_f hadron formation time

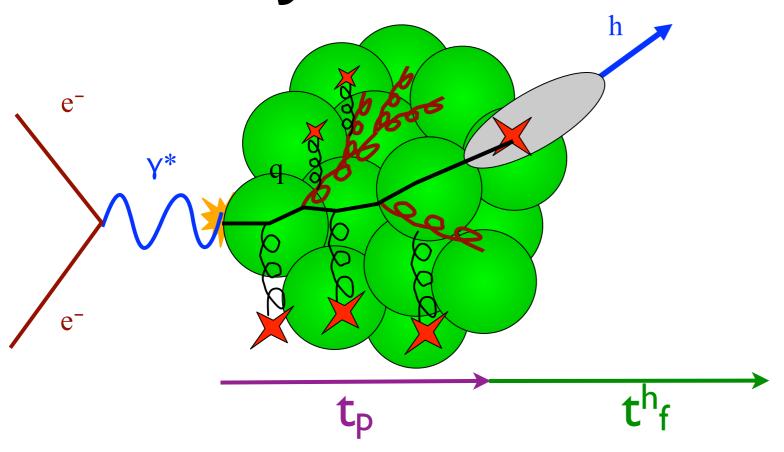


A.Accardi R. Dupre What happens if we add a nuclear medium?

- tp production time of propagating quark
- th_f hadron formation time

Observables:

Broadening: $\Delta p_T^2 = \left\langle p_T^2 \right\rangle_A - \left\langle p_T^2 \right\rangle_p$; direct link to saturation scale



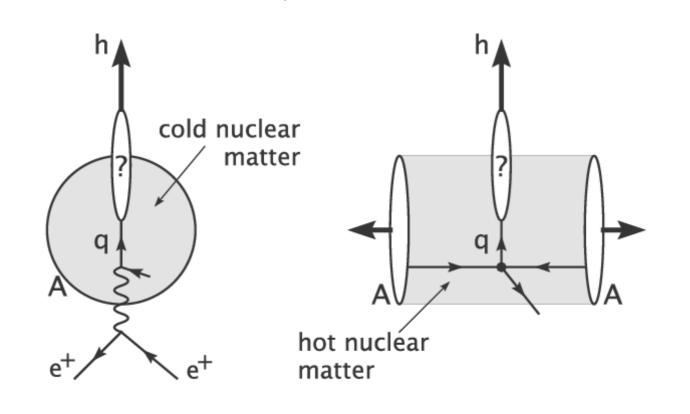
A.Accardi R. Dupre What happens if we add a nuclear medium?

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Broadening: $\Delta p_T^2 = \left\langle p_T^2 \right\rangle_A - \left\langle p_T^2 \right\rangle_p$ direct link to saturation scale

Attenuation: $R_A^h(Q^2, \nu, z_h, p_T^2)$: ratio of hadron production in A to D, modifications of nPDFs cancel out



A. Accardi R. Dupre

- t_p production time of propagating quark
- th_f hadron formation time

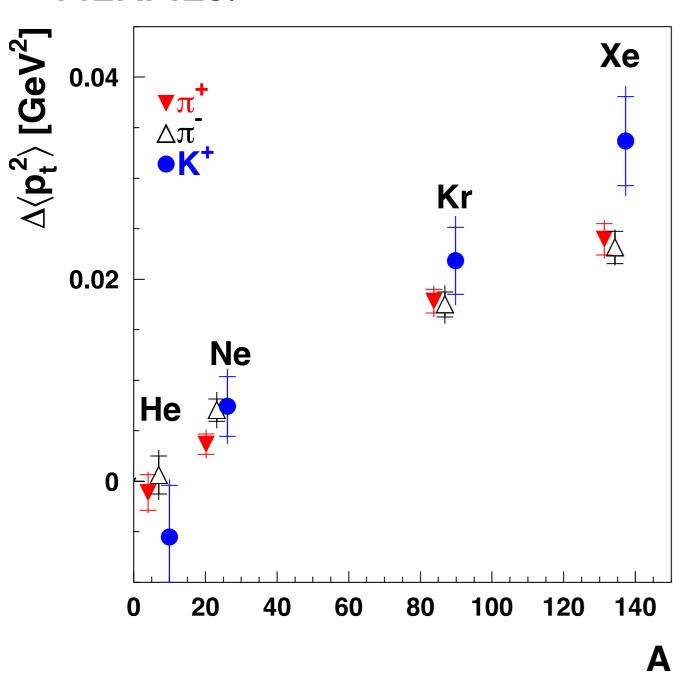
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pt broadening - how can the EIC contribute?

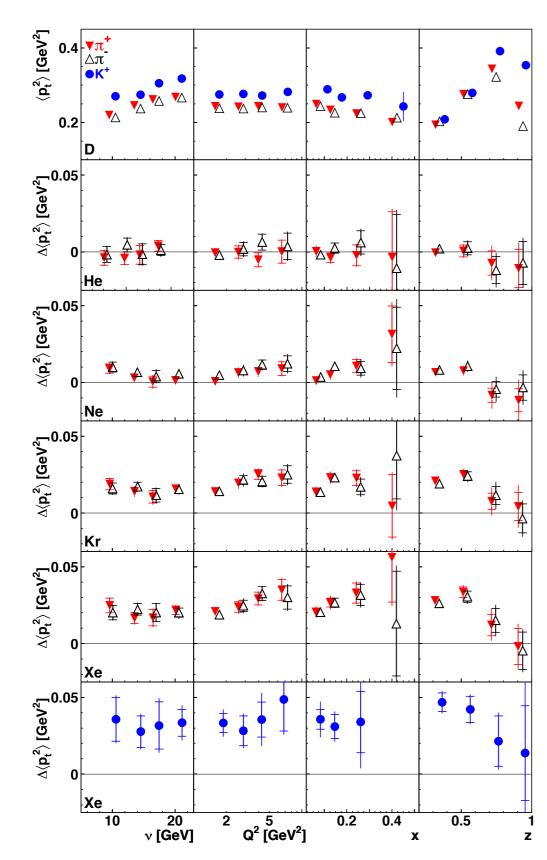
HERMES:



Increase of p_T broadening seen with increasing nuclear size - integrated over all variables

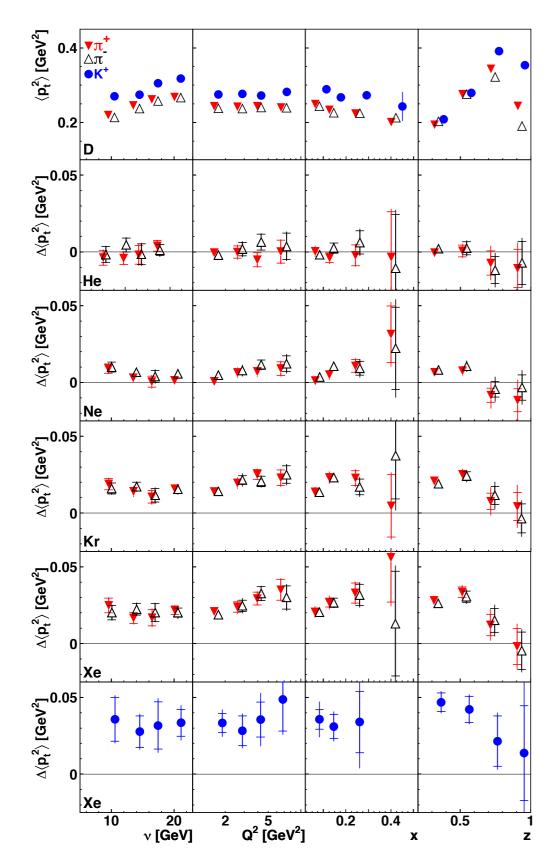
p_T broadening - how can the EIC contribute?

HERMES:

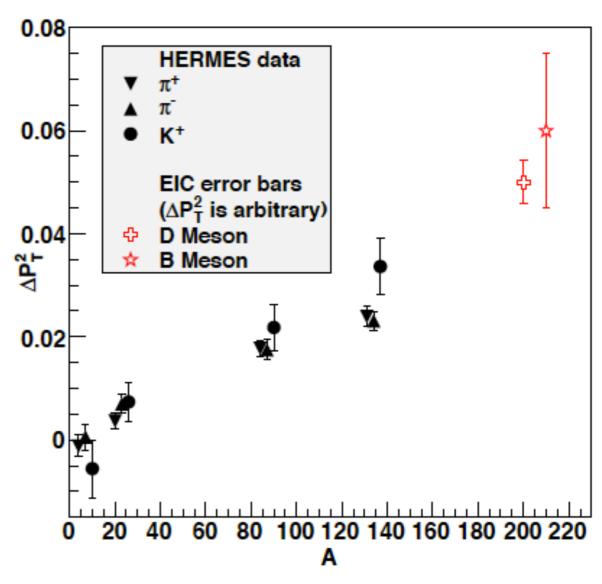


pt broadening - how can the EIC contribute?

HERMES:



EIC:



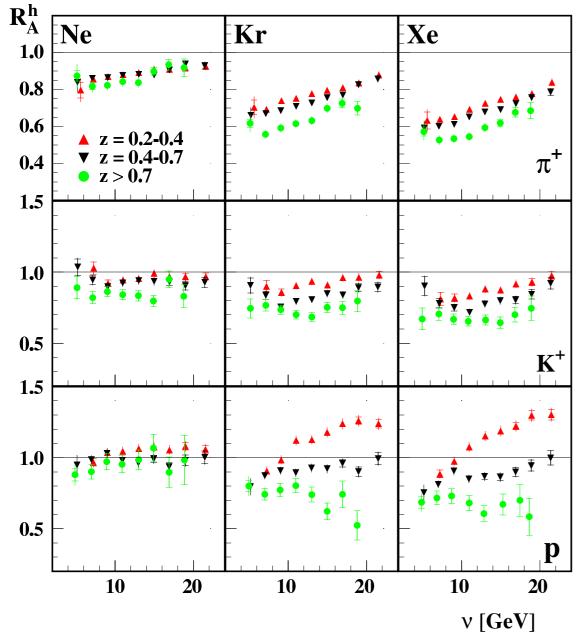
Measurements from HERMES can be repeated, with the addition of heavy quarks

Attenuation - how can the EIC contribute?

HERMES:

$$E_e = 27 \text{ GeV} \rightarrow \sqrt{s} = 7.2 \text{ GeV}$$

$$E_h = 2-15 \text{ GeV}$$



v = virtual photon energy

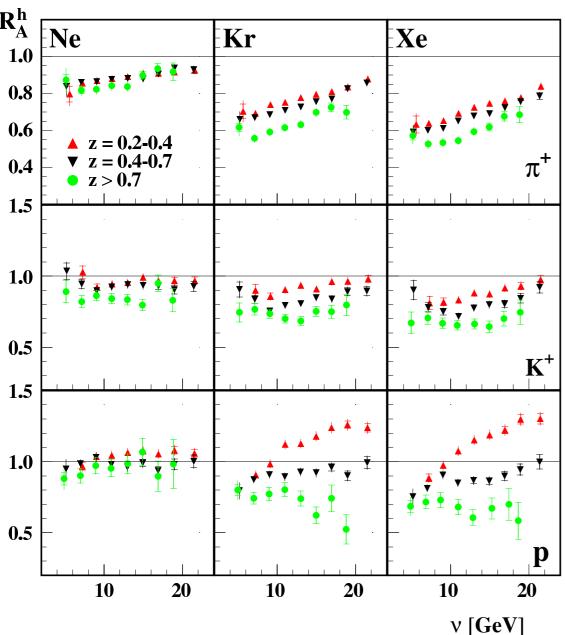
$$Z_h = E_h/v$$

Attenuation - how can the EIC contribute?

HERMES:

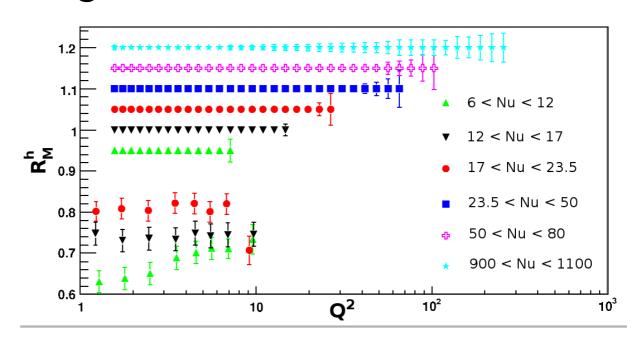
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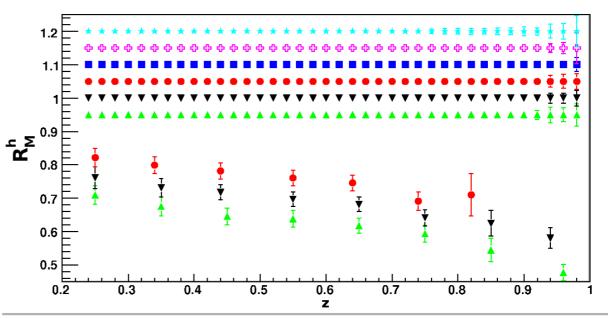
$$E_h = 2-15 \text{ GeV}$$



EIC:

light hadrons:





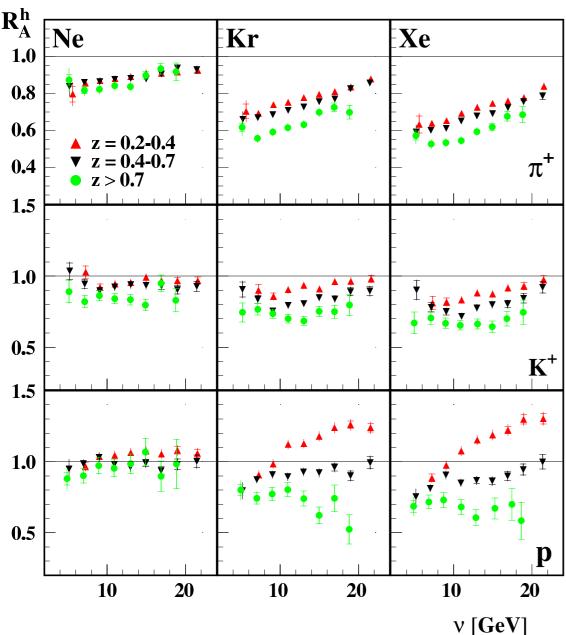
 $V = virtual photon energy large V range \rightarrow boost$ hadronization in and out of nucleus ⁵⁵ $Z_h = E_h/v$

Attenuation - how can the EIC contribute?

HERMES:

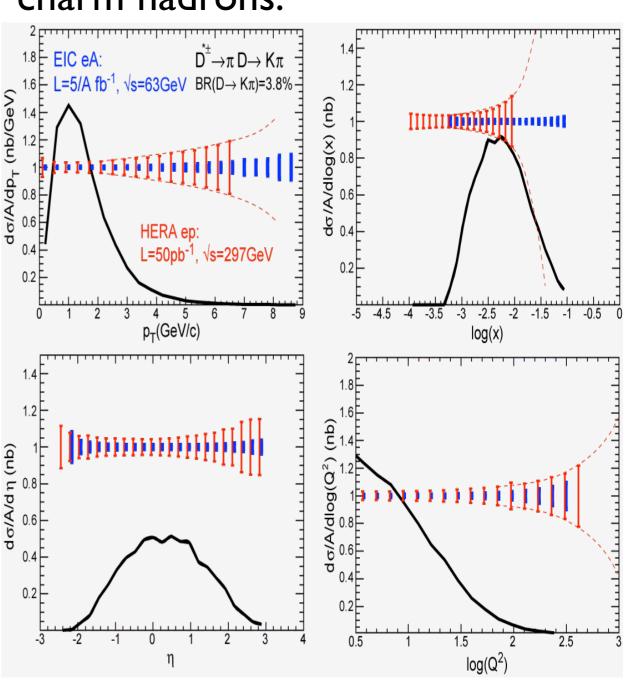
$$E_e = 27 \text{ GeV} \rightarrow \sqrt{s} = 7.2 \text{ GeV}$$

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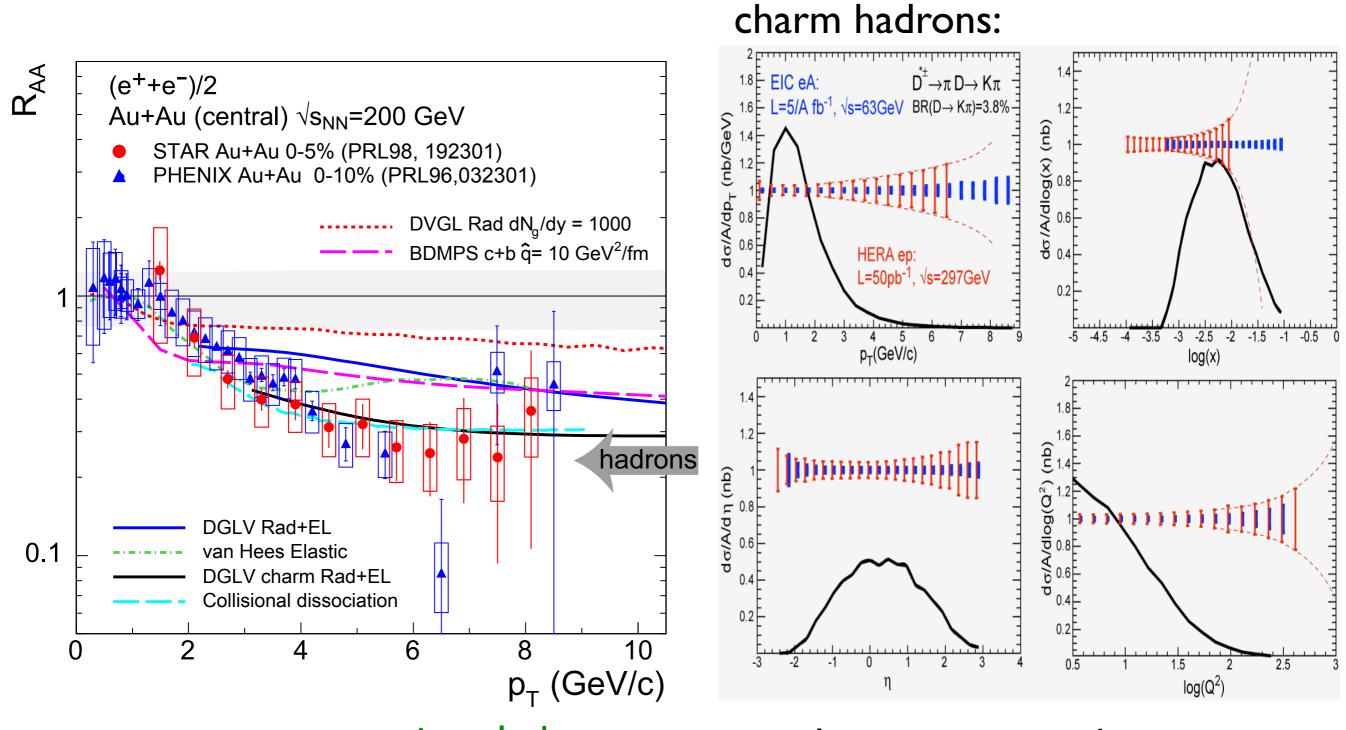
EIC:

charm hadrons:



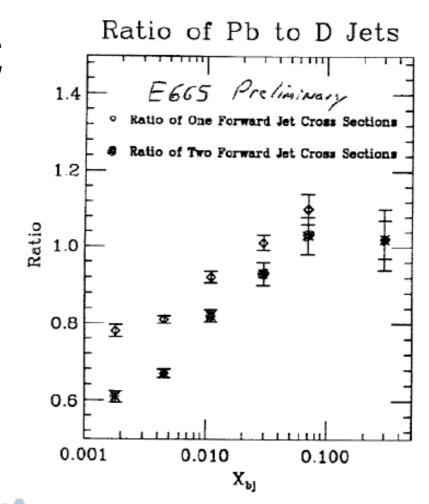
 $V = virtual photon energy large V range \rightarrow boost$ $Z_h = E_h/V \qquad hadronization in and out of nucleus 55$

Attenuation - how can the EIC contribute? RHIC: EIC:



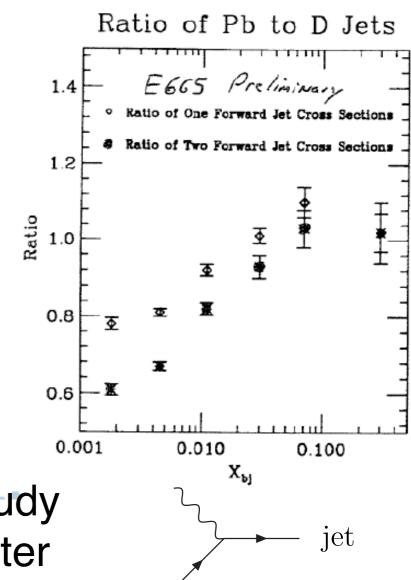
 ν = virtual photon energy large ν range \rightarrow boost $Z_h = E_h/\nu$ hadronization in and out of nucleus 55

- E665 at FNAL have measured jets in μ +A at $\sqrt{s} \sim 30 \text{ GeV}$
 - Feasible to start a jet programme in phase 1
 - caveat that collider kinematics are different to fixed target

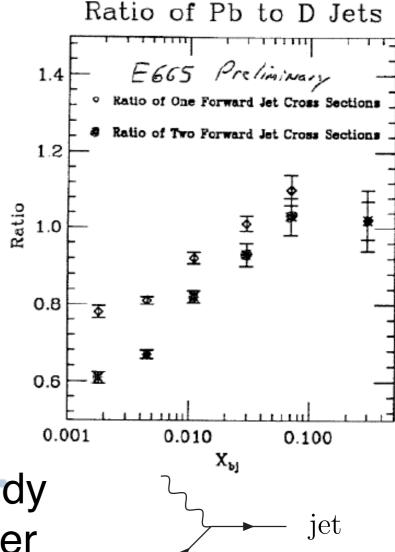


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1+1 jets, dominated by q processes → allow study of parton propagation through cold nuclear matter



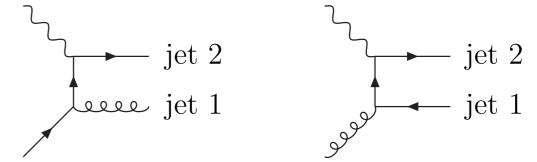
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1+1 jets, dominated by q processes → allow study of parton propagation through cold nuclear matter

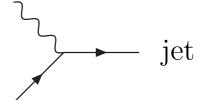
$$\frac{d^2\sigma_{2+1}}{dxdQ^2} = A_q(x, Q^2)q^A(x, Q^2) + A_g(x, Q^2)g_A(x, Q^2)$$

2+1 jets→ sensitive to nuclear gluons



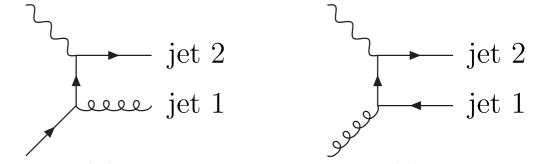
By measuring 1+1 jets, can extract information on gluons

1+1 jets, dominated by q processes → allow study of parton propagation through cold nuclear matter

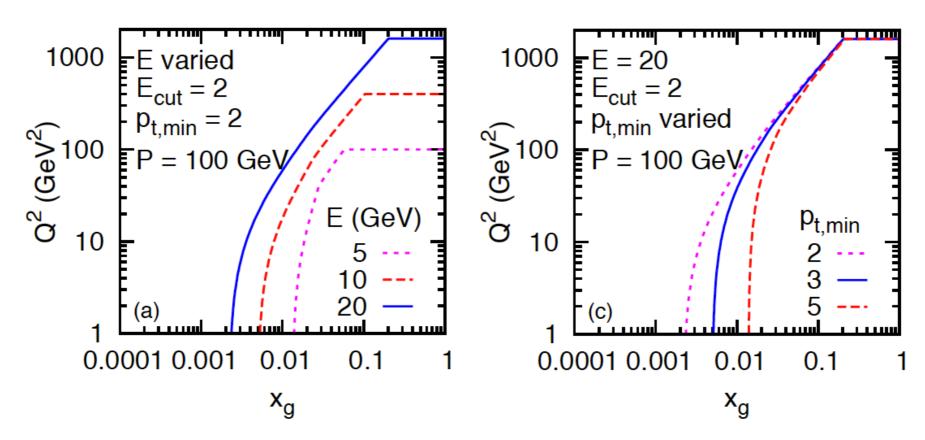


$$\frac{d^2\sigma_{2+1}}{dxdQ^2} = A_q(x, Q^2)q^A(x, Q^2) + A_g(x, Q^2)g_A(x, Q^2)$$

2+1 jets→ sensitive to nuclear gluons



By measuring 1+1 jets, can extract information on gluons



b dependent gluons, gluon correlations from DVCS and diffractive vector meson production

Silver Measurements

Deliverables	Observables	What we learn	Stage-I	Stage-II
b-dependent gluons; gluon correlations	DVCS; diffractive vector mesons	interplay between small- x evolution and confinement	moderate x with light, heavy nuclei	smaller x, saturation

See talk by Tobias Toll on Thursday afternoon

Summary and Conclusions

- The e+p physics programme at an EIC will allow us to study in detail the spin structure of the proton
 - \rightarrow Unprecedented coverage down to small-x, provide constraints on Δg and Δq
- The e+A physics programme at an EIC will give us an unprecedented opportunity to study gluons in nuclei
 - Low-x: Measure the properties of gluons where saturation is the dominant governing phenomena
 - → Higher-x: Understand how fast partons interact as they traverse nuclear matter and provide new insight into hadronization
- Understanding the role of gluons in nuclei is crucial to understanding RHIC (and LHC) heavy-ion results

Good headway can be made on these measurements already with a stage-I eRHIC ($E_e = 5 \text{ GeV}$)

- The INT programme in the Fall of 2010 allowed us to formulate the observables in terms of golden and silver measurements
 - A detailed write-up of the whole programme (encompassing both e+A and e+p) is now published!!