

Transverse (Spin) Structure of Hadrons

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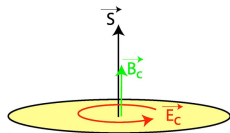
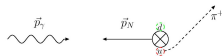
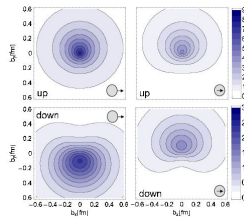
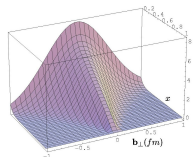
February 2, 2012

- Deeply virtual Compton scattering (DVCS)
- ↳ Generalized parton distributions (GPDs)
- ↳ 'transverse imaging'
- Chromodynamik lensing and \perp single-spin asymmetries (SSA)

transverse distortion of PDFs
+ final state interactions } \Rightarrow

↳ SSA in $\gamma N \rightarrow \pi + X$

- quark-gluon correlations $\rightarrow \perp$ force on q in DIS
- Summary

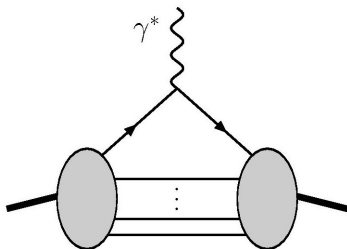


3 D imaging — join the experience!

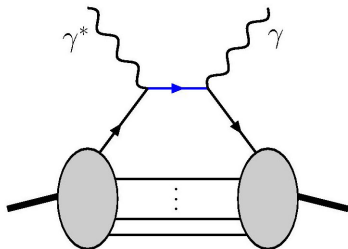


- virtual Compton scattering: $\gamma^* p \rightarrow \gamma p$ (actually: $e^- p \rightarrow e^- \gamma p$)
 - ‘deeply’: $-q_\gamma^2 \gg M_p^2, |t| \rightarrow$ Compton amplitude dominated by (coherent superposition of) Compton scattering off single quarks
- \hookrightarrow only difference between form factor (a) and DVCS amplitude (b) is replacement of photon vertex by two photon vertices connected by **quark** (energy denominator depends on quark momentum fraction x)
- \hookrightarrow DVCS amplitude provides access to momentum-decomposition of form factor = **Generalized Parton Distribution (GPDs)**.

$$\int dx H_q(x, \xi, t) = F_1^q(t)$$



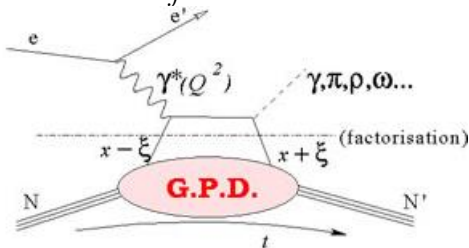
$$\int dx E_q(x, \xi, t) = F_2^q(t)$$



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$$\int dx H_q(x, \xi, t) = F_1^q(t)$$

$$\int dx E_q(x, \xi, t) = F_2^q(t)$$



future experiments

JLab@12GeV, COMPASS II, EIC,
PANDA/FAIR

- form factors: $\overleftarrow{FT} \rho(\vec{r})$
- $GPDs(x, \vec{\Delta})$: form factor for quarks with momentum fraction x
- ↪ suitable FT of $GPDs$ should provide spatial distribution of quarks with momentum fraction x
- careful: cannot measure longitudinal momentum (x) and longitudinal position simultaneously (Heisenberg)
- ↪ consider purely transverse momentum transfer

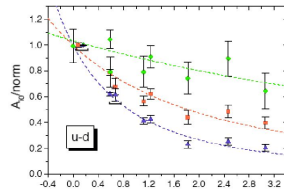
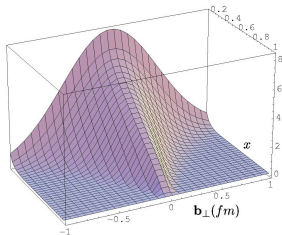
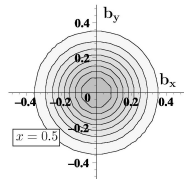
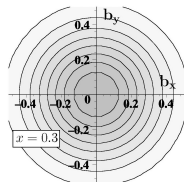
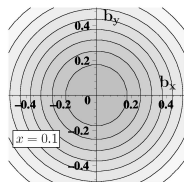
Impact Parameter Dependent Quark Distributions

$$q(x, \mathbf{b}_\perp) = \int \frac{d^2 \Delta_\perp}{(2\pi)^2} H(x, \xi = 0, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp}$$

$q(x, \mathbf{b}_\perp)$ = parton distribution as a function of the separation \mathbf{b}_\perp from the transverse center of momentum $\mathbf{R}_\perp \equiv \sum_{i \in q, g} \mathbf{r}_{\perp, i} x_i$
 MB, Phys. Rev. D62, 071503 (2000)

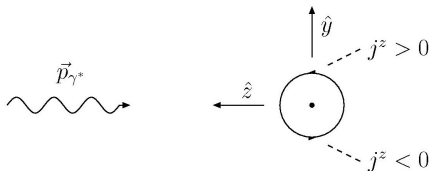
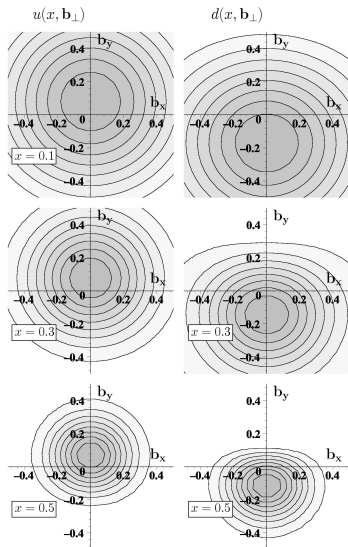
- **No relativistic corrections** (Galilean subgroup!)
- ↪ corollary: interpretation of 2d-FT of $F_1(Q^2)$ as charge density in transverse plane also free of relativistic corrections
- **probabilistic interpretation**

$q(x, \mathbf{b}_\perp)$ for unpol. p



unpolarized proton

- $q(x, \mathbf{b}_\perp) = \int \frac{d^2 \Delta_\perp}{(2\pi)^2} H(x, 0, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp}$
 - x = momentum fraction of the quark
 - \vec{b}_\perp = \perp distance of quark from \perp center of momentum
 - small x : large 'meson cloud'
 - larger x : compact 'valence core'
 - $x \rightarrow 1$: active quark becomes center of momentum
- $\hookrightarrow \vec{b}_\perp \rightarrow 0$ (narrow distribution) for $x \rightarrow 1$

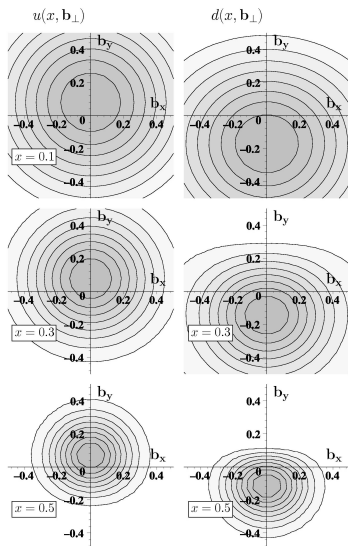


proton polarized in $+\hat{x}$ direction

no axial symmetry!

$$q(x, \mathbf{b}_\perp) = \int \frac{d^2 \Delta_\perp}{(2\pi)^2} H_q(x, 0, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp} \\ - \frac{1}{2M} \frac{\partial}{\partial b_y} \int \frac{d^2 \Delta_\perp}{(2\pi)^2} E_q(x, 0, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp}$$

Physics: relevant density in leading twist DIS is $j^+ \equiv j^0 + j^3$ and left-right asymmetry from j^3



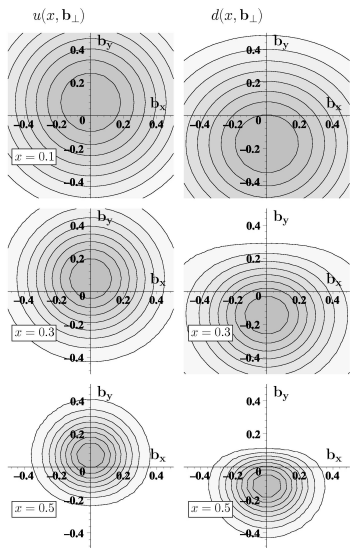
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sign & magnitude of the average shift

model-independently related to p/n
anomalous magnetic moments:

$$\langle b_y^q \rangle \equiv \int dx \int d^2 b_\perp q(x, \mathbf{b}_\perp) b_y \\ = \frac{1}{2M} \int dx E_q(x, 0, 0) = \frac{\kappa_q}{2M}$$



sign & magnitude of the average shift
 model-independently related to p/n
 anomalous magnetic moments:

$$\begin{aligned} \langle b_y^q \rangle &\equiv \int dx \int d^2 b_\perp q(x, \mathbf{b}_\perp) b_y \\ &= \frac{1}{2M} \int dx E_q(x, 0, 0) = \frac{\kappa_q}{2M} \end{aligned}$$

$$\kappa^P = 1.913 = \frac{2}{3} \kappa_u^P - \frac{1}{3} \kappa_d^P + \dots$$

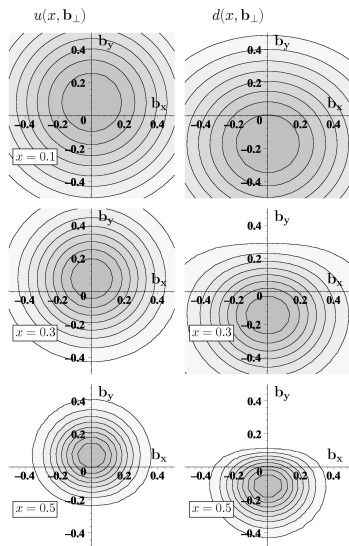
- u -quarks: $\kappa_u^P = 2\kappa_p + \kappa_n = 1.673$

↪ shift in $+\hat{y}$ direction

- d -quarks: $\kappa_d^P = 2\kappa_n + \kappa_p = -2.033$

↪ shift in $-\hat{y}$ direction

- $\langle b_y^q \rangle = \mathcal{O}(\pm 0.2 \text{ fm})$!!!!

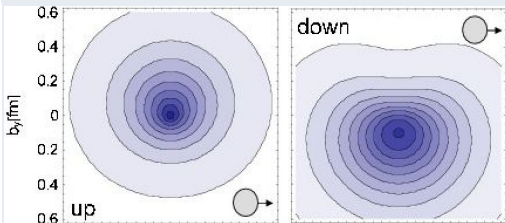


sign & magnitude of the average shift

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$$\begin{aligned} \langle b_y^q \rangle &\equiv \int dx \int d^2 b_\perp q(x, \mathbf{b}_\perp) b_y \\ &= \frac{1}{2M} \int dx E_q(x, 0, 0) = \frac{\kappa_q}{2M} \end{aligned}$$

lattice QCD (QCDSF): lowest moment



transverse images \leftrightarrow Ji relation for quark angular momentum:

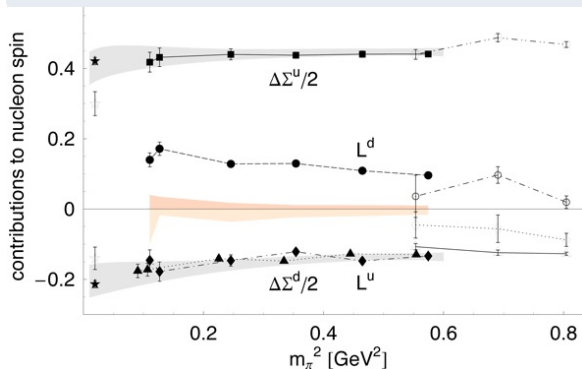
- $J_q^x = m_N \int dx x r^y q(x, \mathbf{r}_\perp)$ with $b^y = r^y - \frac{1}{2m_N}$, where $q(x, \mathbf{r}_\perp)$ is distribution relative to CoM of whole nucleon
- recall: $q(x, \mathbf{b}_\perp)$ for nucleon polarized in $+\hat{x}$ direction

$$q(x, \mathbf{b}_\perp) = \int \frac{d^2 \Delta_\perp}{(2\pi)^2} H_q(x, 0, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp} \\ - \frac{1}{2M_N} \frac{\partial}{\partial b_y} \int \frac{d^2 \Delta_\perp}{(2\pi)^2} E_q(x, 0, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp}$$

$$\Rightarrow J_q^x = M_N \int dx x r^y q(x, \mathbf{r}_\perp) = \int dx x \left(m_N b^y + \frac{1}{2} \right) q(x, \mathbf{r}_\perp) \\ = \frac{1}{2} \int dx x [H(x, 0, 0) + E(x, 0, 0)]$$

- X.Ji(1996): rotational invariance \Rightarrow apply to all components of \vec{J}_q
- partonic interpretation exists only for \perp components!

lattice: QCDSF



- no disconnected quark loops
- chiral extrapolation

$$J^q = \frac{1}{2} \int dx x [H(x, 0, 0) + E(x, 0, 0)]$$

$$L^q = J^q - \frac{1}{2} \Delta\Sigma^q$$

TMDs

- Transverse Momentum Dependent Parton Distributions
- 8 structures possible at leading twist (only 3 for PDFs)
- f_{1T}^\perp and h_1^\perp require both **orbital angular momentum** and **final state interaction**
- can be measured in semi-inclusive deep-inelastic scattering (SIDIS) & Drell-Yan (DY) $q\bar{q} \rightarrow \mu^+\mu^-$

experiments

JLab@6GeV & 12GeV, HERMES, COMPASS I & II, RHIC, EIC, FAIR/PANDA

“TMDs”

nucleon polarisation

Sivers function

correlation between the transverse spin of the nucleon and the transverse momentum of the quark
















sensitive to orbital angular momentum

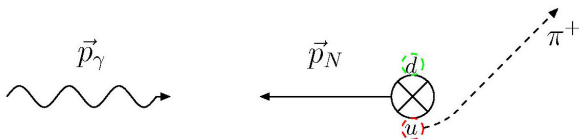
Boer-Mulders function

correlation between the transverse spin and the transverse momentum of the quark in unpol nucleons

T-odd

quark polarisation

	U	L	T
U	f_1  number density q		f_{1T}^\perp  -  Sivers
L		g_1  -  helicity Δq	g_{1T}  - 
T	h_1^\perp  -  Boer Mulders	h_{1L}^\perp  - 	h_1  -  transversity h_{1T}^\perp  - 

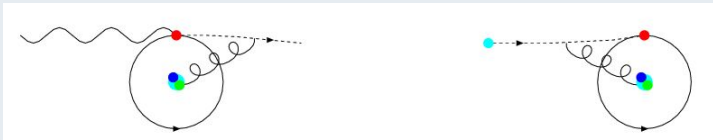
Sivers f_{1T}^\perp in semi-inclusive deep-inelastic scattering (SIDIS) $\gamma p \rightarrow \pi X$ 

- u, d distributions in \perp polarized proton have left-right asymmetry in \perp position space (T-even!); sign 'determined' by κ_u & κ_d
 - attractive FSI deflects active quark towards the CoM
- \hookrightarrow FSI translates position space distortion (before the quark is knocked out) in $+\hat{y}$ -direction into momentum asymmetry that favors $-\hat{y}$ direction \rightarrow 'chromodynamic lensing'

\Rightarrow $\kappa_p, \kappa_n \longleftrightarrow$ sign of SSA!!!!!!! MB, PRD 69, 074032 (2004)

- confirmed by HERMES (and recent COMPASS) p data; consistent with vanishing isoscalar Sivers (COMPASS) \rightarrow G. Schnell

compare FSI for 'red' q that is being knocked out of nucleon with ISI for 'anti-red' \bar{q} that is about to annihilate with a 'red' target q



FSI in SIDIS

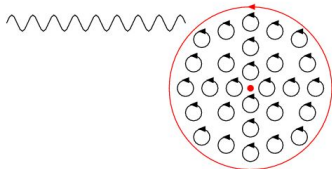
- knocked-out q 'red'
- ↪ spectators 'anti-red'
- ↪ interaction between knocked-out quark and spectators **attractive**

ISI in DY

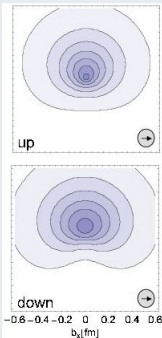
- incoming \bar{q} 'anti-red'
- ↪ struck target q 'red'
- ↪ spectators also 'anti-red'
- ↪ interaction between incoming \bar{q} and spectators **repulsive**

test of $f_{1T}^\perp(x, \mathbf{k}_\perp)_{DY} = -f_{1T}^\perp(x, \mathbf{k}_\perp)_{SIDIS}$ and $h_1^\perp(x, \mathbf{k}_\perp)_{DY} = -h_1^\perp(x, \mathbf{k}_\perp)_{SIDIS}$
critical test of TMD factorization approach

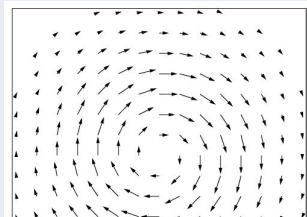
q with polarization \odot



lattice calculation (QCDSF)



unpolarized target



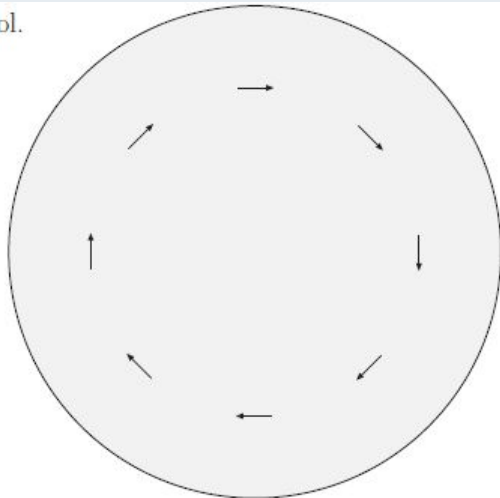
- transversity distribution in unpol. target described by chirally odd GPD \bar{E}_T (M.Diehl & P.Haegler '05)
 - $\bar{E}_T > 0$ for u & d (QCDSF)
 - connection $h_1^\perp(x, \mathbf{k}_\perp) \leftrightarrow \bar{E}_T$ similar to $f_{1T}^\perp(x, \mathbf{k}_\perp) \leftrightarrow E$ (MB '05)
- $\hookrightarrow h_1^\perp(x, \mathbf{k}_\perp) < 0$ for $u/p, d/p, u/\pi, \bar{d}/\pi$
- $h_{1SIDIS}^\perp = -h_{1DY}^\perp$

experiments (no polarization needed!):

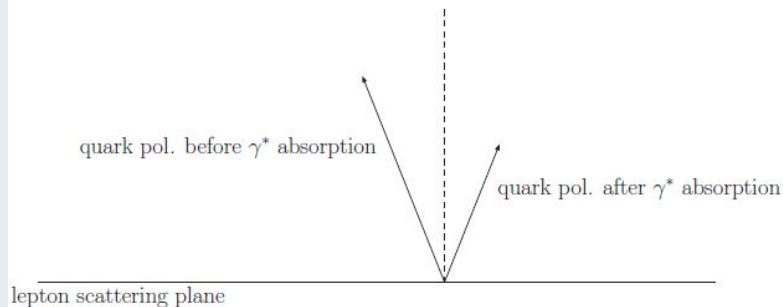
HERMES, COMPASS, RHIC, JLab@12GeV, FAIR/PANDA, EIC

Primordial Quark Transversity Distribution

→ \perp quark pol.



Flip of Quark Transverse Spin Component

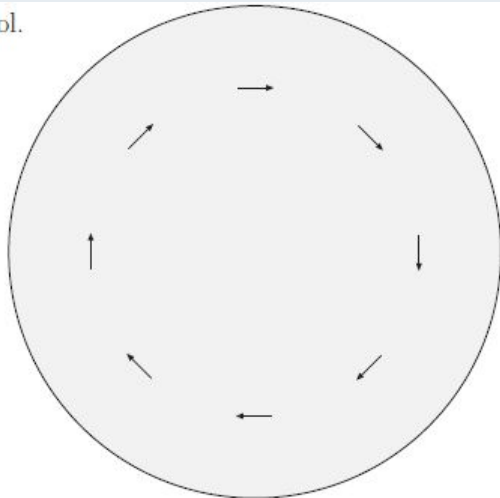


when \perp pol. quark absorbs γ^* , \perp polarization

- gets reduced in size
- tilted symmetrically w.r.t. normal of lepton scattering plane

Primordial Quark Transversity Distribution

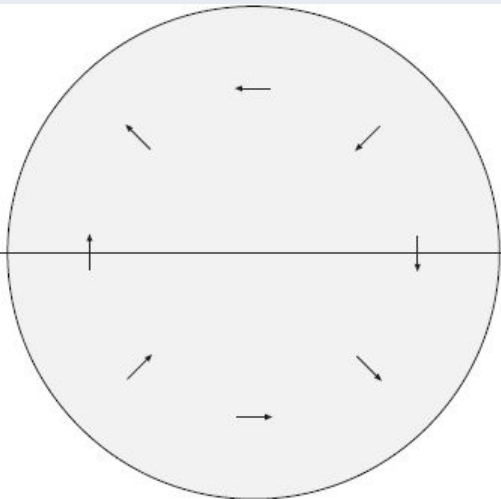
→ \perp quark pol.



Quark Transversity Distribution after γ^* Absorption

→ \perp quark pol.

lepton scattering plane

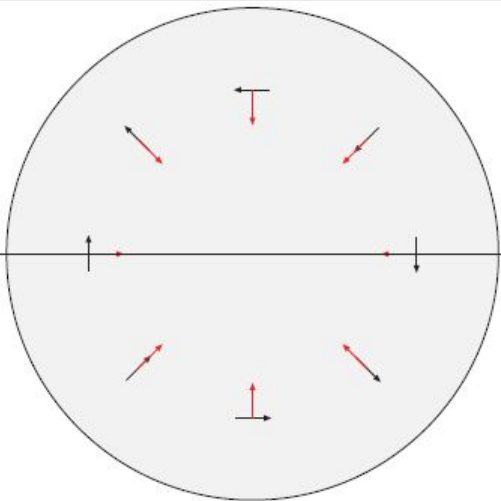


\perp momentum (of q) due to FSI

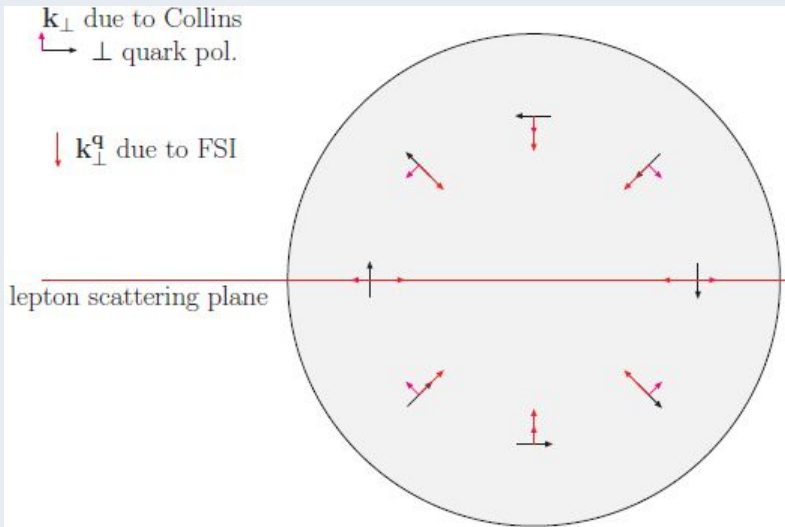
\rightarrow \perp quark pol.

\downarrow k_{\perp}^q due to FSI

lepton scattering plane



additional \perp momentum (of π) due to Collins effect



Collins: favored π momentum preferentially to **left** (quark spin up)

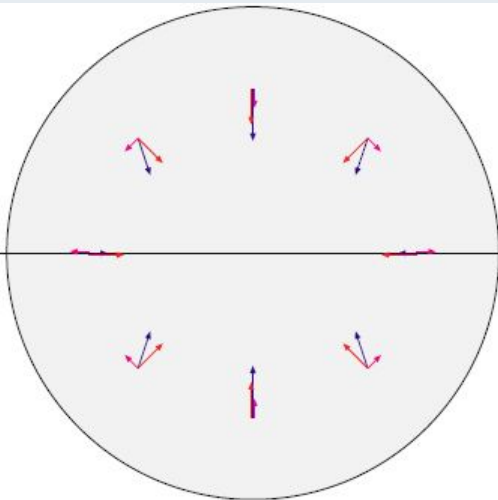
net k_{\perp}^{π} (FSI + Collins)

\downarrow
 k_{\perp} due to Collins

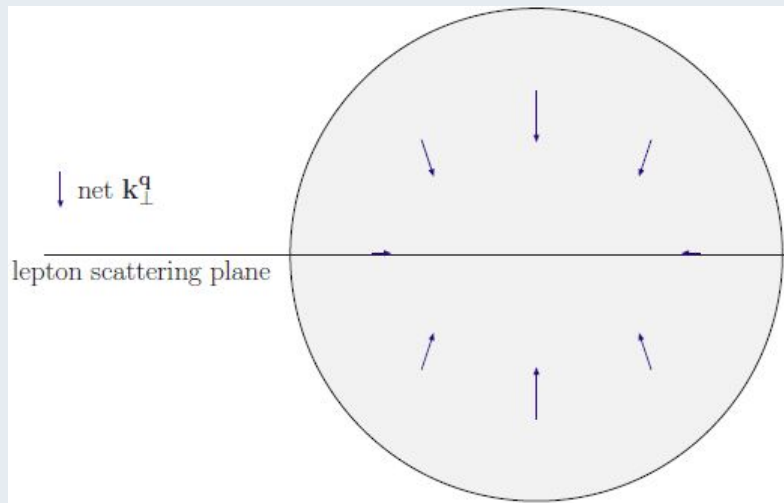
\downarrow k_{\perp}^q due to FSI

\downarrow net k_{\perp}^q

lepton scattering plane



net k_{\perp}^{π} (FSI + Collins)



$\cos 2\pi$ modulation of k_{\perp}^{π}

higher twist in polarized DIS

- $\sigma_{LL} \propto g_1 - \frac{2Mx}{\nu} g_2$
 - $g_1 = \frac{1}{2} \sum_q e_q^2 g_1^q$ with $g_1^q = q^\uparrow(x) + \bar{q}^\uparrow(x) - q^\downarrow(x) - \bar{q}^\downarrow(x)$
 - g_2 involves quark-gluon correlations
- ↪ no parton interpret. as difference between number densities for g_2
- for \perp pol. target, g_1 & g_2 contribute equally

$$\sigma_{LT} \propto g_T \equiv g_1 + g_2$$

- ↪ 'clean' separation between g_2 and $\frac{1}{Q^2}$ corrections to g_1

What can we learn from g_2 ?

- $g_2 = g_2^{WW} + \bar{g}_2$ with $g_2^{WW}(x) \equiv -g_1(x) + \int_x^1 \frac{dy}{y} g_1(y)$

$$d_2 \equiv 3 \int dx x^2 \bar{g}_2(x) = \frac{1}{2MP^{+2}S^x} \langle P, S | \bar{q}(0) g G^{+y}(0) \gamma^+ q(0) | P, S \rangle$$

$$d_2 \equiv 3 \int dx x^2 \bar{g}_2(x) = \frac{1}{2MP^{+2}S^x} \langle P, S | \bar{q}(0) g G^{+y}(0) \gamma^+ q(0) | P, S \rangle$$

$$\sqrt{2}G^{+y} = G^{0y} + G^{zy} = -E^y + B^x$$

$$d_2 \equiv 3 \int dx x^2 \bar{g}_2(x) = \frac{1}{2MP^{+2}S^x} \langle P, S | \bar{q}(0) g G^{+y}(0) \gamma^+ q(0) | P, S \rangle$$

color Lorentz force

MB 2008

$$\sqrt{2}G^{+y} = G^{0y} + G^{zy} = -E^y + B^x = -\left(\vec{E} + \vec{v} \times \vec{B}\right)^y \text{ for } \vec{v} = (0, 0, -1)$$

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\leftrightarrow $d_2 \leftrightarrow$ average **color Lorentz force** acting on quark moving with $v = c$ in $-\hat{z}$ direction in the instant after being struck by γ^*

$$\langle F^y \rangle = -2M^2 d_2 = -\frac{M}{P^{+2}S^x} \langle P, S | \bar{q}(0) g G^{+y}(0) \gamma^+ q(0) | P, S \rangle$$

cf. Qiu-Sterman matrix element $\langle k_\perp^y \rangle \equiv \int_0^1 dx \int d^2 k_\perp k_\perp^2 f_{1T}^\perp(x, k_\perp^2)$

$$\langle k_\perp^y \rangle = -\frac{1}{2p^+} \left\langle P, S \left| \bar{q}(0) \int_0^\infty dx^- g G^{+y}(x^-) \gamma^+ q(0) \right| P, S \right\rangle$$

semi-classical interpretation: average k_\perp in SIDIS obtained by correlating the quark density with the transverse impulse acquired from (color) Lorentz force acting on struck quark along its trajectory to (light-cone) infinity

matrix element defining d_2

\leftrightarrow

1st integration point in QS-integral

$$d_2 \equiv 3 \int dx x^2 \bar{g}_2(x) = \frac{1}{2MP^{+2}S_x} \langle P, S | \bar{q}(0) g G^{+y}(0) \gamma^+ q(0) | P, S \rangle$$

color Lorentz force

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sign of $d_2 \leftrightarrow \perp$ imaging

- $\kappa_q/p \rightarrow$ sign of deformation
- \hookrightarrow direction of average force
- $\hookrightarrow d_2^u > 0, d_2^d < 0$
- cf. $f_{1T}^{\perp u} < 0, f_{1T}^{\perp d} < 0$

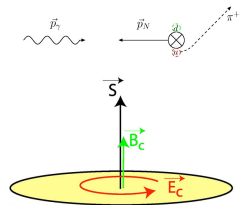
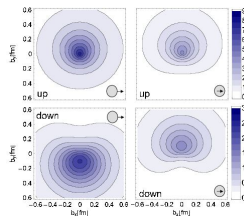
lattice (Göckeler et al., 2005)

$$d_2^u \approx 0.010, d_2^d \approx -0.0056$$

magnitude of d_2

- $\langle F^y \rangle = -2M^2 d_2 = -10 \frac{\text{GeV}}{fm} d_2$
- expect partial cancellation of forces in SSA
- $\hookrightarrow |\langle F^y \rangle| \ll \sigma \approx 1 \frac{\text{GeV}}{fm}$
- $\hookrightarrow d_2 = \mathcal{O}(0.01)$

- Deeply Virtual Compton Scattering \rightarrow GPDs
- \hookrightarrow impact parameter dependent PDFs $q(x, \mathbf{b}_\perp)$
- $E^q(x, 0, -\Delta_\perp^2) \leftrightarrow \kappa_{q/p}$ (contribution from quark flavor q to anomalous magnetic moment)
- $E^q(x, 0, -\Delta_\perp^2) \rightarrow \perp$ deformation of PDFs for \perp polarized target
- \perp deformation \leftrightarrow (sign of) SSA (Sivers; Boer-Mulders)
- parton interpretation for Ji-relation
- higher-twist $(\int dx x^2 \bar{g}_2(x), \int dx x^2 \bar{e}(x)) \leftrightarrow \perp$ force in DIS
- \perp deformation \leftrightarrow (sign of) quark-gluon correlations $(\int dx x^2 \bar{g}_2(x), \int dx x^2 \bar{e}(x))$



combine complementary information from deeply-virtual Compton scattering, semi-inclusive DIS & Drell-Yan to study orbital angular momentum and map the 3-d structure of hadrons

Q^2 scaling for Compton form factor (JLab)