



Optical/IR Observational Astronomy

Spectroscopy



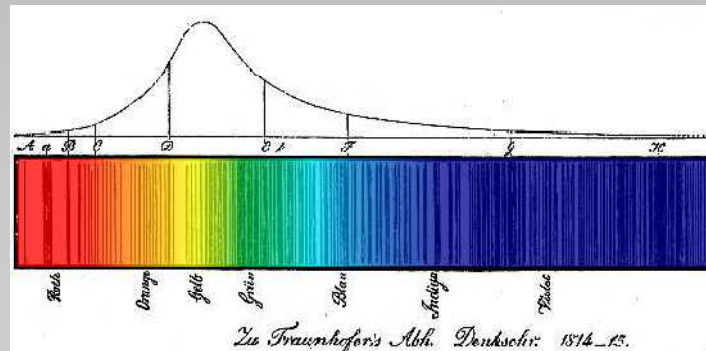
David Buckley, SALT

Background

Spectroscopy is really just monochromatic photometry

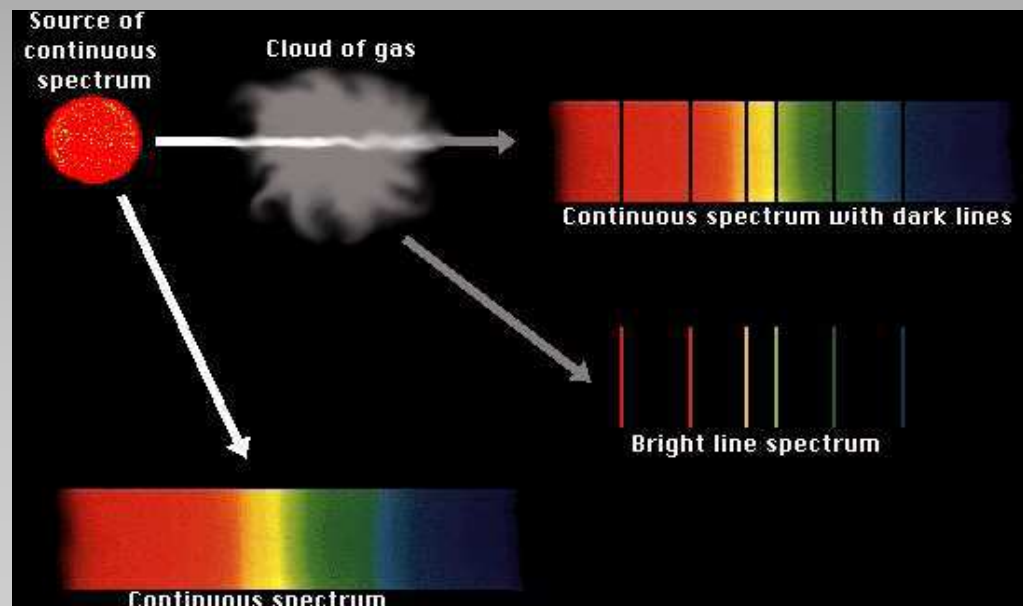
History

- 1637 Descartes explained the origin of the rainbow.
- 1666 Newton's classic experiments on the nature of colour.
- 1752 Melvil discovered the yellow "D lines" emitted by sodium vapour.
- 1802 Wollaston discovered dark lines in the spectrum of the Sun.
- 1814 Fraunhofer used a small theodolite telescope to examine stellar spectra and found similar lines as in the solar spectrum.
 - He catalogued the features as A, B, C, ... etc,
 - some which persist to the present day – the sodium "D" lines, the "G" band (CN - cyanogen).



Background

- 1859 Kirchhoff & Bunsen experiments established that a heated surface emits a continuous (Planck or “black body”) spectrum
 - a heated low pressure gas emits an emission spectrum with discrete lines at wavelengths characteristic of the gas; and a cool low pressure
 - » **Examples include planetary nebula & HII (ionized hydrogen) region**
 - gas in front of a hot source absorbs at those same characteristic wavelengths.
 - » **Examples include the photospheres of stars (black bodies with absorbing “atmospheres”**



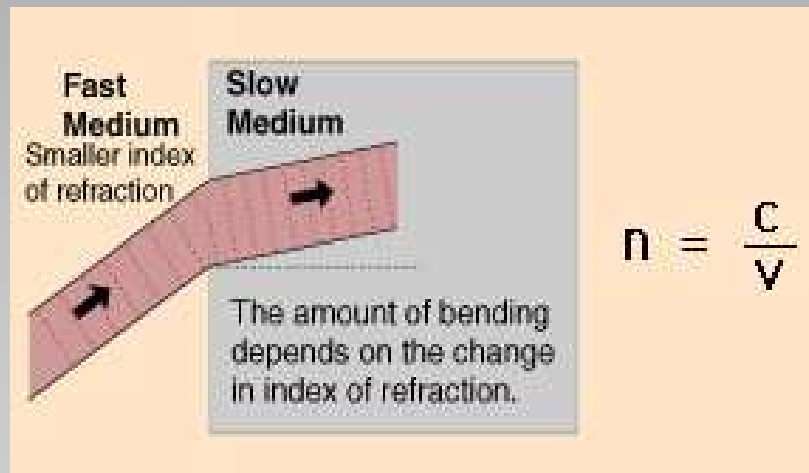


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Spectroscopy

How to disperse light?

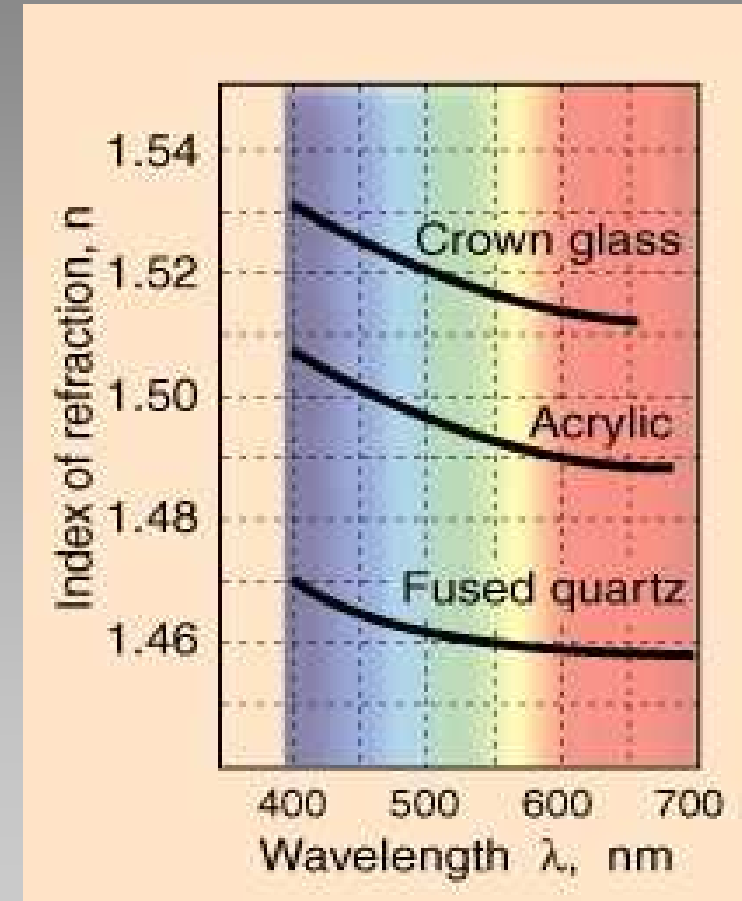
- First achieved with glass prisms
 - Highly transparent
- Made from material with higher index of refraction than air
 - Bends the lights
- Made from material with wavelength dependent index of refraction
 - *Disperses* light
 - Degree of bending is wavelength dependent



Index of refraction, n , is ratio of speed of light in vacuum compared to speed of light in the material

Index of Refraction

- Non-linear with wavelength
 - Degree of dispersion also non-linear
- Original spectrographs used prisms
 - Continued well into 20th C
- But, problems with prisms....
 - They cannot achieve very high dispersion, therefore can't do high resolution ($R = \lambda/\Delta\lambda$) spectroscopy
 - They're quite "lossy":
 - » reflection losses at air-glass surfaces (4% per surface for uncoated prisms)
 - » Absorption losses in the glass, particularly in the UV
 - Temperature changes can cause index changes
 - ⇒ Dispersion changes and line shifts



Dispersion Formula for Prisms

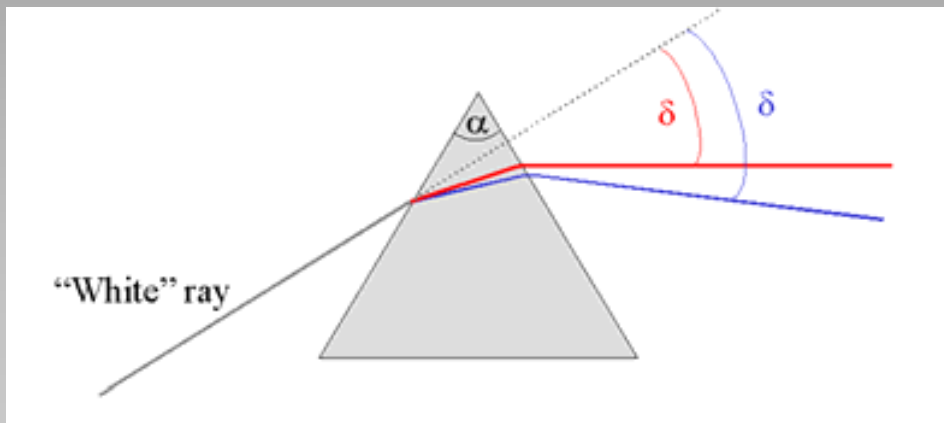
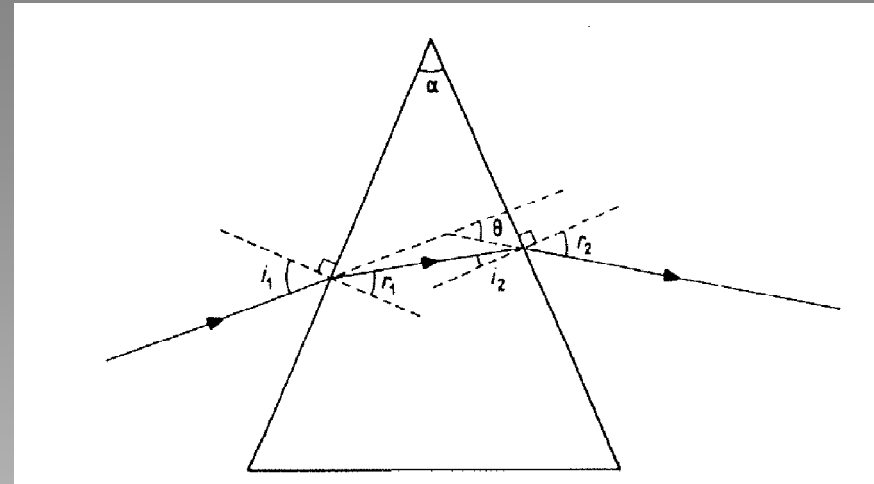
- Just apply Snell's law

$$n_1 \sin i = n_2 \sin r$$

– From air-to-glass ($n_1 = 1, n_2 = 1.5$)

$$\Rightarrow n_2 = \sin i / \sin r$$

- Geometry for generalized case ➔



- Deviation is:

$$\delta = i + r - \alpha$$

- Dispersion is simply:

$$\Delta\theta/\Delta\lambda = d\delta/d\lambda$$

More on Prisms

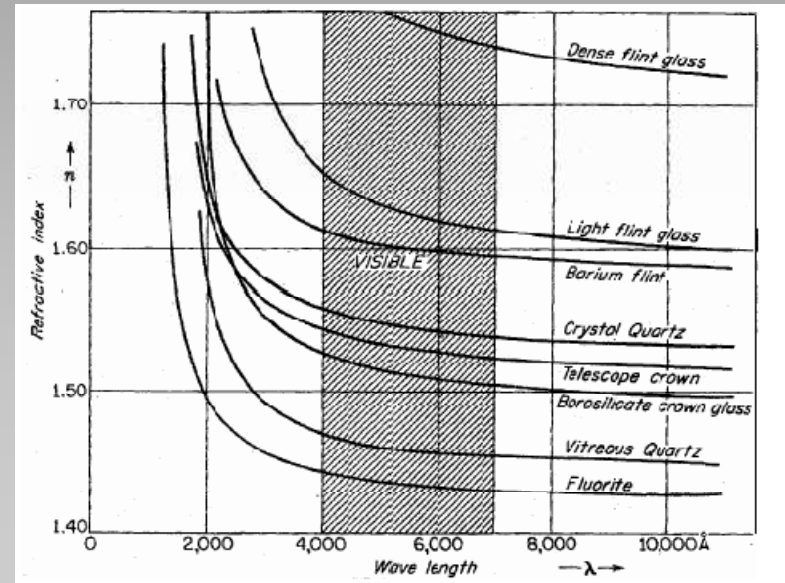
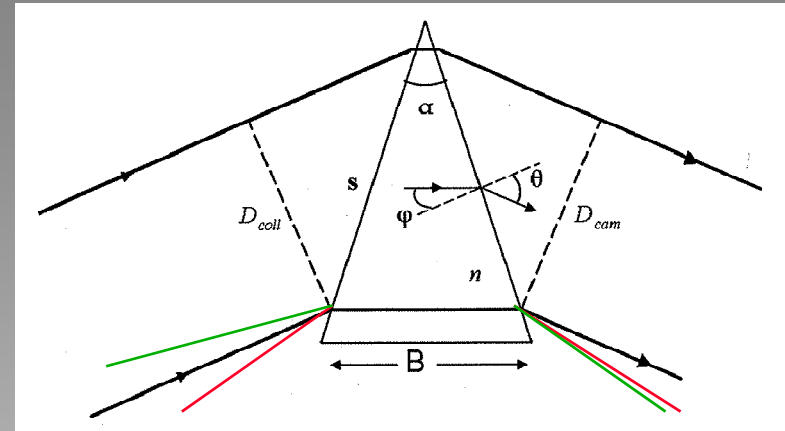
- Prisms
 - Used in *minimum deviation condition*
 - Symmetrical
 - » Beam width constant
 - ⇒ No astigmatism
 - From trig:

$$\frac{d\theta}{d\lambda} = \frac{d\theta}{dn} \frac{dn}{d\lambda} = \frac{B}{D_{\text{cam}}} \frac{dn}{d\lambda} = \frac{2 \sin(\alpha/2)}{\cos \theta} \frac{dn}{d\lambda}$$

- So dispersion depends on
 - » Apex angle of prism
 - » Emergent angle of diffraction
 - » Index of refraction dependency on λ
- Spectral resolution is given as:

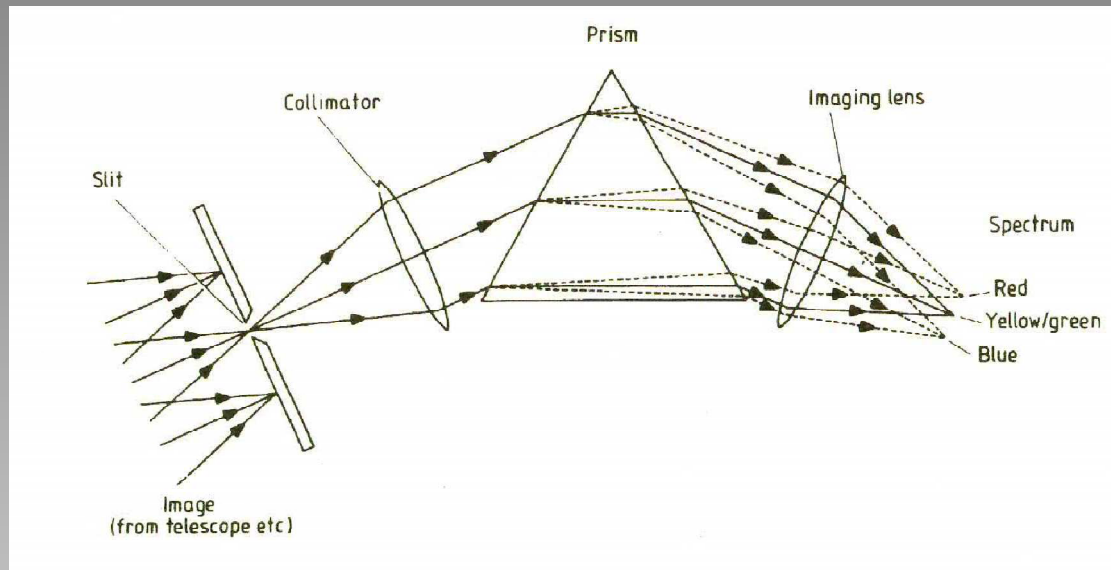
$$R = \lambda / \Delta\lambda = B \, dn/d\lambda \quad (\text{slitless})$$

(higher resolution in blue compared to red)

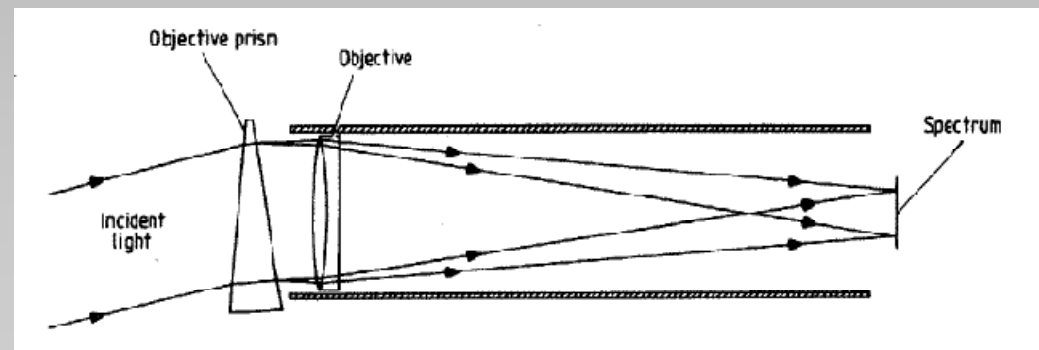


Prism Spectrographs

- This is basis of all prism spectrographs



- First used as objective prisms
 - Prisms placed at telescope aperture



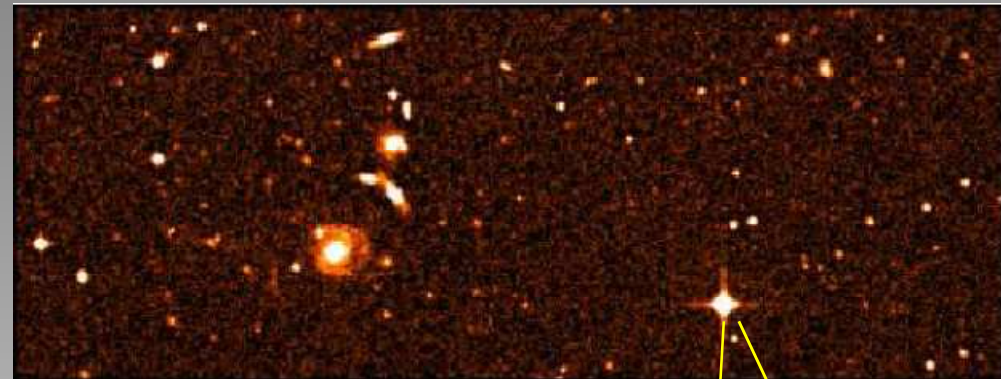


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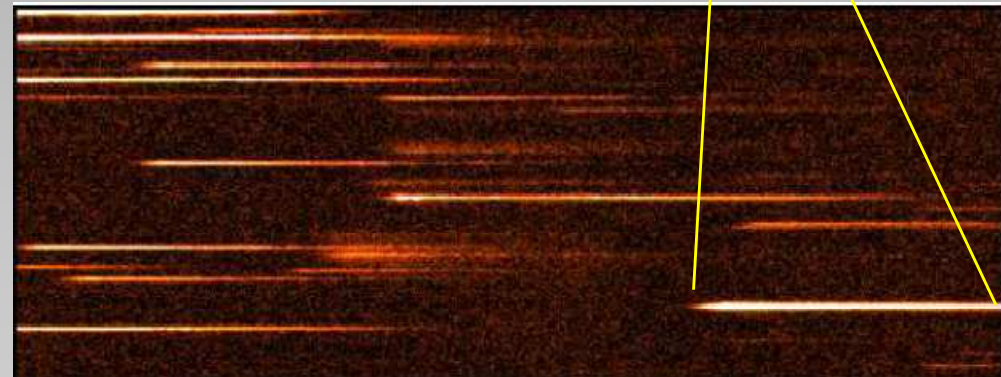
Spectroscopy

Objective Prism Surveys

- Used to disperse images on the sky
- Create multiple low resolution spectra
- Resolution determined by the *prism dispersion, detector resolution and seeing*
- First done with photographic plates
 - e.g. on the Cape Observatory McLean refractor



Image



Spectrogram (dispersed image)

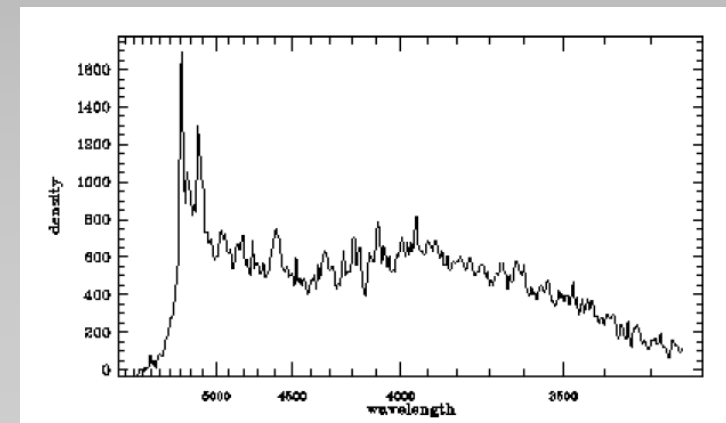
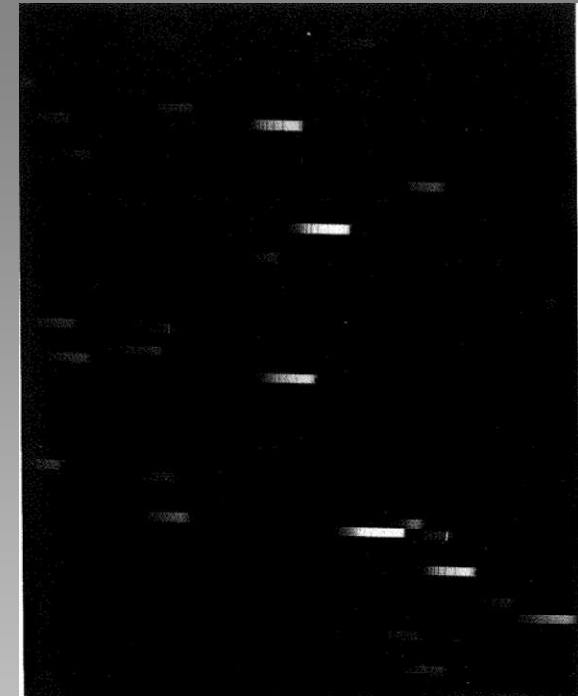


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Spectroscopy

Objective Prisms

- Spectra widened by trailing telescope in direction perpendicular to dispersion
 - Easier to see features by eye (first techniques)
- First catalogue of stellar spectra done at Harvard College Observatory by Annie Cannon
 - Laboriously catalogued 250,000 stars, by eye!



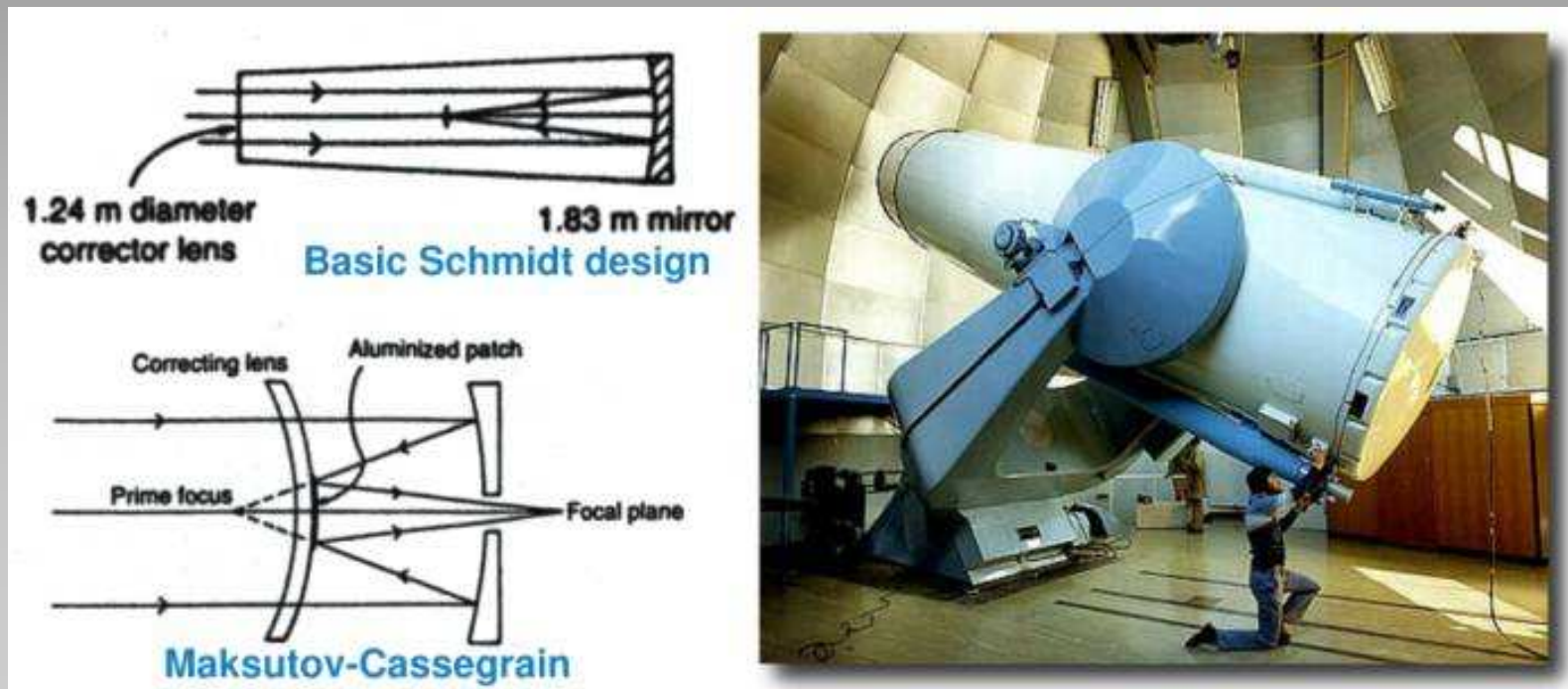


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Spectroscopy

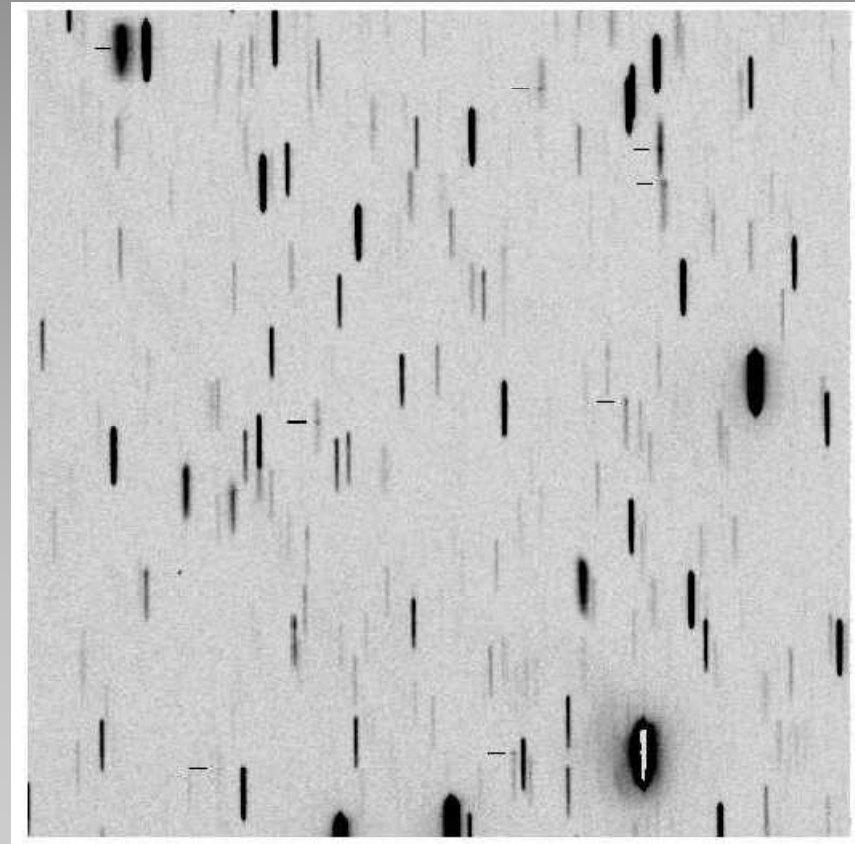
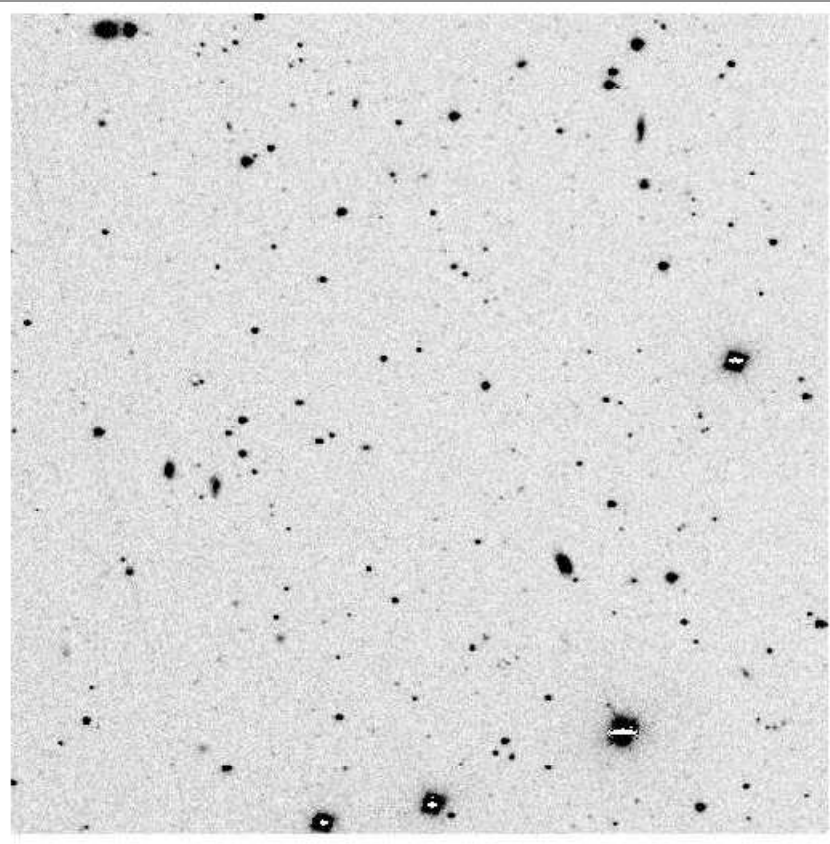
Objective Prisms Surveys

- Combined with Schmidt telescope, objective prisms could be used to spectroscopically survey large areas of sky
- Multiplex advantage: area coverage *and* wavelength information



Example of Objective Prisms Spectra

- Search for emission line objects
 - e.g. star formation in galaxies



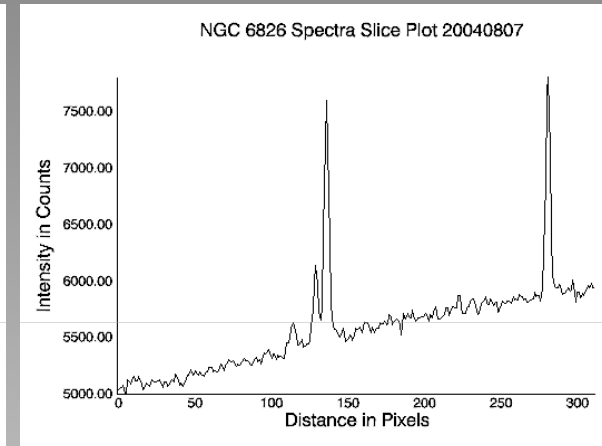
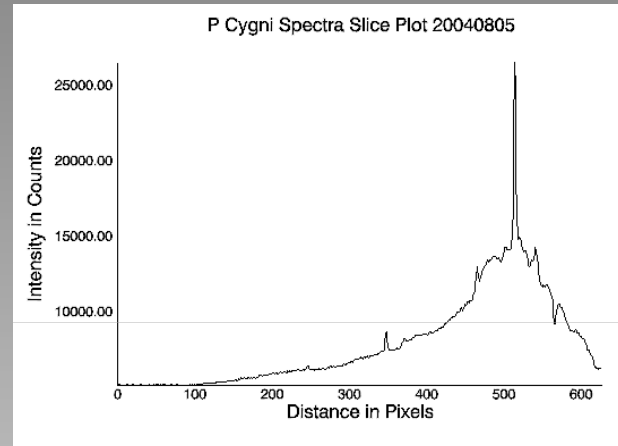
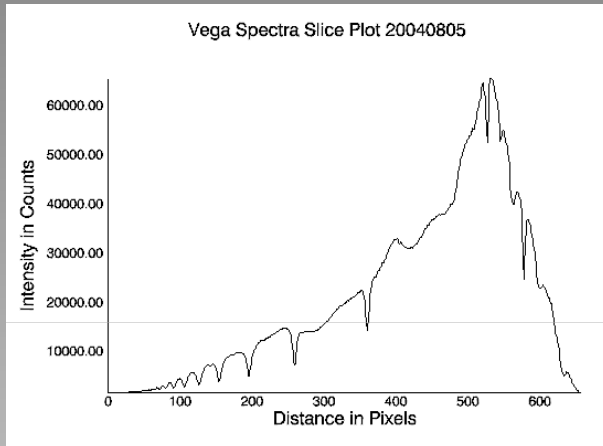


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Spectroscopy

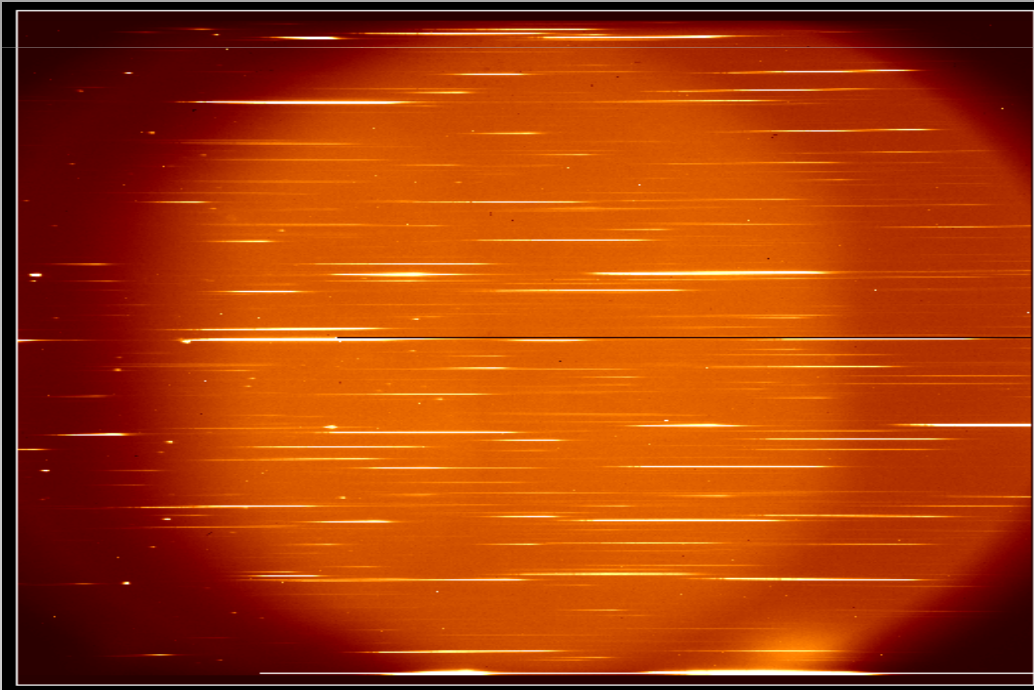
Example of Objective Prisms Spectra

- Spectral classification



Problems with Objective Prisms

- Low spectral resolution
- Spectra are smeared
- Overlapping objects in a dense field
- Non-linear dispersion
- Generally only for relatively bright objects



Diffraction Gratings

- Uses the wave properties of light
- Interference fringes



- Interference from periodic structures
 - Equally spaced grooves
 - Either in reflection or transmission



Optical/IR Observational Astronomy

Spectroscopy

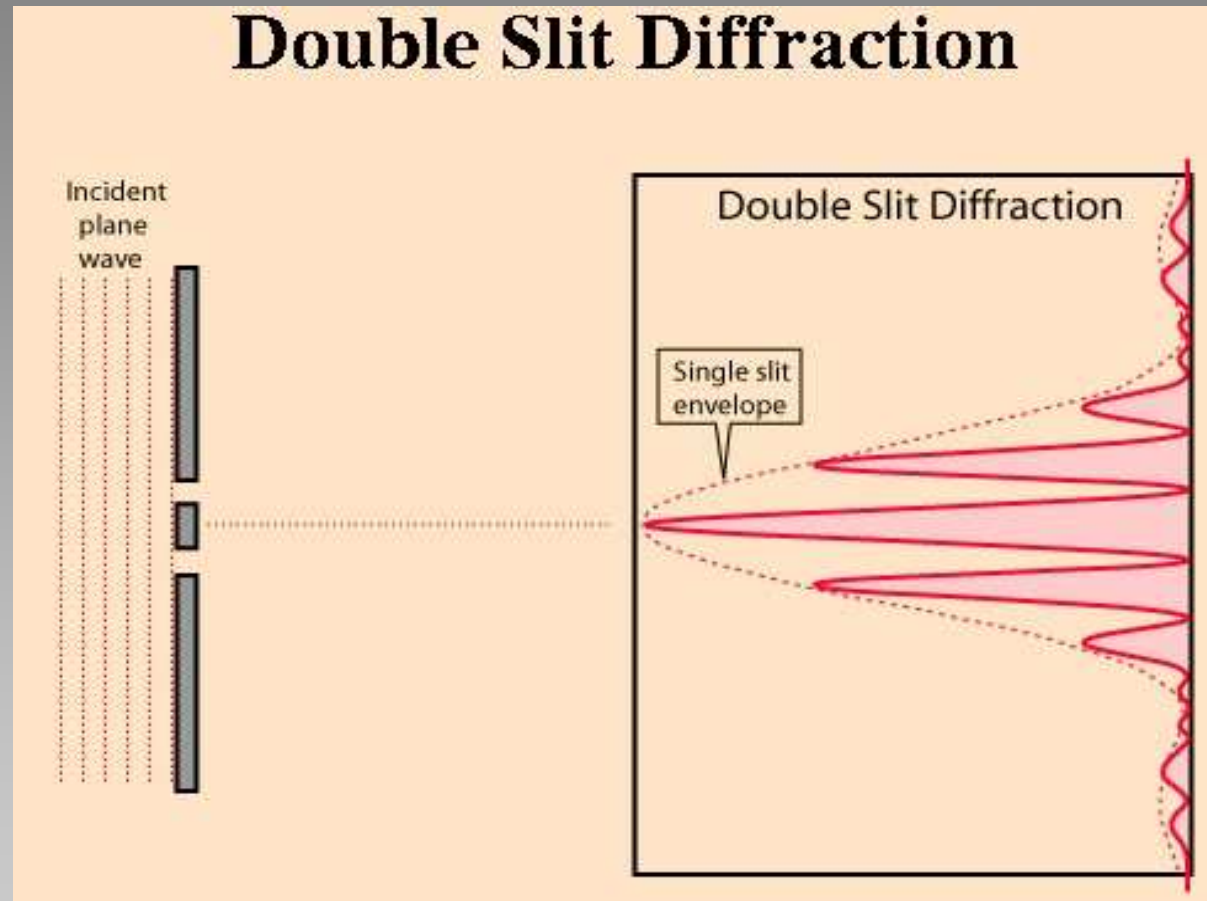
Diffraction Gratings

For constructive interference:

$$m \lambda \simeq d \sin \theta$$

For multiple slits, fringes become narrower

Double Slit Diffraction



Diffraction Gratings: Grating Equation

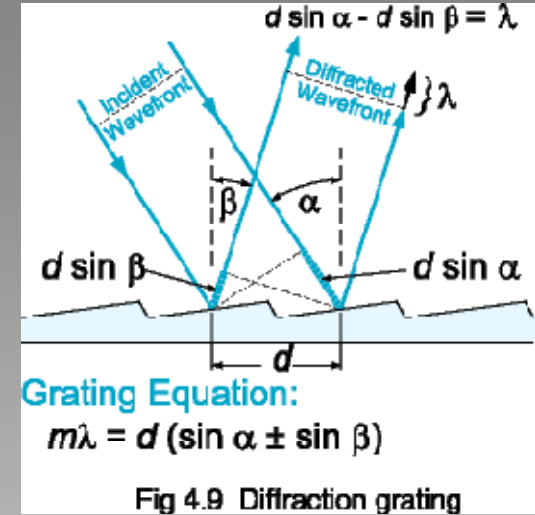
- Constructive interference occurs when the path difference between two diffracted waves differs by an integral number of wavelength

$$m \lambda = d(\sin \alpha \pm \sin \beta)$$

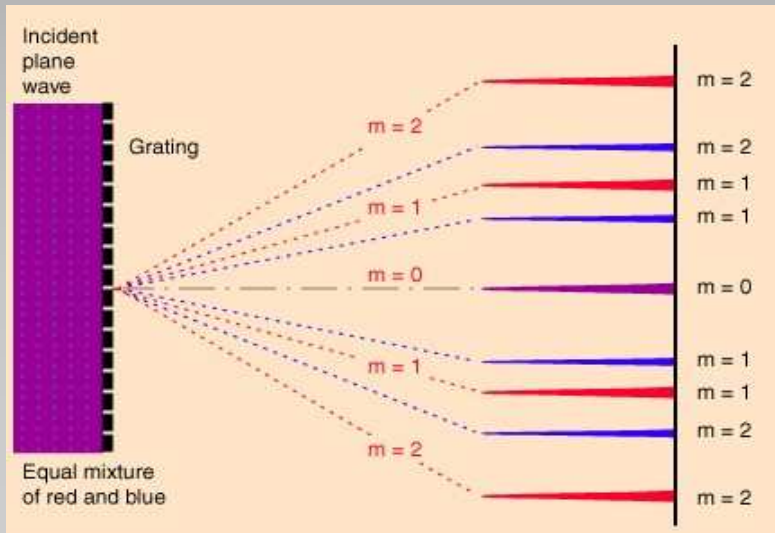
- Dispersion given as

$$\frac{\Delta \beta}{\Delta \lambda} = \frac{m}{d \cos \beta}$$

- Proportional to order, m
- Inversely proportional to spacing, d (proportional to line density, $l = 1/d$)
e.g. 300 l/mm grating is lower resolution than 1200 l/mm



$$m \lambda = d(\sin \alpha - \sin \beta)$$

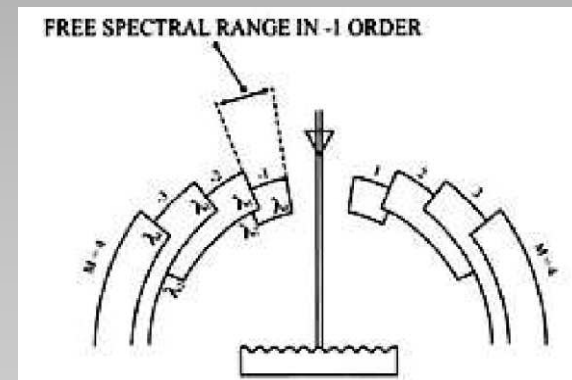


Diffraction Gratings: Free Spectral Range

- For a given diffraction angle different orders can overlap, satisfying the condition:

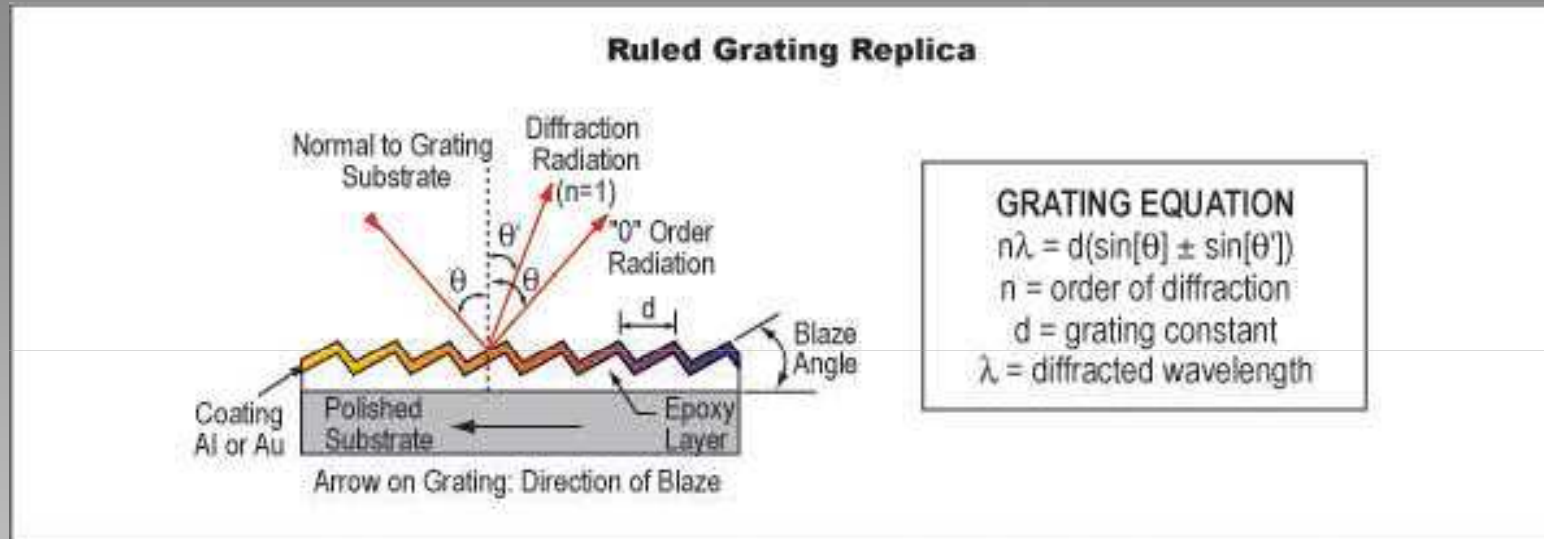
$$m \lambda = 2 m \frac{\lambda}{2} = 3 m \frac{\lambda}{3} \dots$$

- So a first order ($m = 1$) diffracted wavelength at 800 nm overlaps with the second order ($m=2$) diffracted wavelength at 400 nm and third order ($m=3$) diffracted wavelength at 266.7 nm
- Typically we use order separating filters if detectors are sensitive to these higher order diffracted wavelengths
- Free Spectral Range is that part of a spectrum free of overlapping orders
- More of an issues for gratings designed to work at high orders (e.g. echelle gratings)



Grating Blaze

- The tilt angle of the grating facets with respect to grating surface normal



- Most efficient angle, since diffraction at the blaze angle obeys reflection condition of $i = r$
- Diffraction efficiency drops off moving away from the blaze angle/wavelength

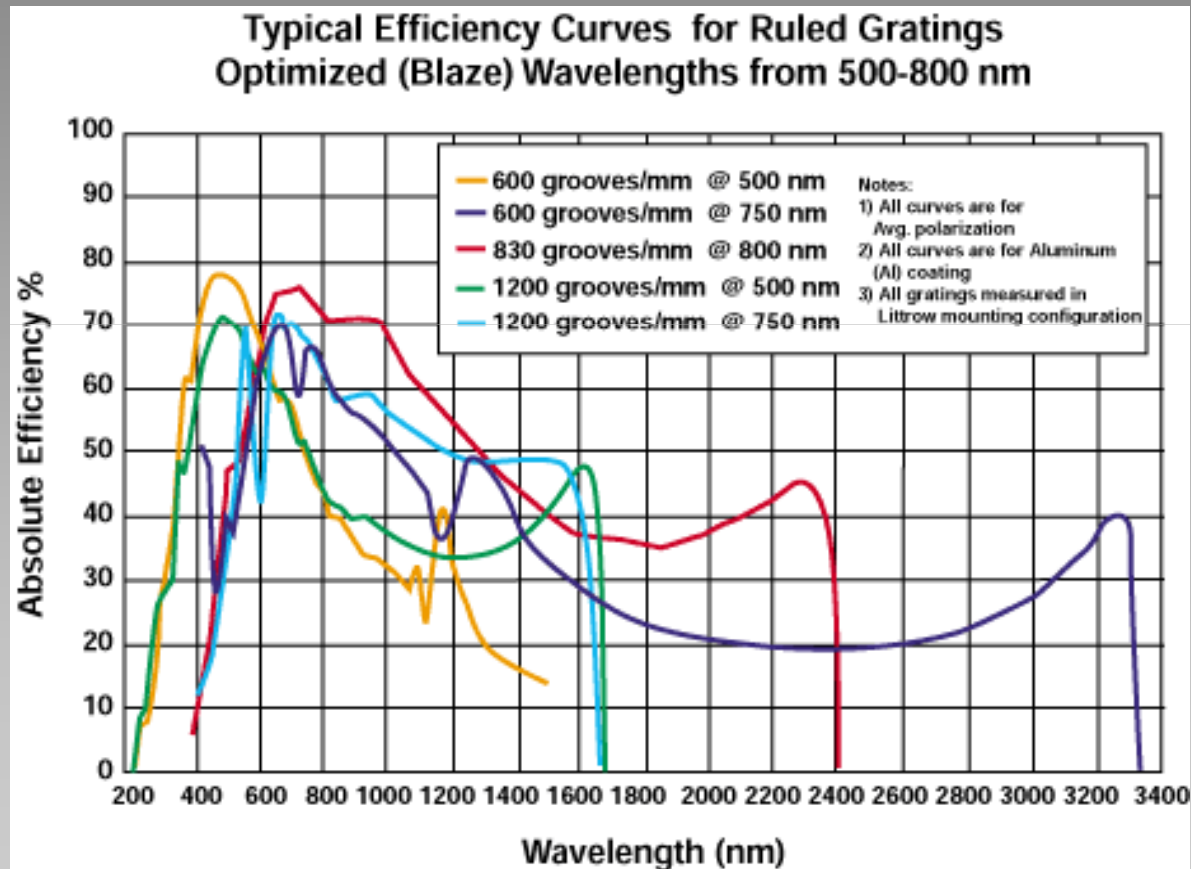


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Spectroscopy

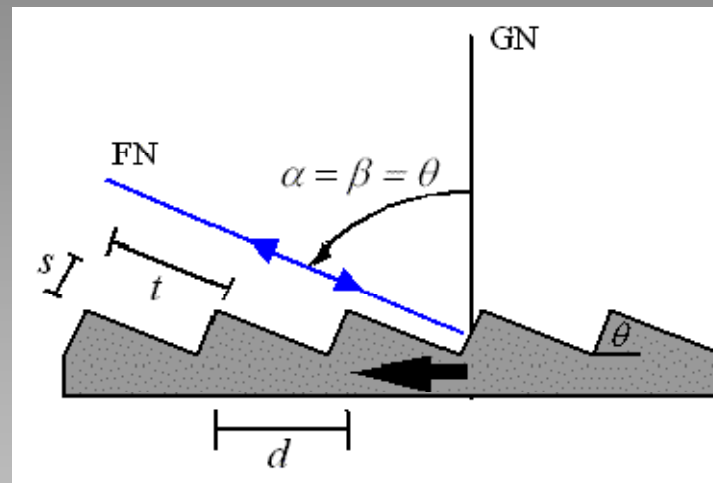
Grating Efficiencies

- Grating efficiency drops off moving away from blaze wavelength:



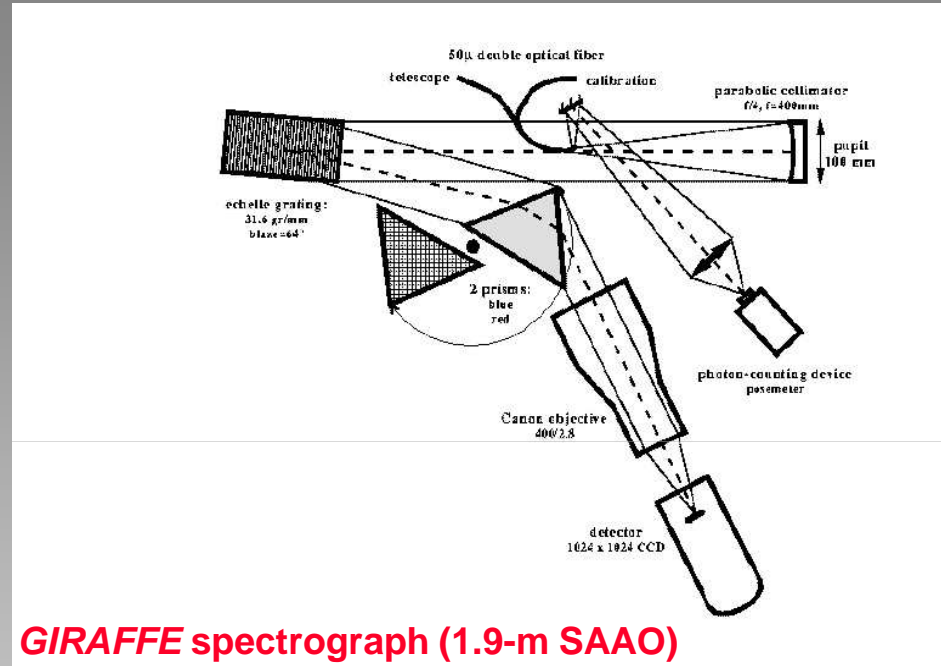
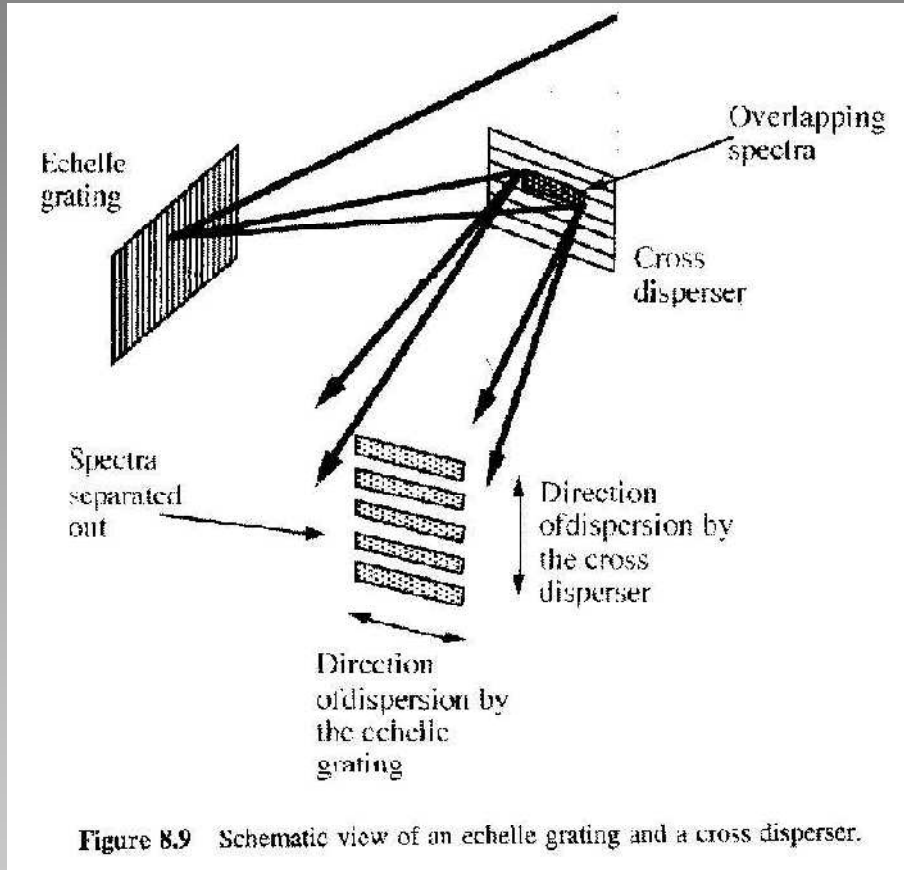
Echelle Gratings

- High resolution gratings achieved by making the blaze condition (highest efficiency) satisfied and high orders (e.g. $m = 50 - 100$)



- High orders overlap, so need to separate them
 - Cross dispersion with another low resolution grating or prism

Echelle Gratings: Cross Dispersion



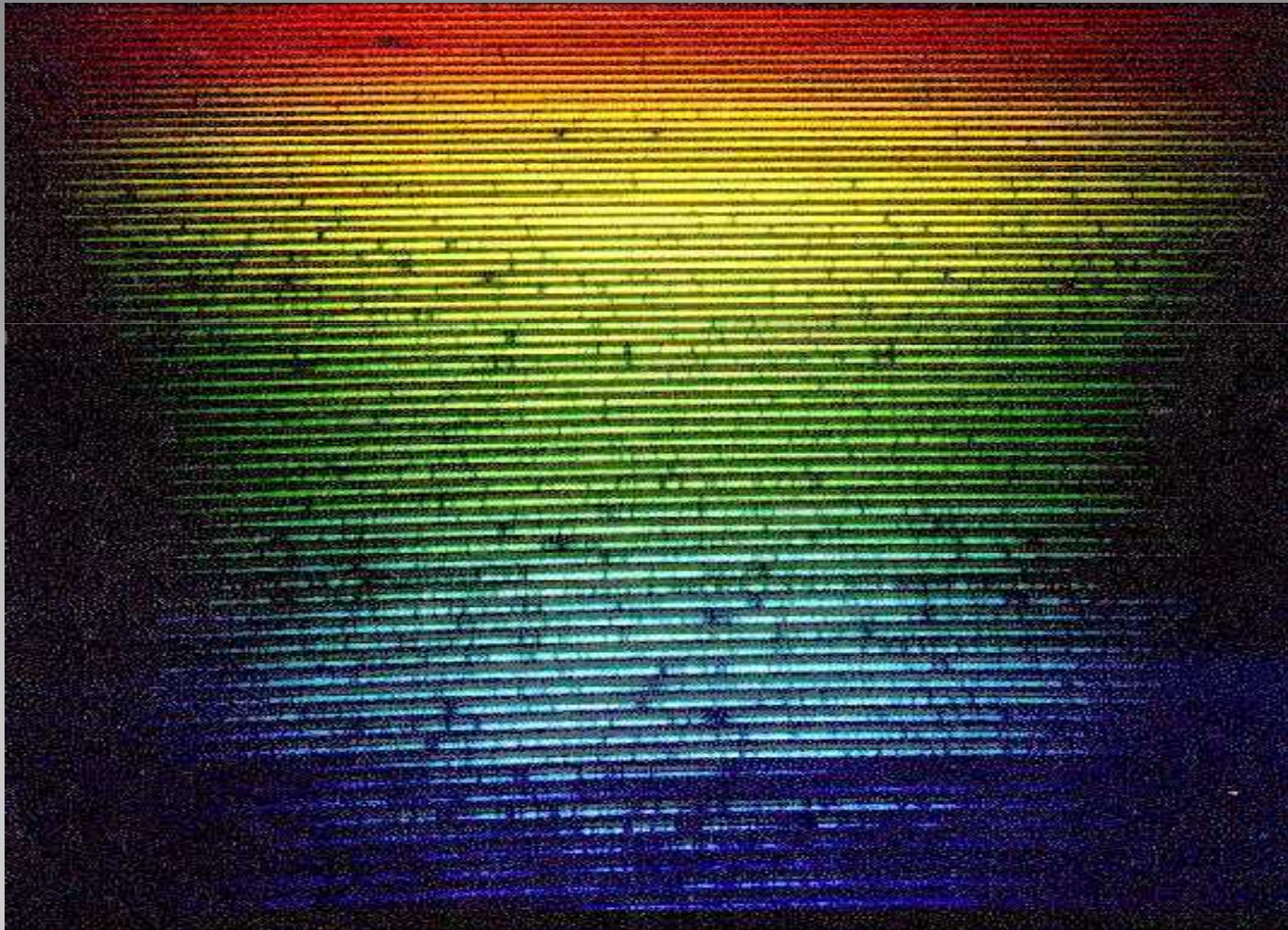


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Spectroscopy

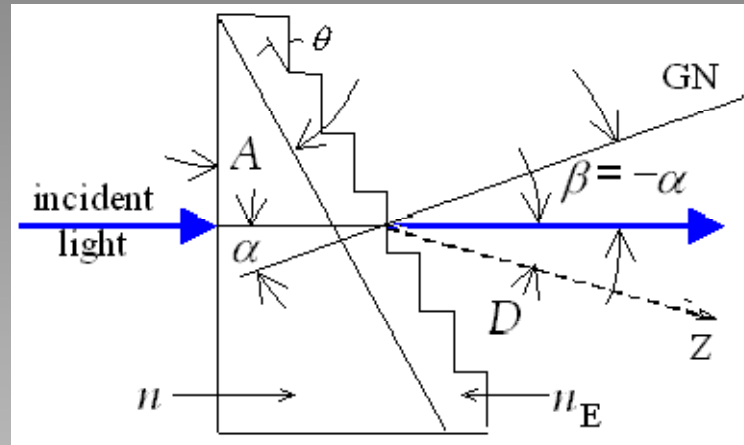
Echelle Gratings

- High resolution echellograms



Grisms

- Uses properties of both gratings and prisms



- Useful for high efficiency & low resolution
- Can keep the diffracted rays “straight through” (no deviation angle)



Optical/IR Observational Astronomy

Spectroscopy

Grisms

- Gratings

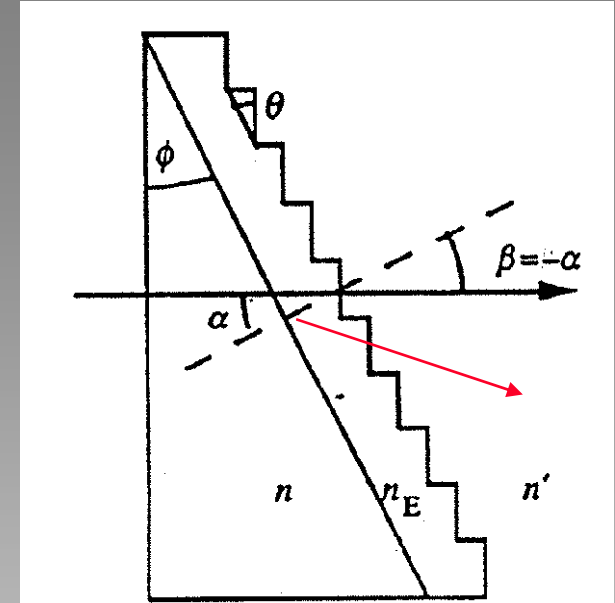
- All things comes from the basic grating equation:

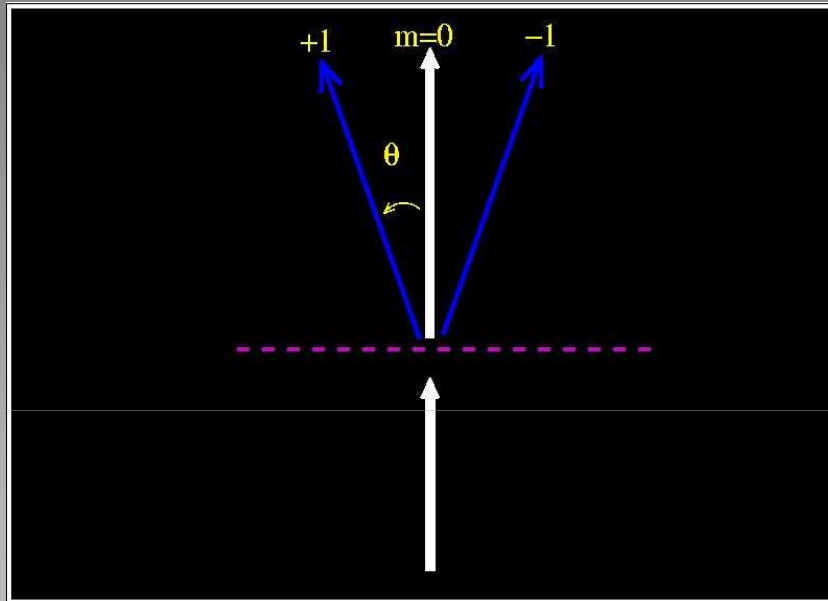
$$m \lambda = d(\sin \alpha \pm \sin \beta)$$

- Grisms (combined grating+prism)

- Blaze angle (facet angle to normal) = incidence angle (on grating)
- Normal incidence on prism face
- Blaze angle = angle of refraction of prism
- Resolution given as:

$$R \propto (n - 1) \tan \varphi$$





Transmission Grating

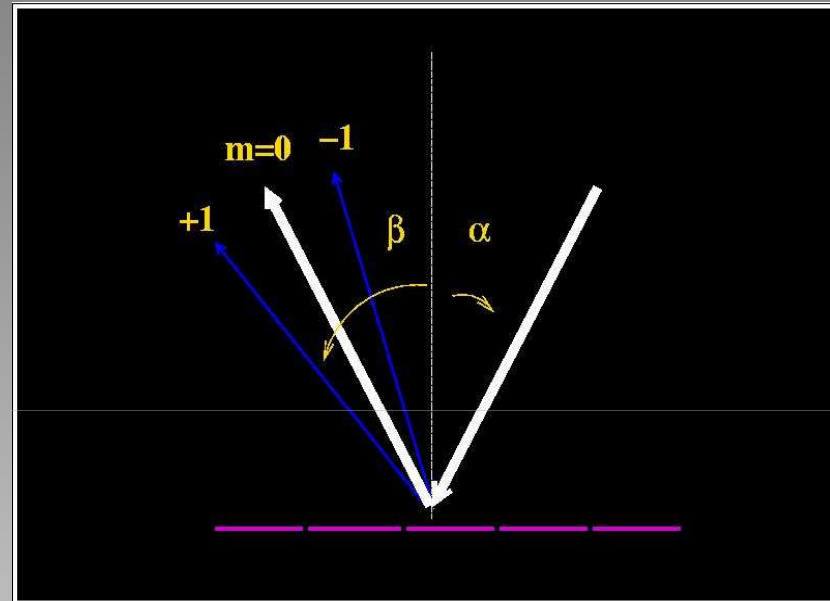
$$\sin \theta = m \cdot n \cdot \lambda$$

n = groove density (grooves or lines / mm)

m = spectral order number ($m = 1, 2, 3 \dots$)

α = angle of incidence (independent of λ)

B = angle of diffraction (λ dependent)



Reflection Grating

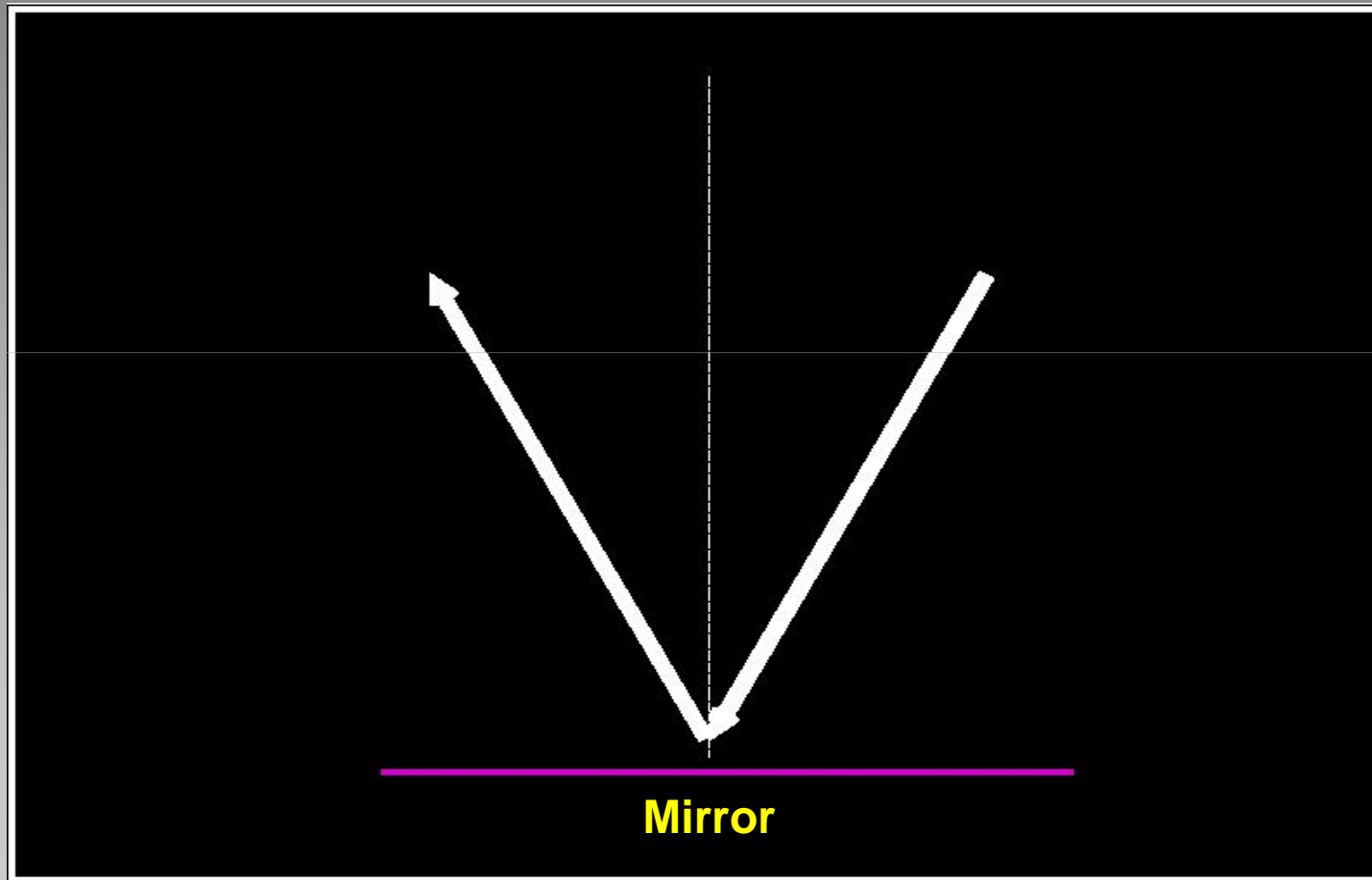
$$\sin \alpha + \sin \beta = m \cdot n \cdot \lambda$$

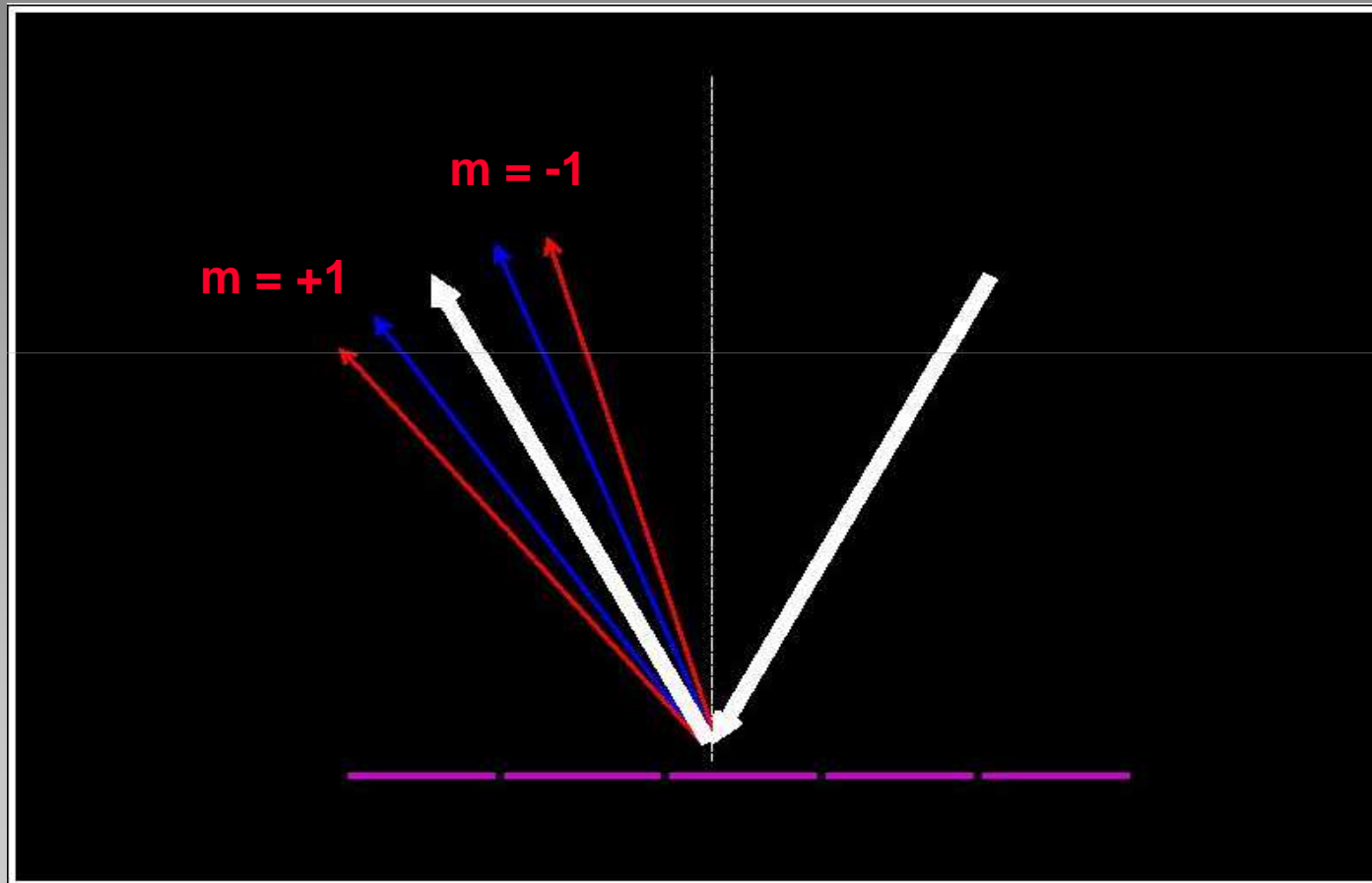


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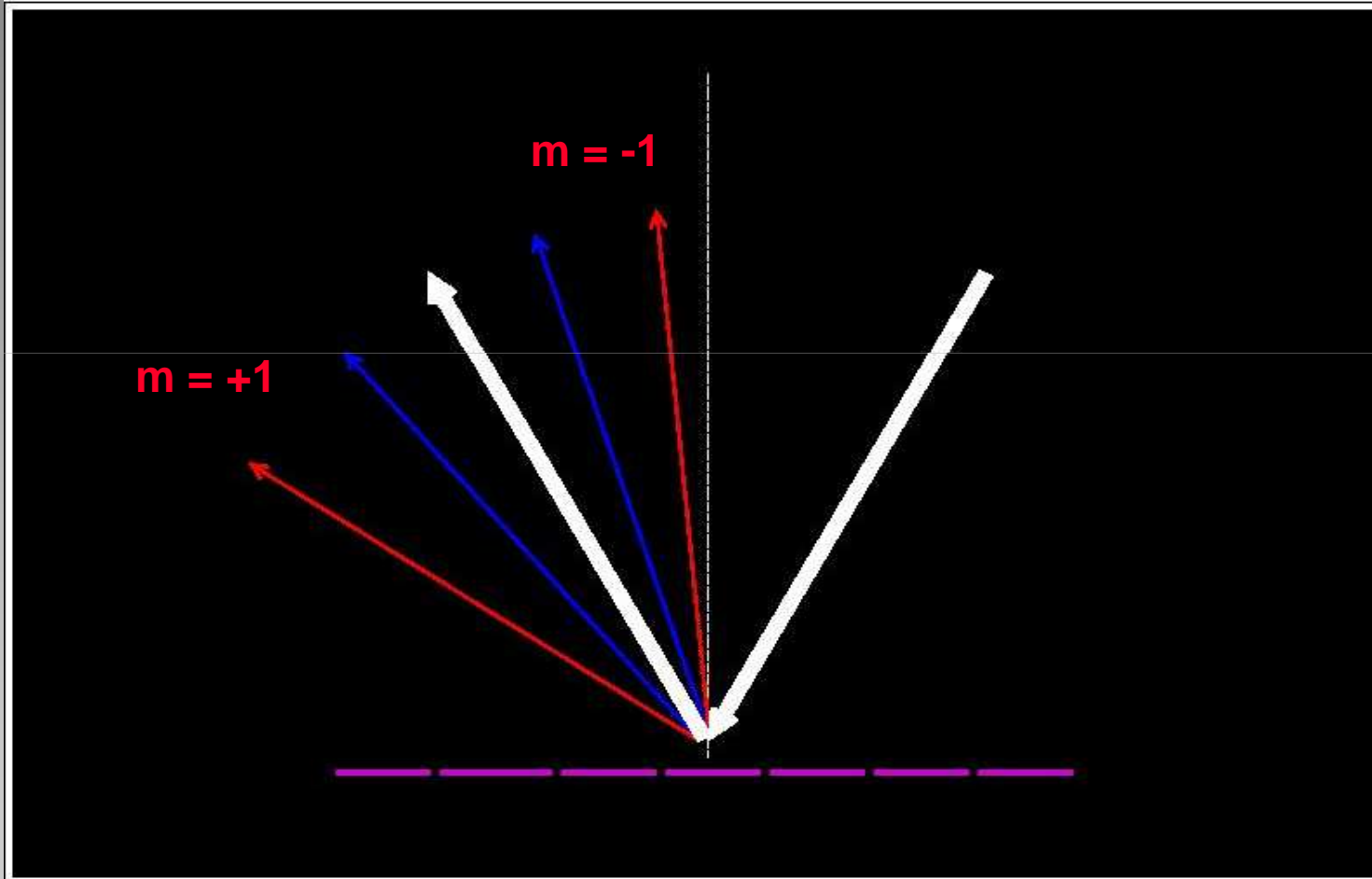
Spectroscopy

Reflection Gratings

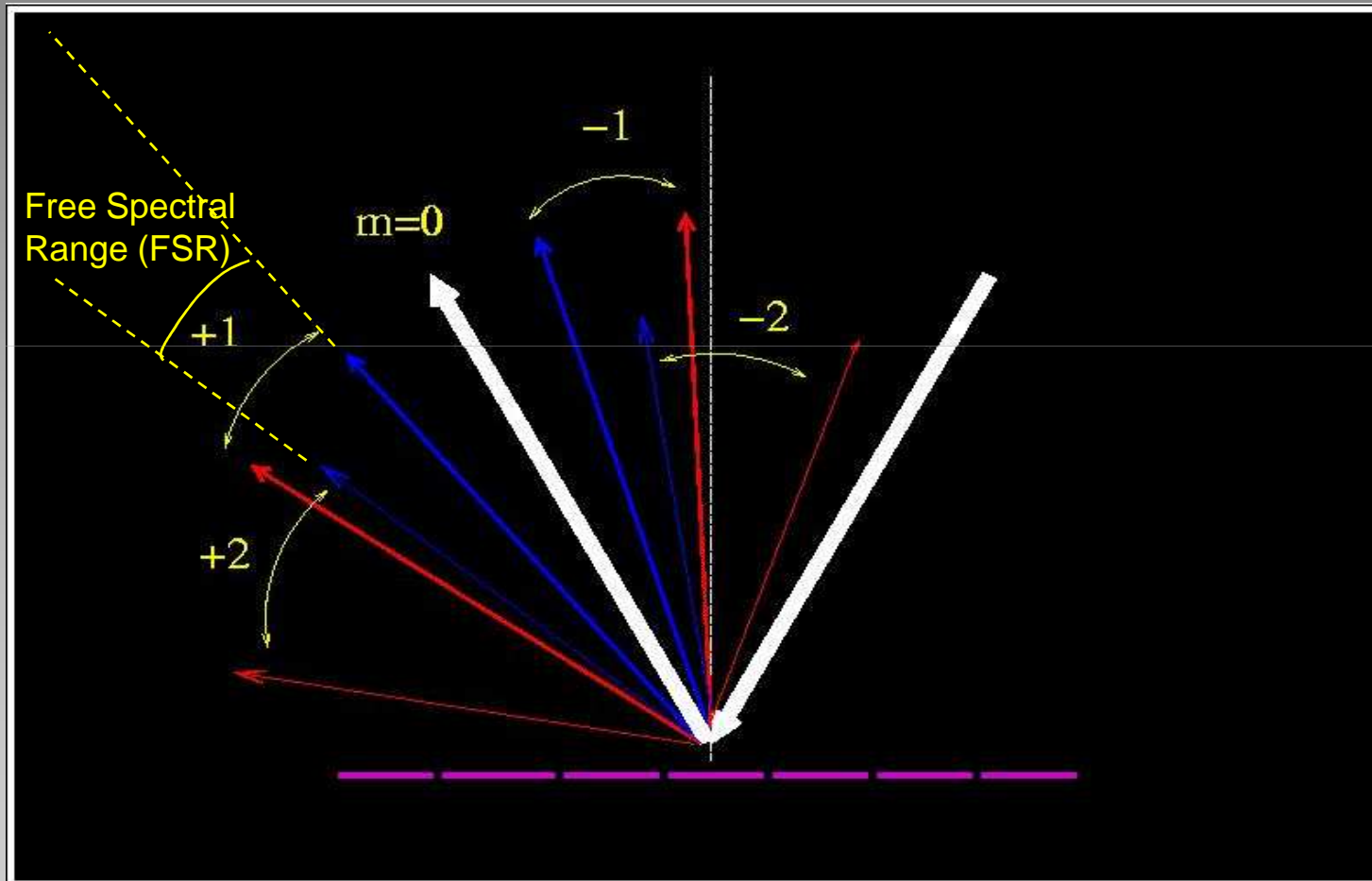


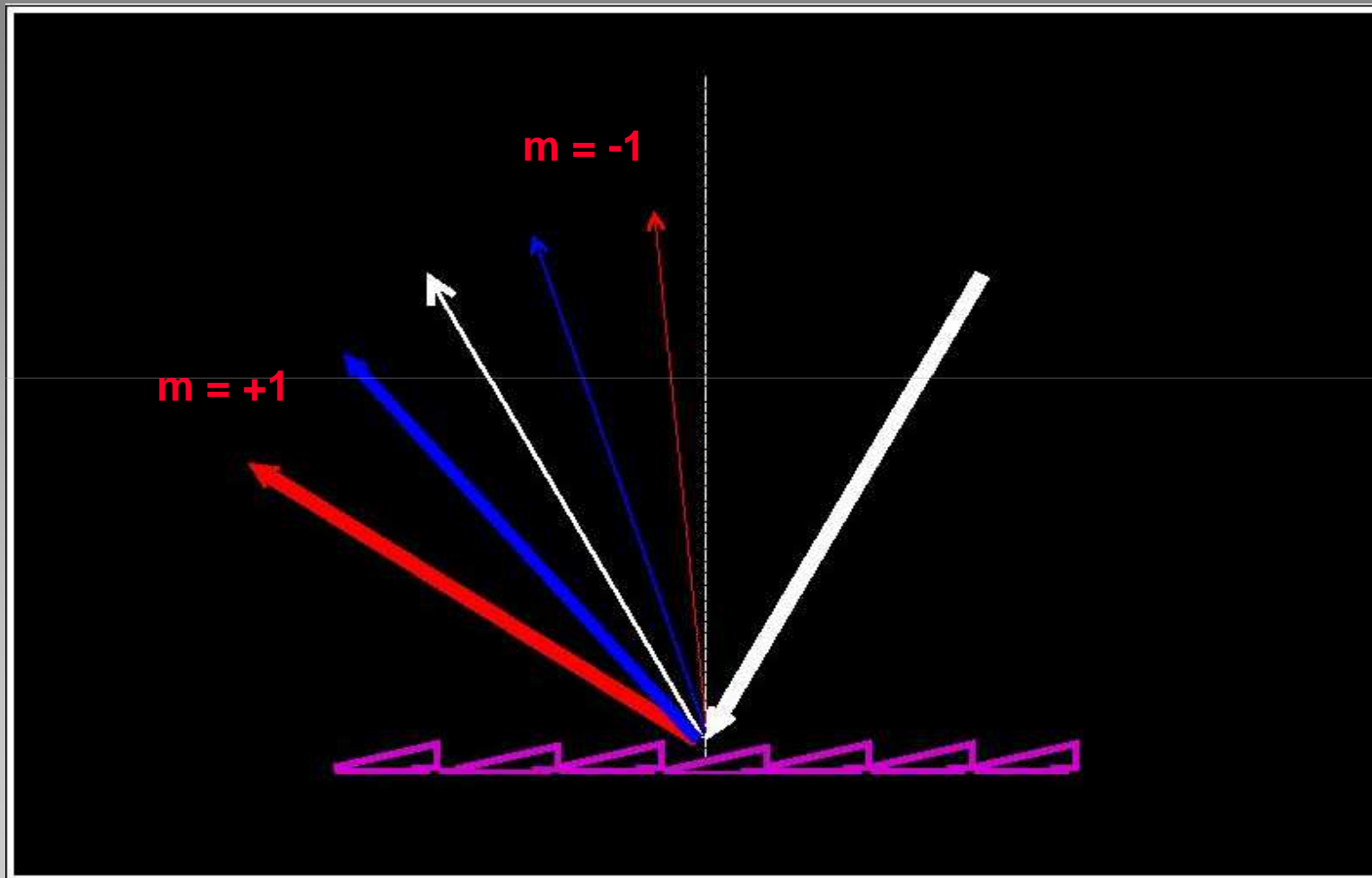


Reflection grating (n =small, *low dispersion*)

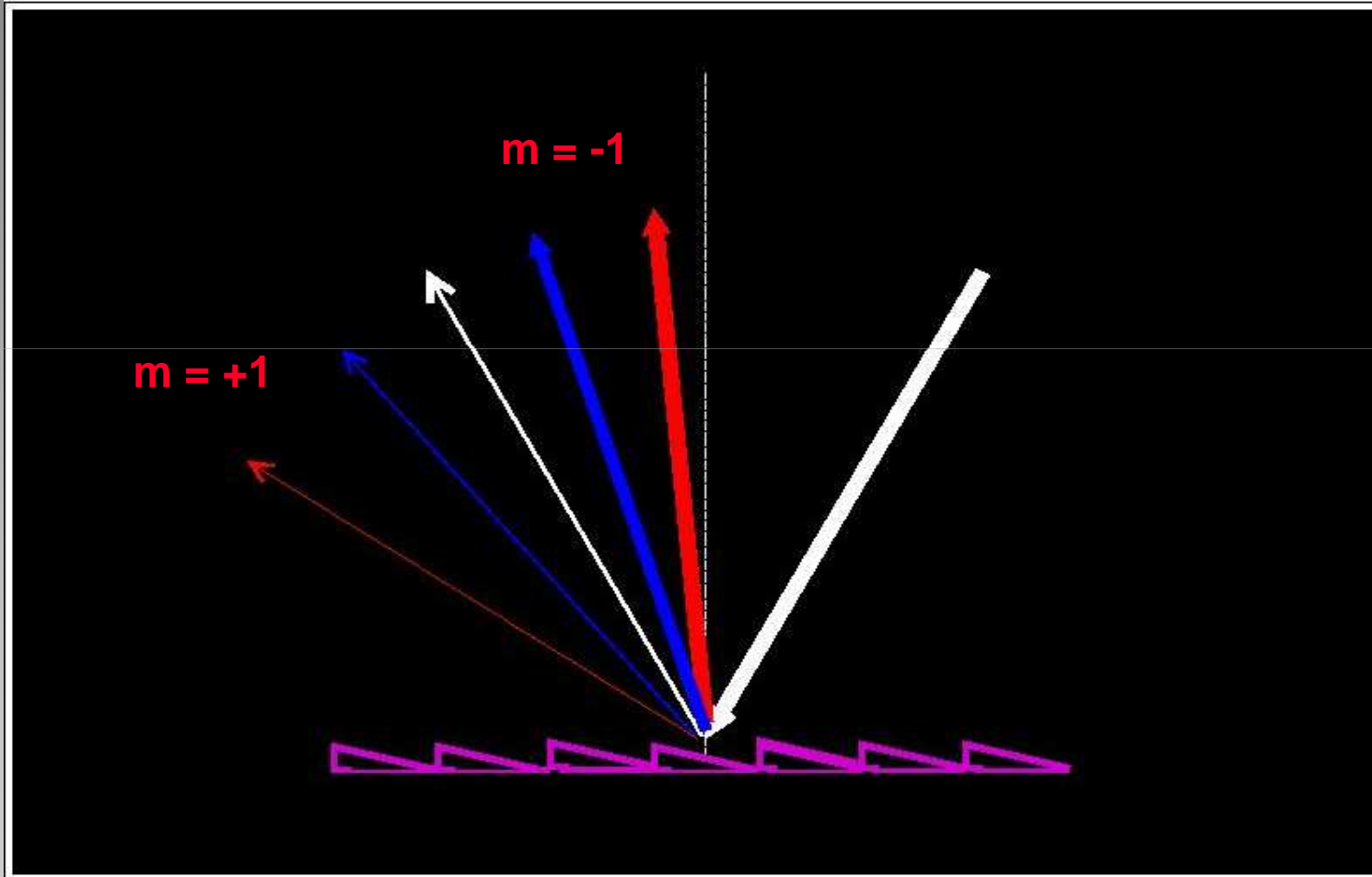


Reflecting grating (n =large, *high* dispersion)





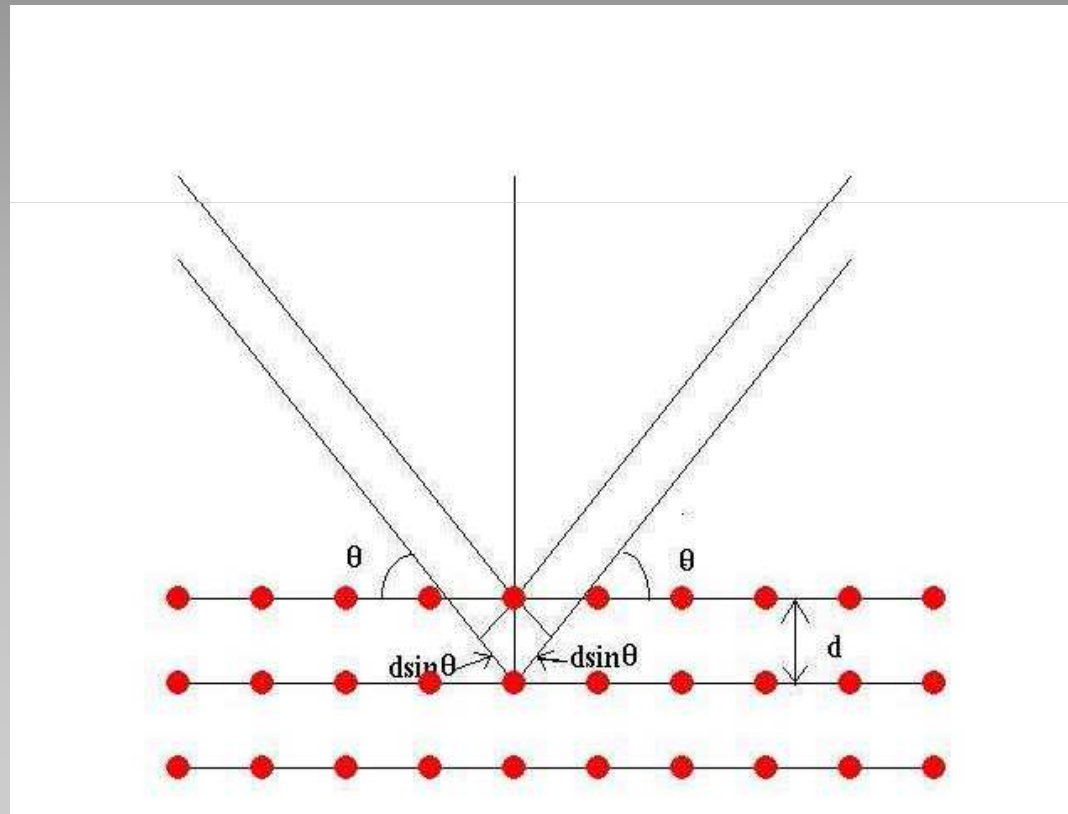
Blazed reflection grating (blazed at +1 order red)



Blazed reflecting grating (blazed at -1 order red)

“New” Grating Technology: Volume Phase Holographic Gratings

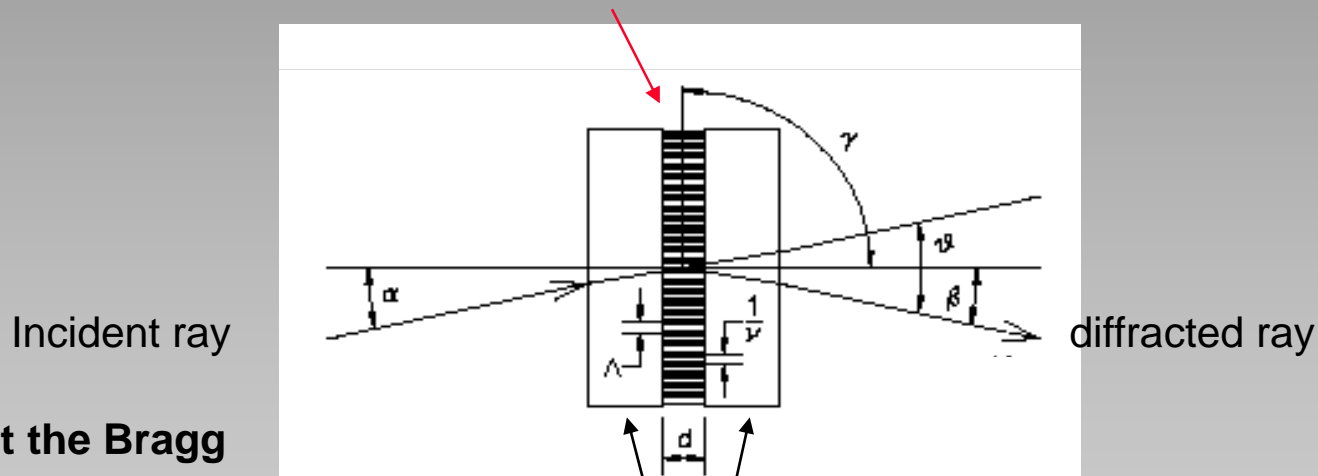
- Uses properties of Bragg diffraction in crystals



Volume Phase Holographic Gratings

- Uses a “volume” (depth) of material
- Photosensitive material (*dichromated gelatin*) has laser hologram exposed in it
- Processing produces modulation of refractive index
- Diffraction occurs in transmission through material

Dichromated gelatin layer (exposed to laser hologram)



Incident ray

diffracted ray

Work at the Bragg condition:

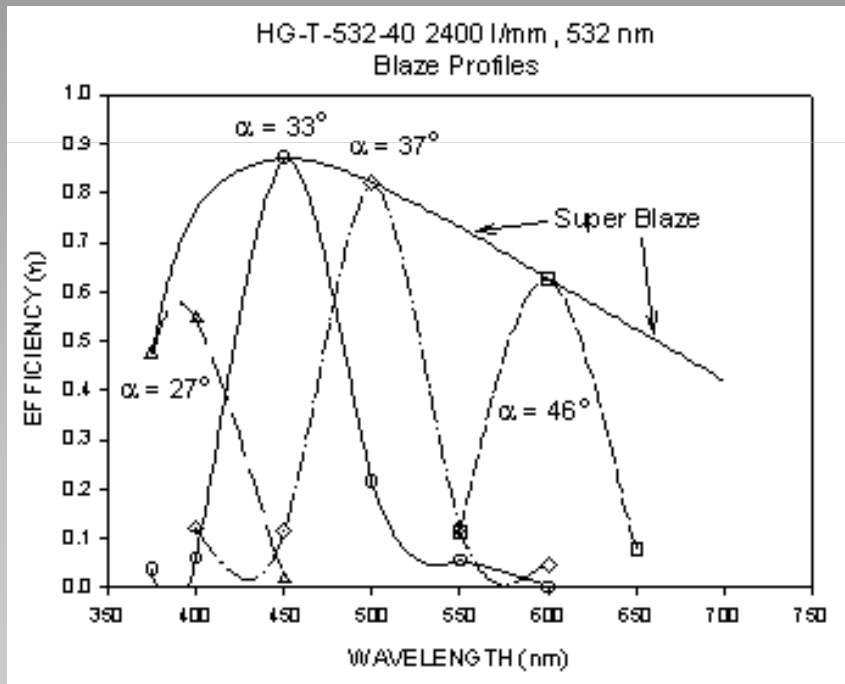
$$\alpha = \beta$$

Optically flat fused silica
or BK7 glass (can AR coat)

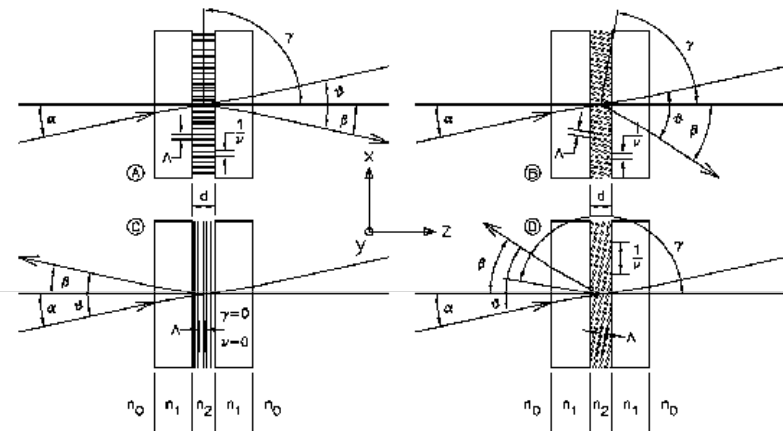
**Thickness (d) is > than depth of groove in surface-relief grating.
Implies efficiency profile governed by Bragg diffraction.**

Volume Phase Holographic Gratings

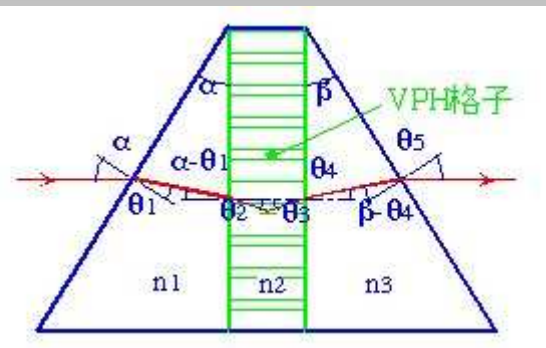
- Wavelength of operation is “tuneable”
- Efficiencies can be very high
- Now used in most modern astronomical spectrographs



VP Grating Configurations



- A. Littrow transmission configuration. B. Non-Littrow transmission configuration.
C. Non-dispersive reflection (notch filter). D. Reflection grating configuration.

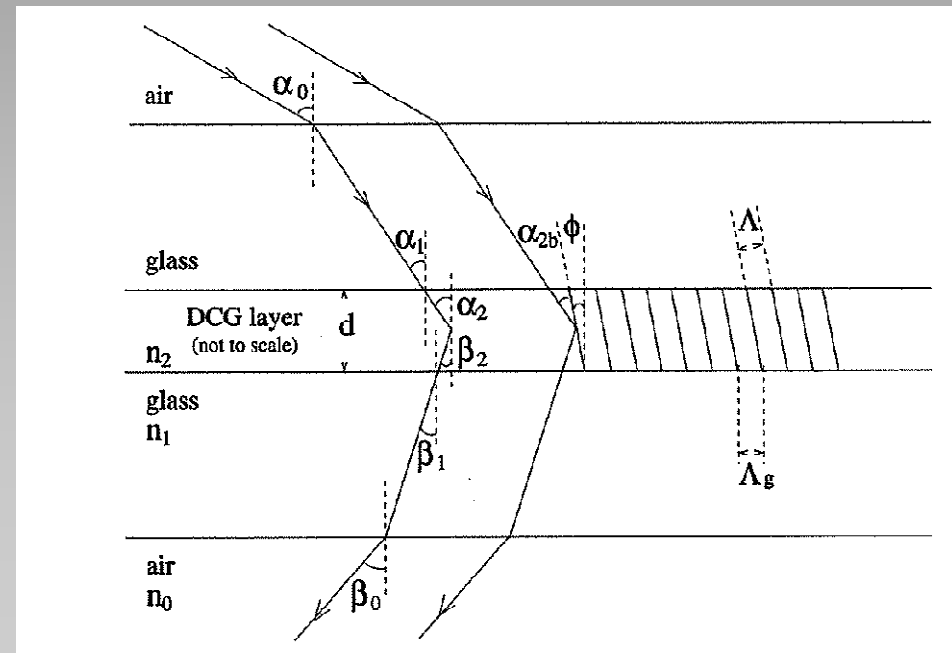


- Volume Phase Holographic gratings (VPHGs)
 - Diffraction at the Bragg condition ($\alpha = \beta$)
 - Grating equation for a VPHG is:

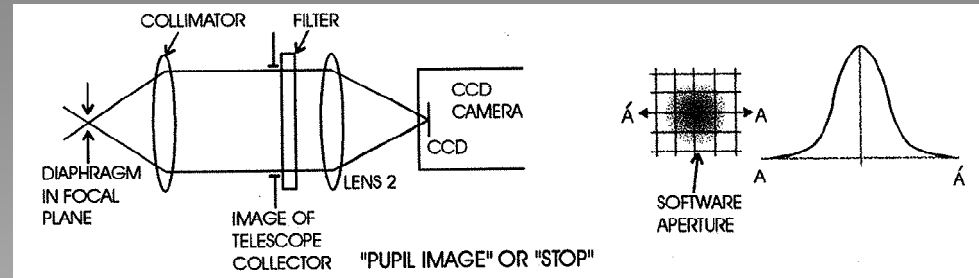
$$m\lambda = 2n_2\Lambda \sin \alpha_{2b}$$

Where Λ = index plane (“groove”) density, α_{2b} = Bragg angle w.r.t. index planes, n_2 = index of refraction of DCG layer

- Used extensively on the SALT RSS

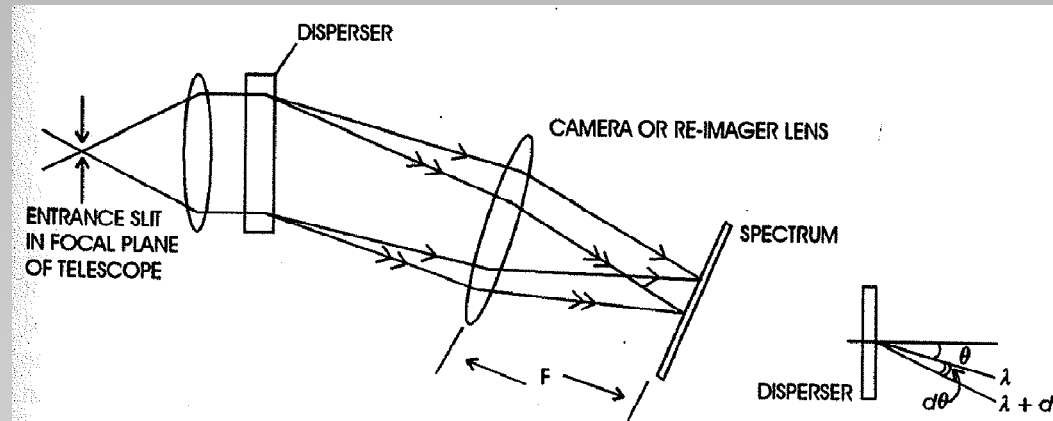


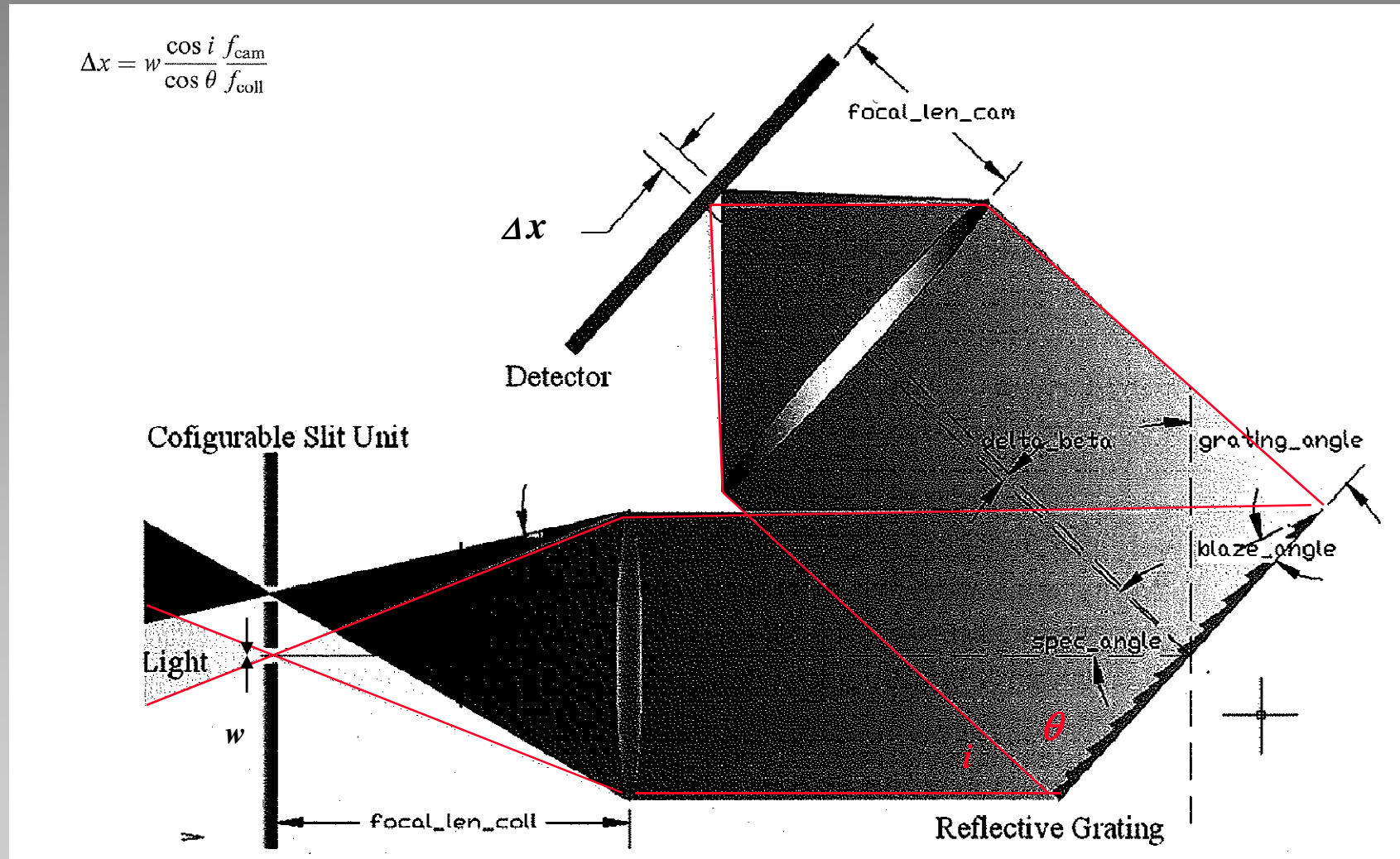
- Imaging camera (like SALTICAM, *but not quite*)



- SALTICAM does *not* collimate (beam at $f/2$)
 - » Not so important for broad band filters
- RSS imaging mode does collimate
 - » Much more important for narrow band interference filters

- Spectrometer



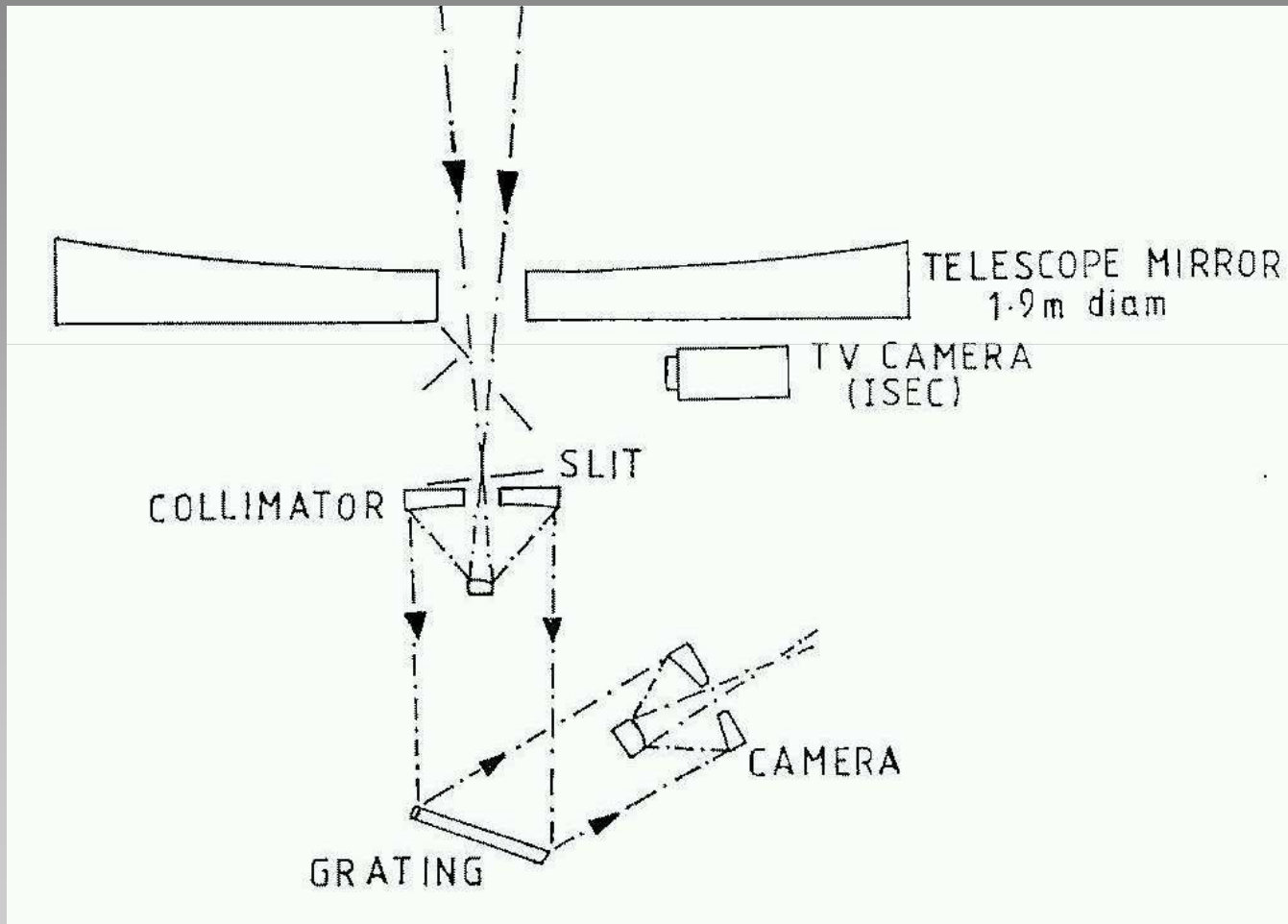




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SAAO 1.9-m Cassegrain Spectrograph

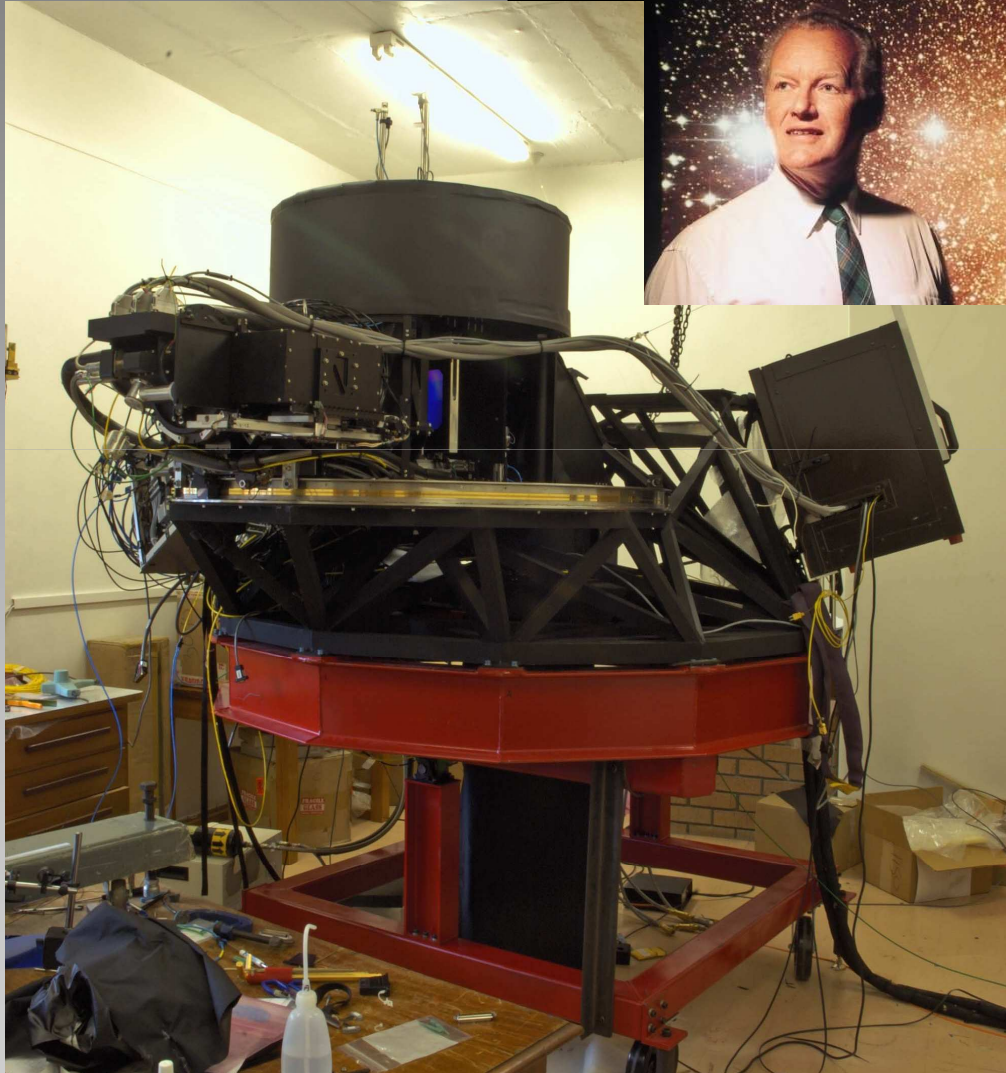




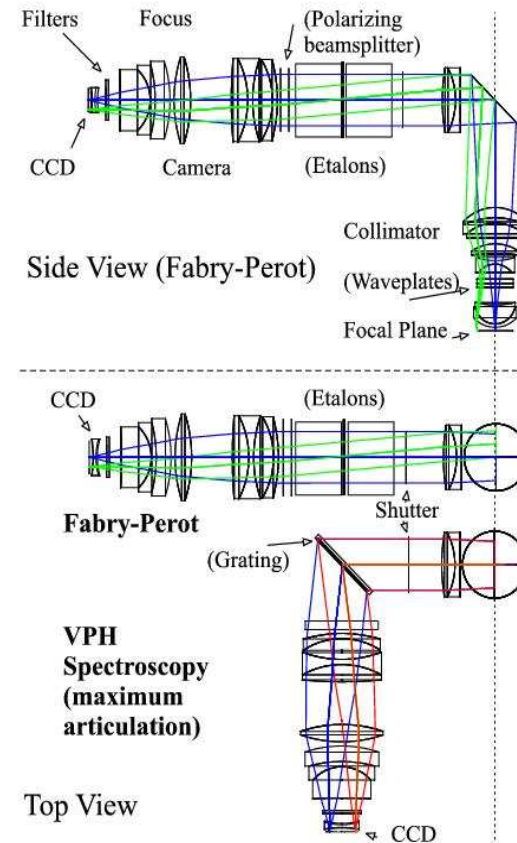
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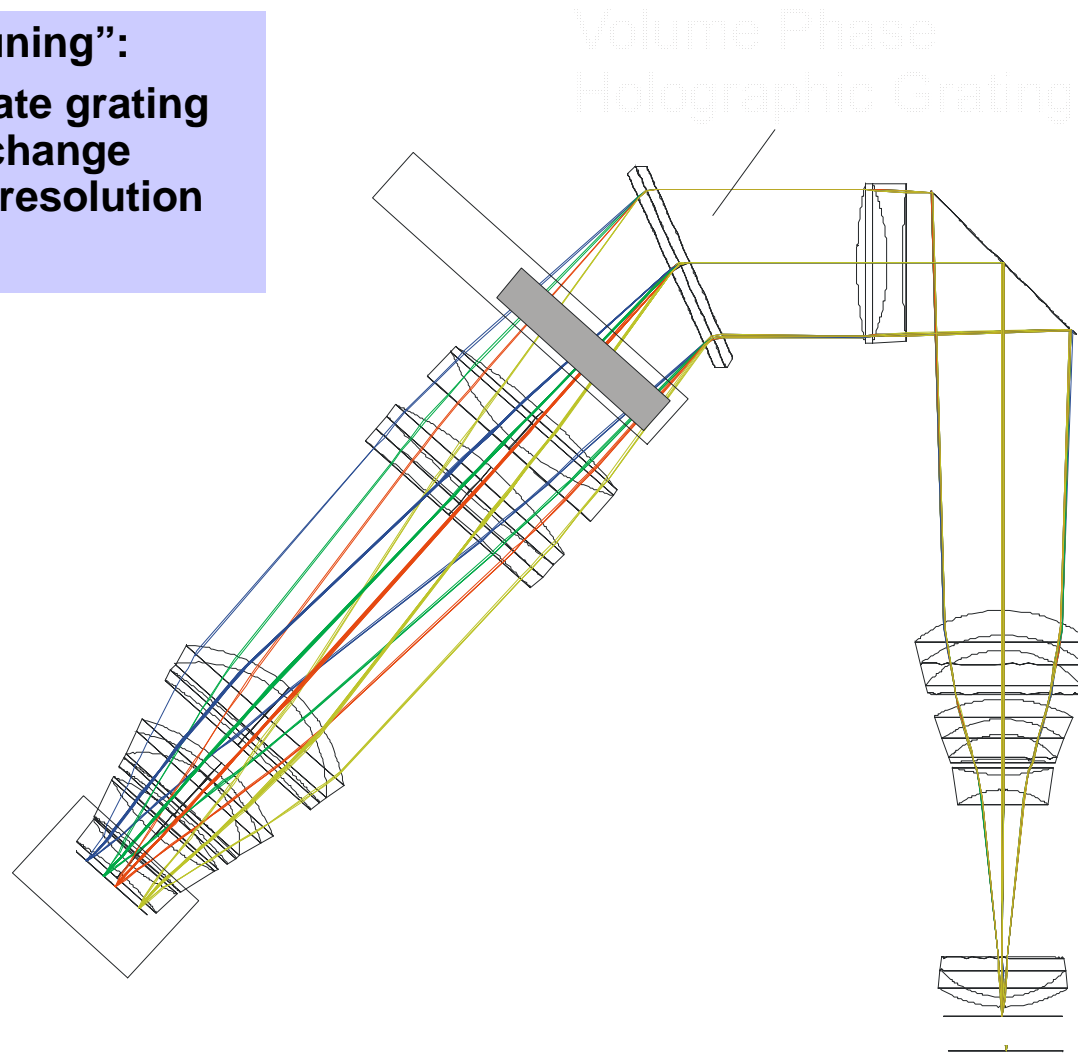
RSS: Robert Stobie Spectrograph



PFIS re-named in memory of Bob Stobie, previous SAAO Director and one of the instigators of SALT and first Chairman of SALT Board.



**RSS VPHG “tuning”:
Articulate (rotate grating
& camera) to change
wavelength & resolution**

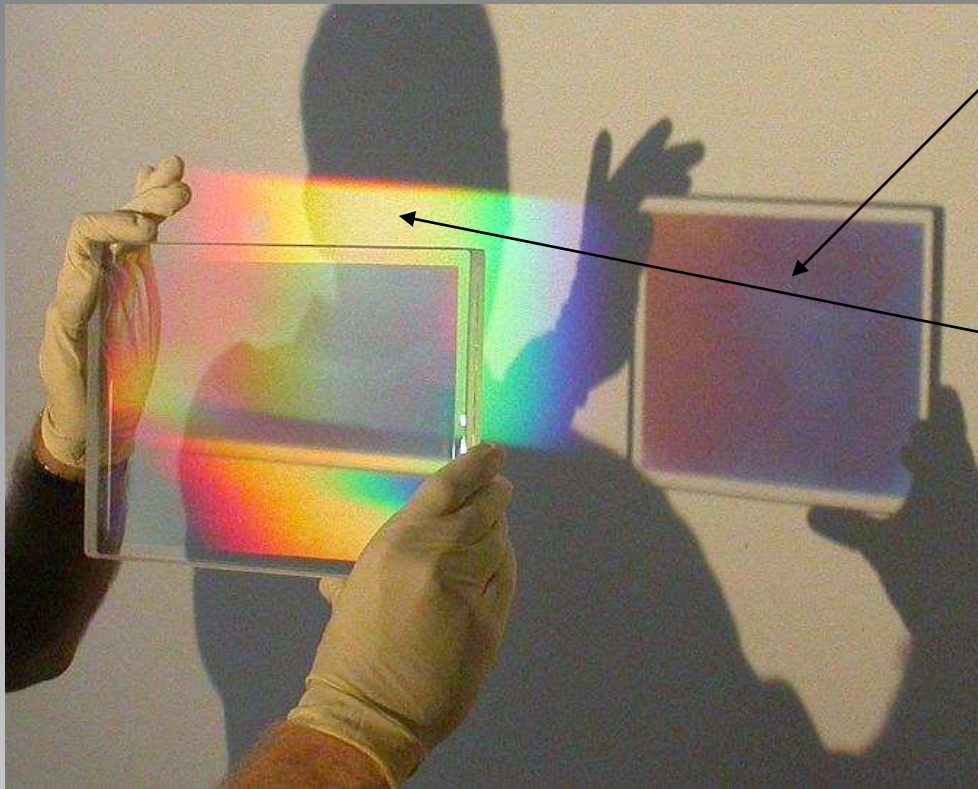




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Spectroscopy

Example of an RSS VPHG



Note the dark shadow:
implies little zero order
throughput.

Most of the light is diffracted into
the first order

1300 gr/ mm VPH

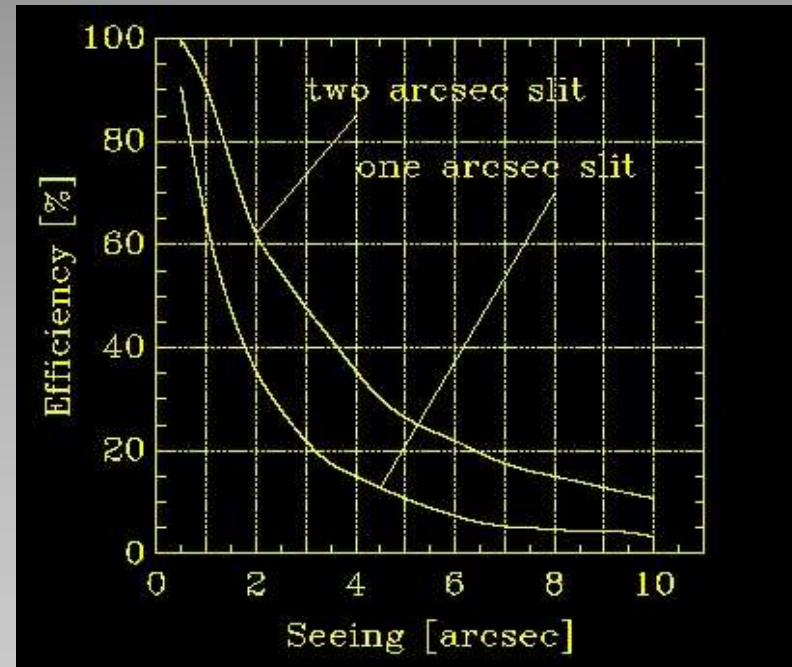
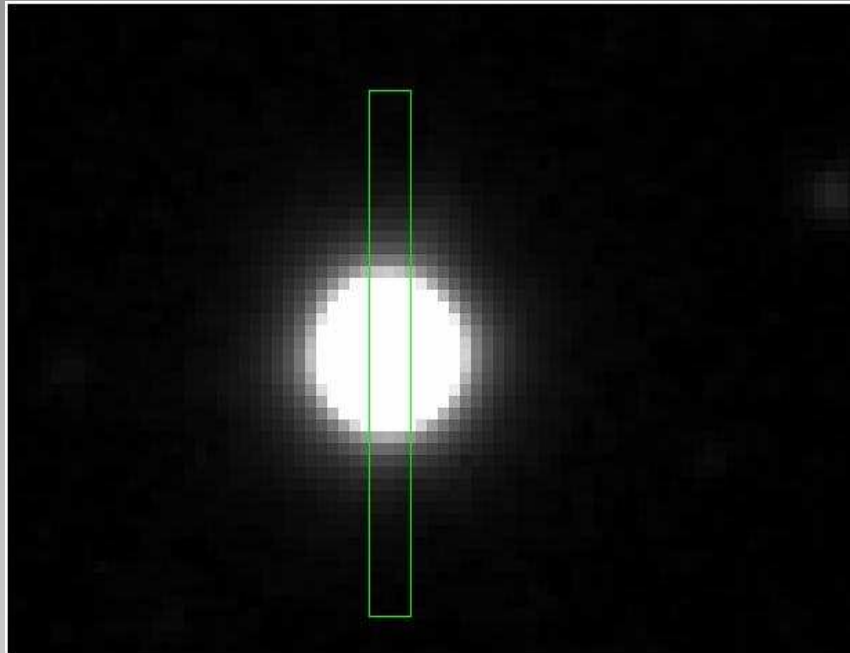


Optical/IR Observational Astronomy

Spectroscopy

Effect of bad focus and seeing in Spectroscopy

Slit width/ Slit loss



$$\Delta x = w \frac{\cos i}{\cos \theta} \frac{f_{\text{cam}}}{f_{\text{coll}}}$$

The Spectrograph Slit

The Entrance slit

- Selects light from a particular object in the telescope's focal plane.

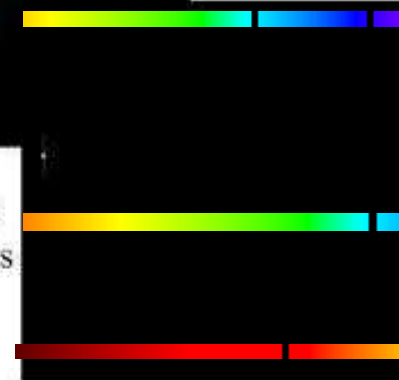


Telescope
FOV

...masked by a
long slit (focus
on galaxy)

- Helps to give good "shape" and resolution to the spectrum produced by the instrument.

...masked by a
short slit (focus
on a star)





Optical/IR Observational Astronomy

Spectroscopy

Signal to Noise

S/N for object measured in aperture with radius r : $n_{\text{pix}} = \#$
of pixels in the aperture = πr^2

Signal $\longleftrightarrow R_* t$

Noise $\leftrightarrow \left[R_* \cdot t + R_{\text{sky}} \cdot t \cdot n_{\text{pix}} + \left(RN + \frac{\text{gain}}{2} \right)^2 \cdot n_{\text{pix}} + \text{Dark} \cdot t \cdot n_{\text{pix}} \right]^{\frac{1}{2}}$

$\sqrt{(R_* \cdot t)^2}$

Noise from the dark current in aperture

Readnoise in aperture

Noise from sky e^- in aperture

All the noise terms added in quadrature
Note: always calculate in e^-



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Spectroscopy

Signal to Noise in mags

Side Issue: $S/N \Leftrightarrow \delta mag$

$$\begin{aligned}
 m \pm \delta(m) &= c_o - 2.5 \log(S \pm N) \\
 &= c_o - 2.5 \log\left[S\left(1 \pm \frac{N}{S}\right)\right] \\
 &= \underbrace{c_o - 2.5 \log(S)}_m - \underbrace{2.5 \log\left(1 \pm \frac{N}{S}\right)}_{\delta m}
 \end{aligned}$$

$$\delta(m) \approx 2.5 \log\left(1 + \frac{1}{S/N}\right)$$

$$= \frac{2.5}{2.3} \left[\frac{N}{S} - \frac{1}{2} \left(\frac{N}{S}\right)^2 + \frac{1}{3} \left(\frac{N}{S}\right)^3 - \dots \right]$$

$$\approx 1.087 \left(\frac{N}{S}\right) \leftarrow \text{Fractional error}$$

Note: in log +/- not symmetric

This is the basis of people referring to +/- 0.02mag error as "2%"

S/N	δmag
2	0.44
10	0.10
100	0.01



Signal to Noise

S/N Calculations

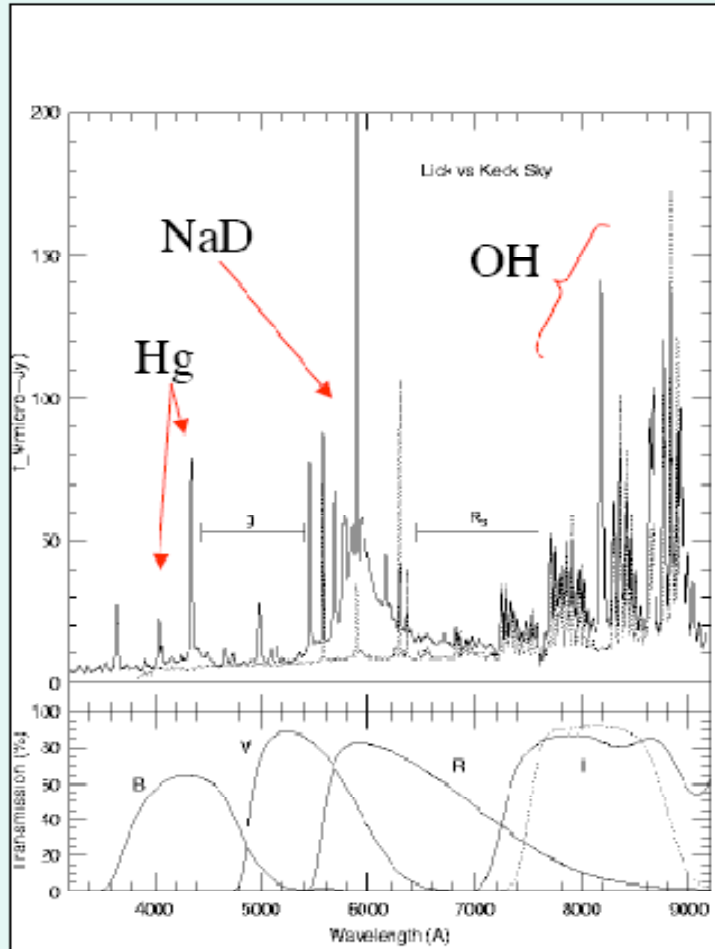
- So, what do you do with this?
 - Demonstrate feasibility
 - Justify observing time requests
 - Get your observations right
 - Estimate limiting magnitudes
 - Discover problems with instruments, telescopes or observations
- So, how do you calculate it?
 - Could use the S/N formula
 - need to know which regime you're operating in*
 - need to know the various parameter (pixel scale, gain, RN, Dark count, etc.)*
 - Usually there are simple tools available to calculate all of this for you!

Sky Background

$$R_{\text{sky}}$$

Signal from the sky background is present in every pixel of the aperture. Because each instrument generally has a different pixel scale, the sky brightness is usually tabulated for a site in units of $\text{mag}/\text{arcsecond}^2$.

$$(\text{mag}/\square)$$



Lunar age (days)	U	B	V	R	I
0	22.0	22.7	21.8	20.9	19.9
3	21.5	22.4	21.7	20.8	19.9
7	19.9	21.6	21.4	20.6	19.7
10	18.5	20.7	20.7	20.3	19.5
14	17.0	19.5	20.0	19.9	19.2



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Two most common cases

S/N - some limiting cases. Let's assume CCD with Dark=0, well sampled read noise.

$$\frac{R_* t}{\left[R_* \cdot t + R_{\text{sky}} \cdot t \cdot n_{\text{pix}} + (RN)^2 \cdot n_{\text{pix}} \right]^{\frac{1}{2}}}$$

Bright Sources: $(R_* t)^{1/2}$ dominates noise term

$$S/N \approx \frac{R_* t}{\sqrt{R_* t}} = \sqrt{R_* t} \propto t^{\frac{1}{2}}$$

Sky Limited ($\sqrt{R_{\text{sky}} t} > 3 \times RN$): $S/N \propto \frac{R_* t}{\sqrt{n_{\text{pix}} R_{\text{sky}} t}} \propto \sqrt{t}$

Note: seeing comes in with n_{pix} term

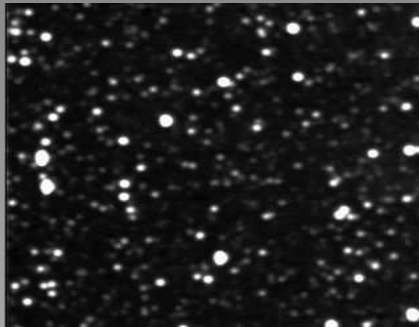


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Spectroscopy

Effect of bad focus and seeing in Spectroscopy

Signal-to-Noise for typical SALT case

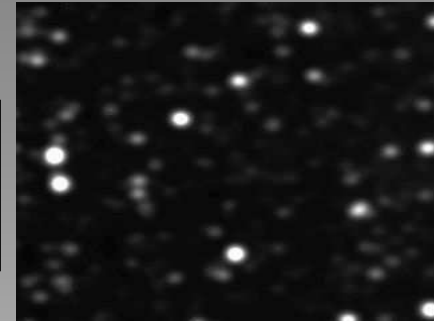


Good seeing

Bright object: not so affected

$$\frac{S}{N} \approx \frac{O}{\sqrt{O}}$$

~



Bad seeing:
n times bigger seeing psf

$$\frac{S}{N} \approx \frac{O}{\sqrt{O}}$$

Faint object [Background (sky) dominating]: badly affected

$$\frac{S}{N} \approx \frac{O}{\sqrt{B}}$$

>

$$\frac{S}{N} \approx \frac{O}{\sqrt{Bn^2}} = \frac{O}{\sqrt{B}} \frac{1}{n}$$



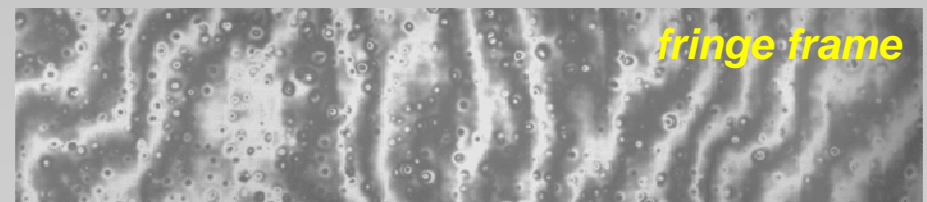
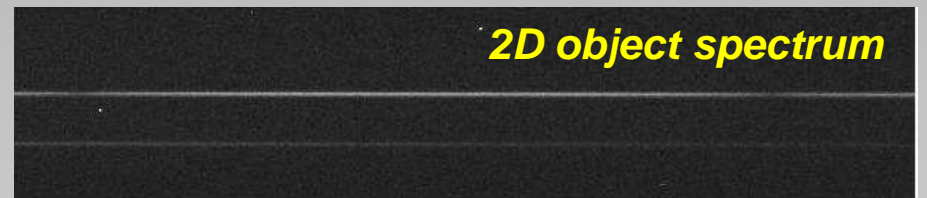
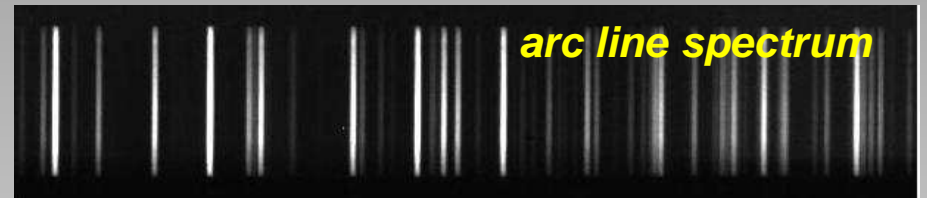
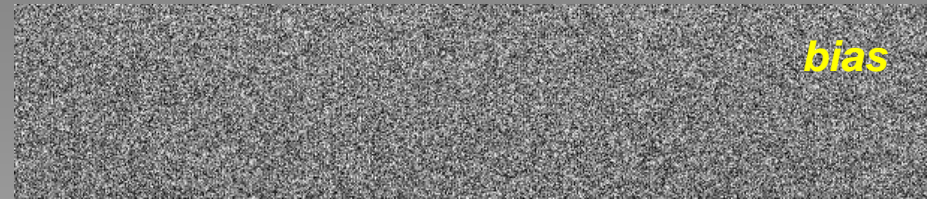
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Spectroscopy

Reducing Spectroscopic Data

Calibration frames:

- * Bias frame (electronic DC level)
 - sometimes derived from “overscan” or “prescan”
- * Flat field (detector sensitivity)
- * Arc Lamp (wavelength calibration)
- * Flux Standard
 - flux from $electrons\ cm^{-2}\ sec^{-1}\ \text{\AA}^{-1}$ to $ergs/cm^{-2}\ sec^{-1}\ \text{\AA}^{-1}$
- * Fringe frame correction



Basic reduction procedures:

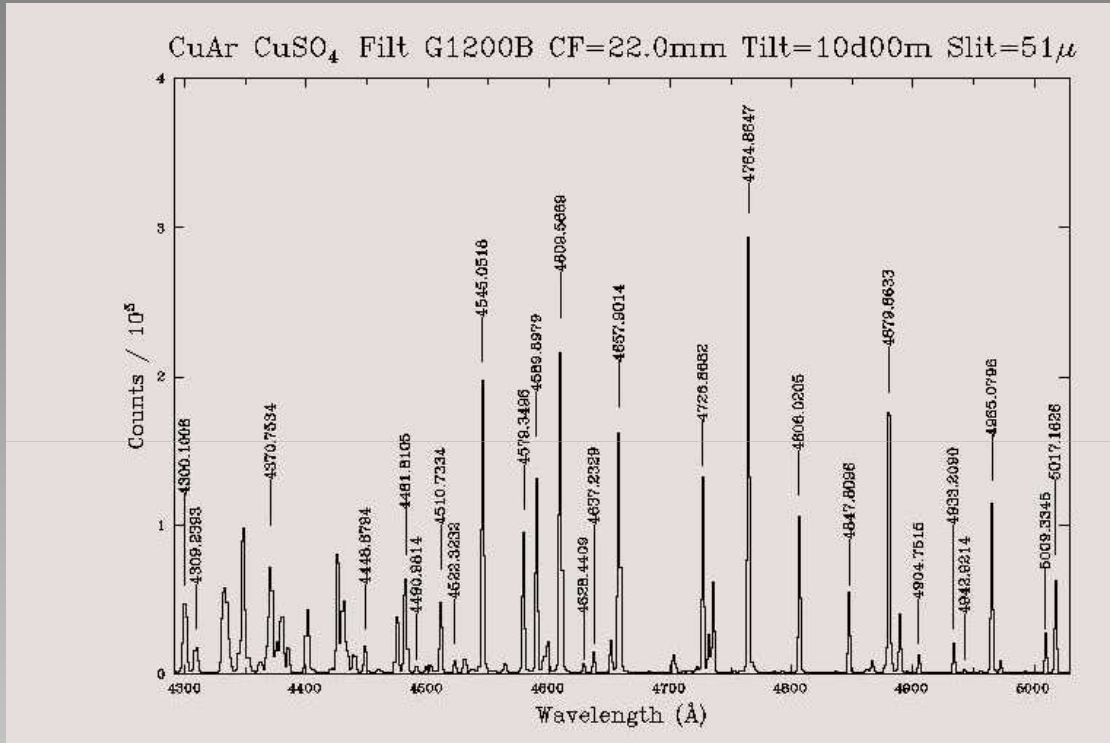
1. Bias correction
2. Flat fielding
3. Cosmic ray removal
4. Wavelength determination
5. Background subtraction
6. Spectrum extraction
7. Flux calibration (remove atm. effects)



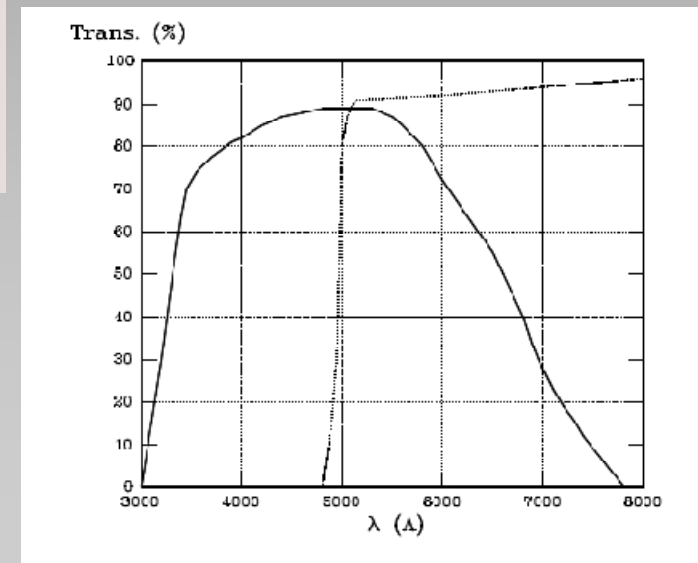
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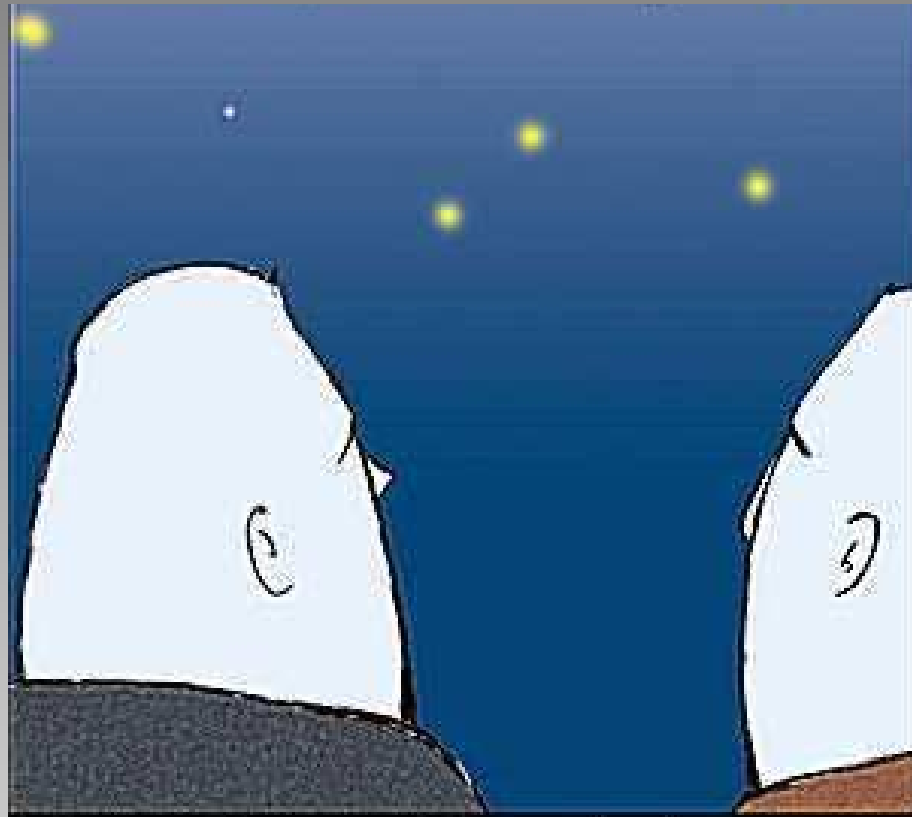
Spectroscopy

Wavelength Calibrations



- Determine wavelength as a function of pixel number on the detector
- Typically use 3rd order polynomial fit to known wavelengths
- Need to have order blocking filters to exclude 2nd order overlap in the 1st order, 3rd in 2nd, etc.





"I've always found the stars
dissapointingly small."