

Worked Examples Set 2

Q.1. Application of Maxwell's eqns. [Griffiths Problem 7.42]

In a perfect conductor the conductivity σ is infinite, so from 'Ohm's law' $\vec{J} = \sigma \vec{E}$, $\vec{E} = 0$. Any net charge must be on the surface (as in electrostatics for an imperfect conductor).

(a) Show that the magnetic field is constant inside a perfect conductor.

Solution: Faraday's law: $\nabla \times \vec{E} = -\partial \vec{B} / \partial t$ so if $\vec{E} = 0$ then $\partial \vec{B} / \partial t = 0$, i.e. $\vec{B}$ is independent of t , i.e. constant.

(b) Show that the magnetic flux through a perfectly conducting loop is constant.

Solution: Faraday's law in integral form is $\oint_C \vec{E} \cdot d\vec{l} = -\frac{d\Phi_m}{dt}$

Q.1. Solution (b) [continued]

Inside the perfectly conducting loop $\vec{E} = 0$, so $\frac{d\Phi_m}{dt} = 0$

i.e. Magnetic flux Φ_m is independent of t , i.e. constant.

(c) A superconductor is a perfect conductor with the additional property that the constant $\vec{B}$ -field inside is actually zero.

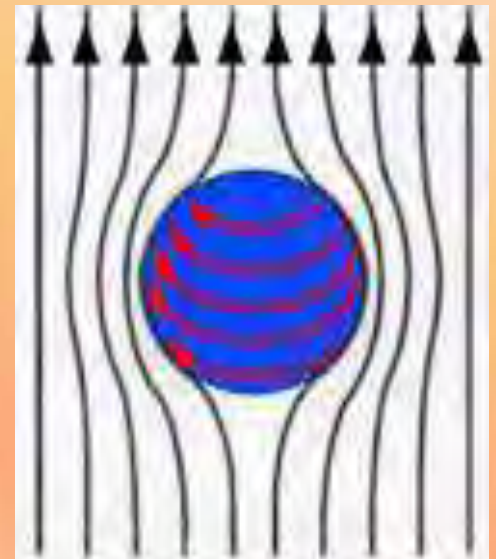
Show that the current in a superconductor is confined to the surface (there are no volume currents).

Solution: Maxwell-Ampere law:

$$\nabla \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \partial \vec{E} / \partial t$$

so if $\vec{E} = 0$ and $\vec{B} = 0$ (given) then $\vec{J} = 0$

i.e. there are no volume currents and so any currents must be surface currents ($\vec{K}$).



Q.2. Potentials and Gauge Transformations

[Griffiths Problems 10.3 and 10.5]

- (a) Find the fields, and the charge and current distributions, corresponding to potentials $V(\vec{r}, t) = 0$, $\vec{A}(\vec{r}, t) = -\frac{1}{4\pi\epsilon_0} \frac{qt}{r^2} \hat{r}$
- (b) Use the gauge function $\lambda = -\frac{1}{4\pi\epsilon_0} \frac{qt}{r}$ to transform the potentials in (a), and comment on the results.

Solution: (a) $\vec{E} = -\nabla V - \frac{\partial \vec{A}}{\partial t} = 0 + \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{r} = \frac{q}{4\pi\epsilon_0 r^2} \hat{r}$

$\vec{B} = \nabla \times \vec{A} = 0$ [$\vec{A} = A\hat{r}$ (radial component only) which does not depend on θ or ϕ , i.e. $\partial A_r / \partial \theta = 0$, $\partial A_r / \partial \phi = 0$]

$$\nabla \times \mathbf{A} = \left(\frac{1}{r \sin \theta} \left[\frac{\partial}{\partial \theta} (\sin \theta A_\phi) - \frac{\partial A_\theta}{\partial \phi} \right], \frac{1}{r} \left[\frac{1}{\sin \theta} \frac{\partial A_r}{\partial \phi} - \frac{\partial}{\partial r} (r A_\phi) \right], \frac{1}{r} \left[\frac{\partial}{\partial r} (r A_\theta) - \frac{\partial A_r}{\partial \theta} \right] \right)$$

Fields of a stationary point charge, strange potentials.

Q.2. Solution [continued]

Charge distribution: we have a point charge q , which can be written as a (3D) delta function : $\rho = q\delta^3(\vec{r})$;
current density $\vec{J} = 0$ (since $\vec{B} = 0$)

(b) Gauge transformation: $\lambda = -\frac{1}{4\pi\epsilon_0} \frac{qt}{r}$

$$V' = V - \frac{\partial\lambda}{\partial t} = 0 - \left(-\frac{1}{4\pi\epsilon_0} \frac{q}{r}\right) = \frac{q}{4\pi\epsilon_0 r}$$

$$\begin{aligned}\vec{A}' &= \vec{A} + \nabla\lambda = -\frac{1}{4\pi\epsilon_0} \frac{qt}{r^2} \hat{r} + \left(-\frac{1}{4\pi\epsilon_0} qt\right) \left(-\frac{1}{r^2} \hat{r}\right) \\ &= -\frac{1}{4\pi\epsilon_0} \frac{qt}{r^2} \hat{r} + \frac{1}{4\pi\epsilon_0} \frac{qt}{r^2} \hat{r} = 0\end{aligned}$$

So this transformation transforms the strange potentials of (a) into the more standard potentials of a static point charge.

Q.3. Retarded Potentials

[Griffiths, Example 10.2]

A long straight wire has $I(t) = 0$, $t \leq 0$ and $I(t) = I_0$, $t > 0$ ('switch on' at $t = 0$). Find the resulting $\vec{E}$ and $\vec{B}$ fields.

Solution: Assume wire is neutral, so $V = 0$.

Retarded vector potential at P (see figure) is

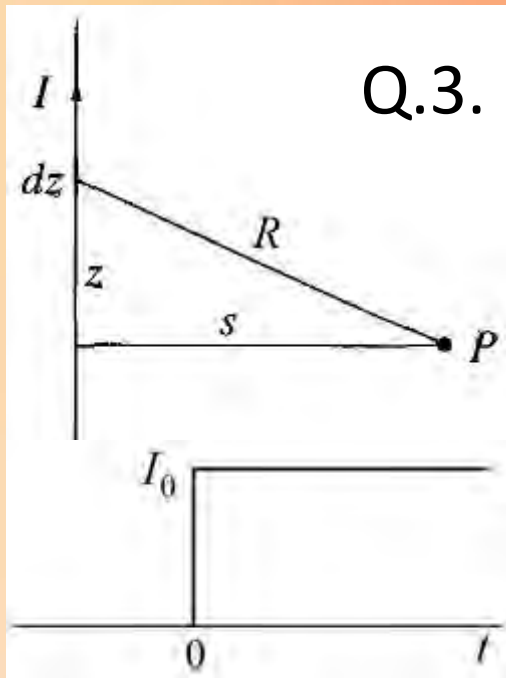
$$\vec{A}(s, t) = \frac{\mu_0}{4\pi} \hat{z} \int_{-\infty}^{+\infty} \frac{I(t_r)}{R} dz \quad [\text{in same direction } (\hat{z}) \text{ as current } I]$$

For the field to reach P takes time s/c , so for $t < s/c$, $\vec{A} = 0$

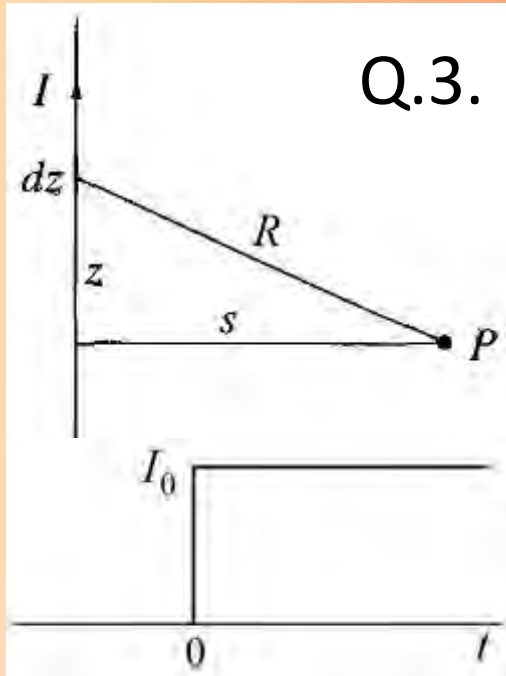
For $t > s/c$ the 'signal' takes time R/c to travel along R , so range of z contributing to $\vec{A}(s, t)$ at P is $|z| \leq \sqrt{(ct)^2 - s^2}$

Outside this range $t_r = t - R/c < 0$ so $I(t_r) = 0$

[Note the $|z| \Rightarrow$ contributions from above & below the $\perp$ point]



Q.3. Solutions [continued]



$$\begin{aligned}
 \Rightarrow \vec{A}(s, t) &= \left(\frac{\mu_0 I_0}{4\pi} \hat{z} \right) 2 \int_0^{\sqrt{(ct)^2 - s^2}} \frac{dz}{\sqrt{s^2 + z^2}} \\
 &= \left(\frac{\mu_0 I_0}{2\pi} \hat{z} \right) \left[\ln(\sqrt{s^2 + z^2} + z) \right]_0^{\sqrt{(ct)^2 - s^2}} \\
 &= \left(\frac{\mu_0 I_0}{2\pi} \hat{z} \right) \ln \left(\frac{ct + \sqrt{(ct)^2 - s^2}}{s} \right)
 \end{aligned}$$

The derivatives of $\vec{A}$ ($= A\hat{z}$) give the fields:

$$\vec{E}(s, t) = -\frac{\partial \vec{A}}{\partial t} = -\frac{\mu_0 I_0 c}{2\pi \sqrt{(ct)^2 - s^2}} \hat{z}$$

$$\vec{B}(s, t) = \nabla \times \vec{A} = -\frac{\partial A_z}{\partial s} \hat{\phi} = \frac{\mu_0 I_0}{2\pi s} \frac{ct}{\sqrt{(ct)^2 - s^2}} \hat{\phi}$$

Note that as $t \rightarrow \infty$, $\vec{E} \rightarrow 0$, $\vec{B} \rightarrow \frac{\mu_0 I_0}{2\pi s} \hat{\phi}$ (the static fields)

Q.4. Poynting Vector

[*Feynman Lectures in Physics*, Vol. II, 27.5]

A capacitor with circular plates, radius R , separation d , is being charged. Find the rate at which its stored energy is increasing and relate this to the Poynting vector flux.

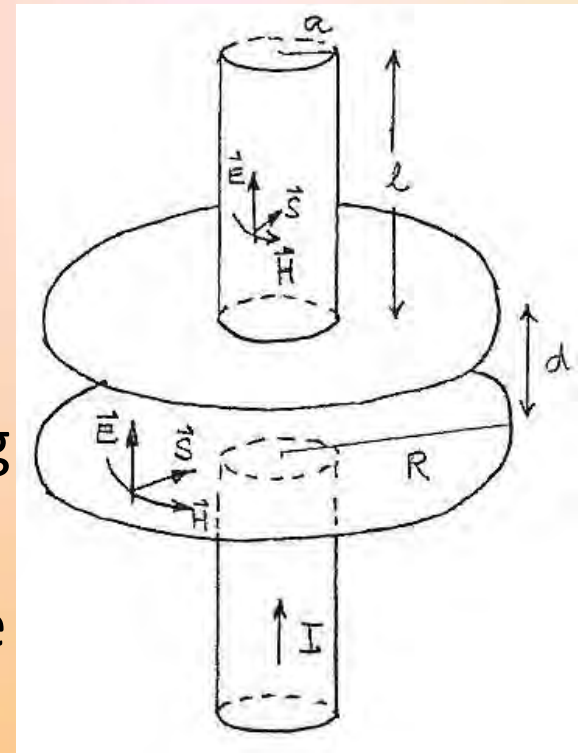
Solution: We assume the field between the plates is uniform. The energy stored is

$$U = \int_V \frac{1}{2} \epsilon_0 E^2 dV = \frac{1}{2} \epsilon_0 E^2 \pi R^2 d \text{ and this increases at the rate } \frac{\partial U}{\partial t} = \pi R^2 d \frac{\partial}{\partial t} \left(\frac{1}{2} \epsilon_0 E^2 \right) = \pi R^2 d \epsilon_0 E \frac{\partial E}{\partial t}$$

Now in the capacitor space $\vec{J}_f = 0$ and Maxwell-Ampere $\rightarrow$

$$\oint \vec{H} \cdot d\vec{l} = \int \frac{\partial \vec{D}}{\partial t} \cdot d\vec{a} \text{ i.e. } H \cdot 2\pi R = \epsilon_0 \frac{\partial E}{\partial t} \cdot \pi R^2 \text{ so } H = \frac{R}{2} \epsilon_0 \frac{\partial E}{\partial t}$$

[$\vec{H}$ is azimuthal, field lines around the cylindrical air space.]



Q.4. *Solution* [continued]

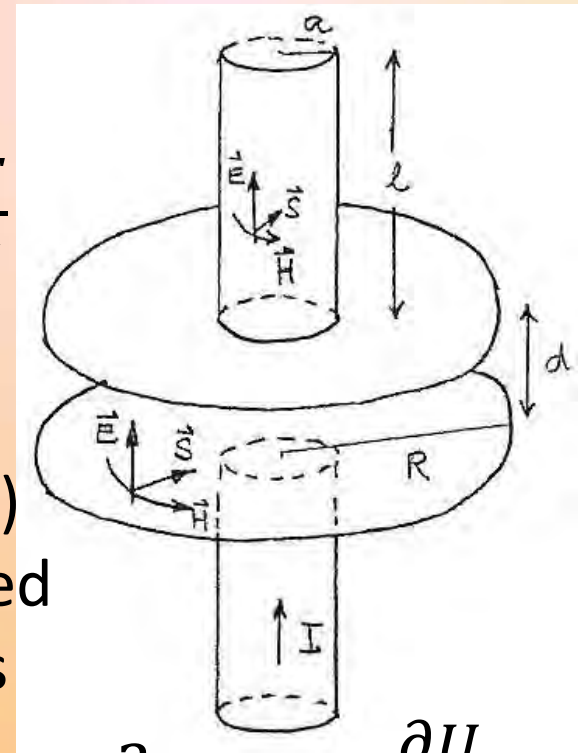
We have $\frac{\partial U}{\partial t} = \pi R^2 d \epsilon_0 E \frac{\partial E}{\partial t}$, $H = \frac{R}{2} \epsilon_0 \frac{\partial E}{\partial t}$

Now the Poynting vector is $\vec{S} = \vec{E} \times \vec{H}$
and we see that it points inwards; this is
the direction of energy flow! (note $\vec{E} \perp \vec{H}$)
The Poynting vector flux, through the curved
surface that defines the capacitor space, is

$$\oint \vec{S} \cdot d\vec{a} = -EH \cdot 2\pi R d = -\epsilon_0 E \frac{\partial E}{\partial t} \pi R^2 d = -\frac{\partial U}{\partial t}$$

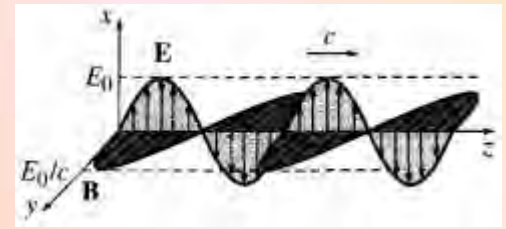
[we have to compute the integral over a closed surface, but
the flux through the ends of the 'cylinder' (the plates) is zero,
since $\vec{S} \parallel$ surface; the integral is negative because flux is **in**.]

This is just **Poynting's theorem** with $\vec{J}_f = 0$: the rate at which
energy flows in = the rate at which the (electric in this case)
field energy is increasing. Energy flows in from outside!



Q.5. EM Waves

[Griffiths, Problem 9.9]



Write down the (real) electric and magnetic fields for a monochromatic plane wave of amplitude E_0 , frequency ω , (and phase angle zero) that is: (a) travelling in the $-x$ -direction and polarized in the $+z$ -direction; (b) travelling in the direction from the origin to the point $(1,1,1)$, with polarization parallel to the xz -plane.

In each case, sketch the wave, and give the explicit Cartesian components of $\vec{k}$ and $\hat{n}$.

Solution: (a) The frequency is specified so we should give k in terms of ω ; the wave direction is $-\hat{x}$, so $\vec{k} = -(\omega/c) \hat{x}$

Then $\vec{k} \cdot \vec{r} = -(\omega/c) \hat{x} \cdot (x\hat{x} + y\hat{y} + z\hat{z}) = -(\omega/c) x$

$\vec{E}$ is in the given direction of polarization, $\hat{z}$

$\vec{B}$ is in the direction $\hat{k} \times \hat{n} = (-\hat{x}) \times \hat{z} = \hat{y}$

Q.5. *Solution [continued]*

Thus the fields are

$$\vec{E}(x, t) = E_0 \hat{z} \cos [(\omega/c)x + \omega t]$$

$$\vec{B}(x, t) = (E_0/c) \hat{y} \cos [(\omega/c)x + \omega t]$$

(b) The unit vector from the origin to the point (1,1,1) is $(\hat{x} + \hat{y} + \hat{z})/\sqrt{3}$, so

$$\vec{k} = (\omega/c) (\hat{x} + \hat{y} + \hat{z})/\sqrt{3}$$

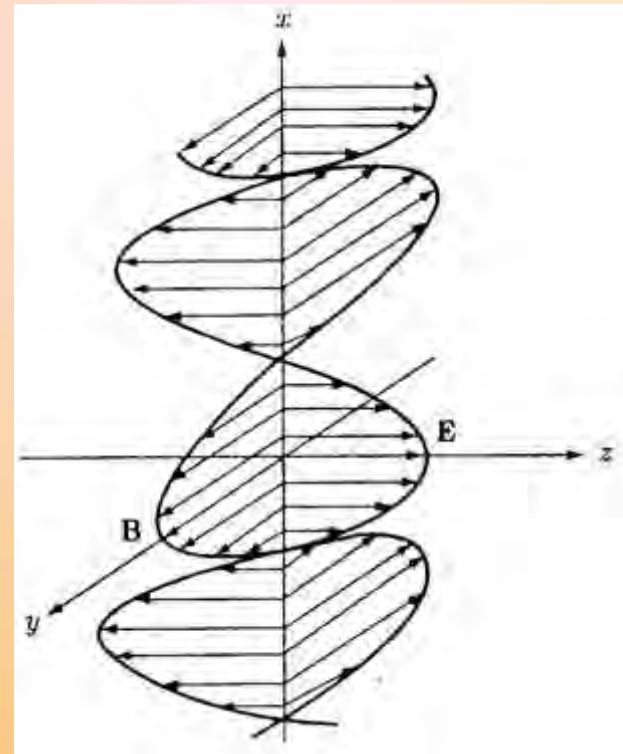
Polarization direction $\hat{n}$ is $\parallel$ xz -plane, so

must have the form $\alpha \hat{x} + \beta \hat{z}$; $\hat{n} \cdot \vec{k} = 0$, so must have $\alpha = -\beta$;

$\hat{n}$ is a unit vector, so $\alpha = 1/\sqrt{2}$ and thus $\hat{n} = (\hat{x} - \hat{z})/\sqrt{2}$

This is the direction of $\vec{E}$; the direction of $\vec{B}$ is $\vec{k} \times \hat{n}$:

$$\vec{k} \times \hat{n} = \frac{1}{\sqrt{6}} \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ 1 & 1 & 1 \\ 1 & 0 & -1 \end{vmatrix} = \frac{(-\hat{x} + 2\hat{y} - \hat{z})}{\sqrt{6}}$$



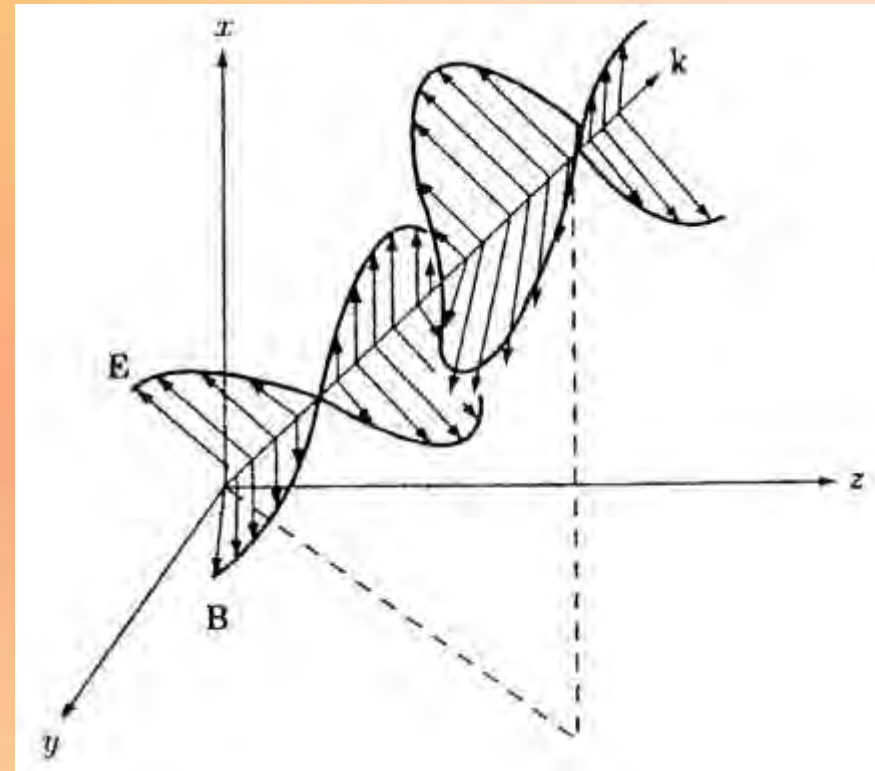
Q.5. *Solution [continued]*

Thus the fields are

$$\vec{E}(x, y, z, t) = E_0 \cos \left[\frac{\omega(x + y + z)}{\sqrt{3}c} - \omega t \right] \left(\frac{\hat{x} - \hat{z}}{\sqrt{2}} \right)$$
$$\vec{B}(x, t) = \frac{E_0}{c} \cos \left[\frac{\omega(x + y + z)}{\sqrt{3}c} - \omega t \right] \left(\frac{(-\hat{x} + 2\hat{y} - \hat{z})}{\sqrt{6}} \right)$$

This is what this wave would look like; sometimes this might be quite difficult to draw!

[Part (b) is far more intricate than any question you would get in a test or exam.]



Q.6. Energy in EM Waves

The “solar constant”, the average intensity of radiation from the Sun at the top of the Earth’s atmosphere, is 1.34 kW/m^2 . If this radiation was a linearly polarized monochromatic (single frequency) plane wave (propagating in free space, of course), find the amplitudes of $\vec{E}$ and $\vec{B}$ for this wave.

Solution: Intensity (in W/m^2) is the (time averaged) Poynting vector: $I = \langle S \rangle = \frac{1}{2} c \epsilon_0 E_0^2$ Inserting the values,

$$E_0 = \sqrt{\frac{2I}{c\epsilon_0}} = \sqrt{\frac{2 \times 1.34 \times 10^3}{3.0 \times 10^8 \times 8.85 \times 10^{-12}}} = 1.0 \times 10^3 \text{ V/m}$$

$E_0 = 1 \text{ kV/m}$! The magnetic field is

$$B_0 = \frac{E_0}{c} = \frac{1.0 \times 10^3}{3.0 \times 10^8} = 3.3 \times 10^{-6} \text{ T} = 3.3 \mu\text{T}$$

