

Tools for measuring galaxy space densities from HI surveys

Martin Zwaan - ESO

HI stacking at $z=1.3$

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THE GAS MASS OF STAR-FORMING GALAXIES AT $Z \approx 1.3$

NISSIM KANEKAR¹, SHIV SETHI², K. S. DWARAKANATH²

Draft version February 1, 2016

ABSTRACT

We report a Giant Metrewave Radio Telescope (GMRT) search for HI 21 cm emission from a large sample of star-forming galaxies at $z \approx 1.18 - 1.34$, lying in sub-fields of the DEEP2 Redshift Survey. The search was carried out by co-adding (“stacking”) the HI 21 cm emission spectra of 857 galaxies, after shifting each galaxy’s HI 21 cm spectrum to its rest frame. We obtain the 3σ upper limit $S_{\text{HI}} < 2.5 \mu\text{Jy}$ on the average HI 21 cm flux density of the 857 galaxies, at a velocity resolution of $\approx 315 \text{ km s}^{-1}$. This yields the 3σ constraint $M_{\text{HI}} < 2.1 \times 10^{10} \times [\Delta V / 315 \text{ km/s}]^{1/2} M_{\odot}$ on the average HI mass of the 857 stacked galaxies, the first direct constraint on the atomic gas mass of galaxies at $z > 1$. The implied limit on the average atomic gas mass fraction (relative to stars) is $M_{\text{GAS}}/M_{*} < 0.5$, comparable to the cold molecular gas mass fraction in similar star-forming galaxies at these redshifts. We find that the cosmological mass density of neutral atomic gas in star-forming galaxies at $z \approx 1.3$ is $\Omega_{\text{GAS}} < 3.7 \times 10^{-4}$, significantly lower than Ω_{GAS} estimates in both galaxies in the local Universe and damped Lyman- α absorbers at $z \geq 2.2$. Blue star-forming galaxies thus do not appear to dominate the neutral atomic gas content of the Universe at $z \approx 1.3$.

Subject headings: galaxies: high-redshift — galaxies: ISM — galaxies: star formation — radio lines: galaxies

1. INTRODUCTION

Over the last decade, optical and infrared studies of the so-called “deep fields” (e.g. [Giavalisco et al. 2004](#); [Scoville et al. 2007](#); [Newman et al. 2013](#)) have revolutionized our understanding of the early Universe. Detailed stud-

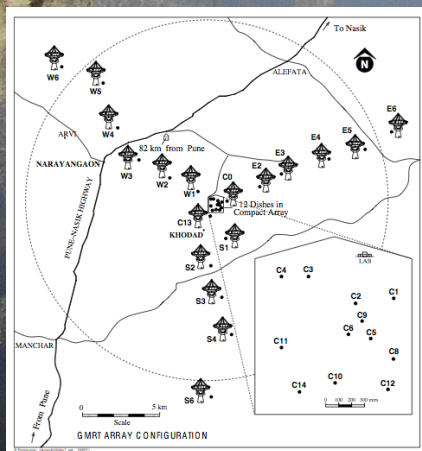
of the molecular component, studies of CO in emission at $z \approx 1.5 - 3$ have indeed found evidence for massive reservoirs of molecular gas ($M_{\text{H}_2} \gtrsim 10^{10} M_{\odot}$, comparable to the stellar mass) in star-forming galaxies (e.g. [Daddi et al. 2010](#); [Tacconi et al. 2010, 2013](#)). Unfortunately, CO is only a tracer of the molecular gas, and the conversion factor

Biggs, Ivison, Zwaan (BIZ)
HI stacking at $z=1.3$

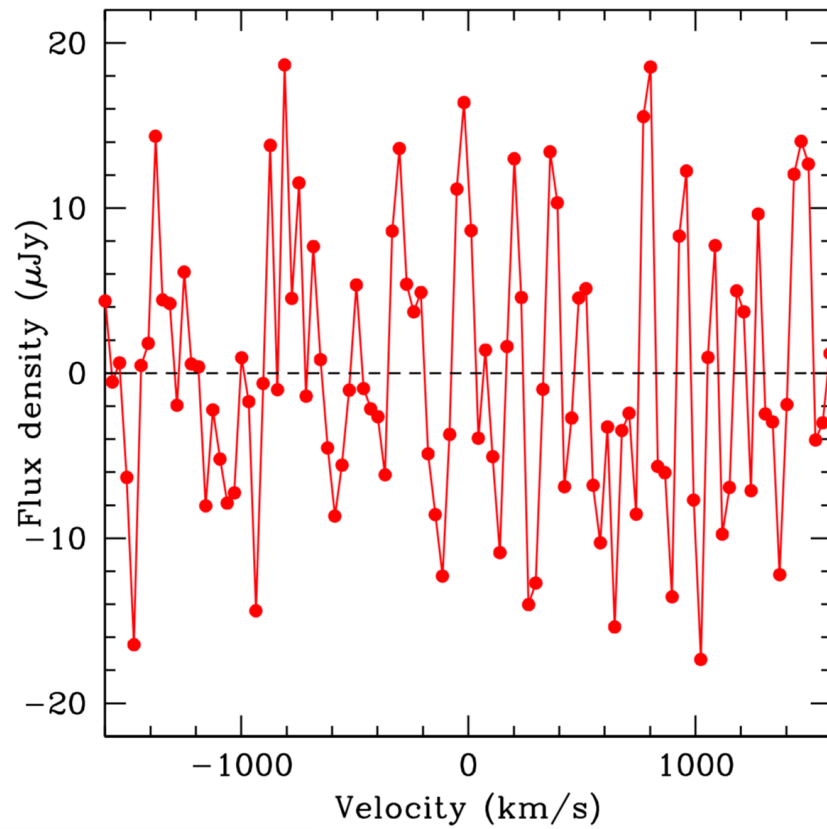
The GMRT

610 MHz corresponds to HI at a redshift of ≈ 1.3
Angular resolution = 4-5"
Primary beam = 44'

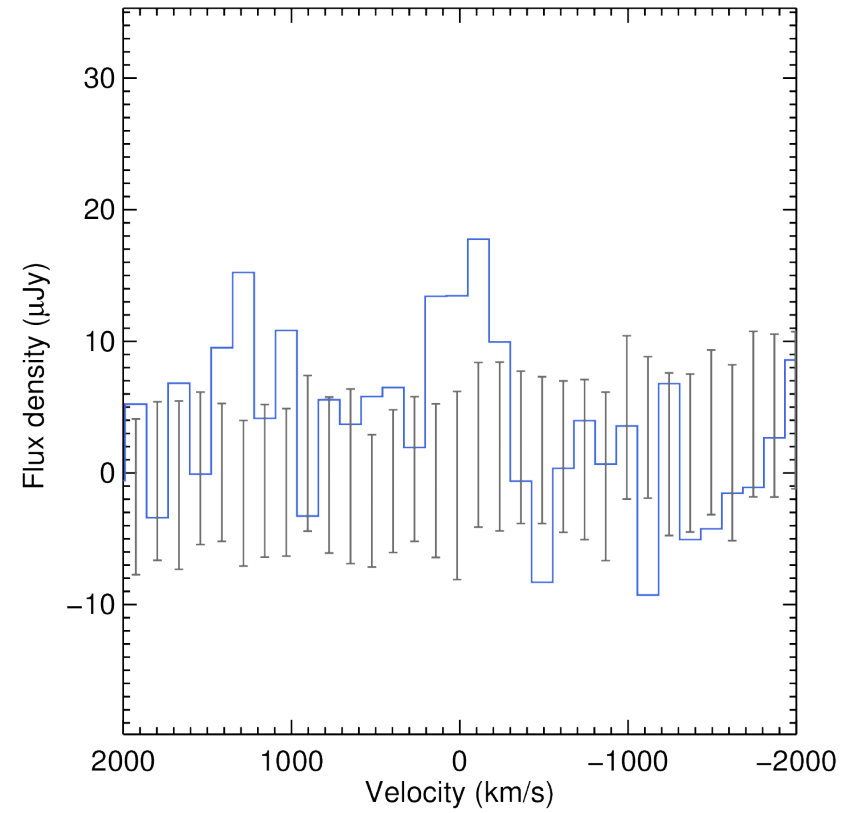
30 45-m dishes
Longest baseline = 25 km



Compare stacked HI spectrum of Kanekar and BIZ

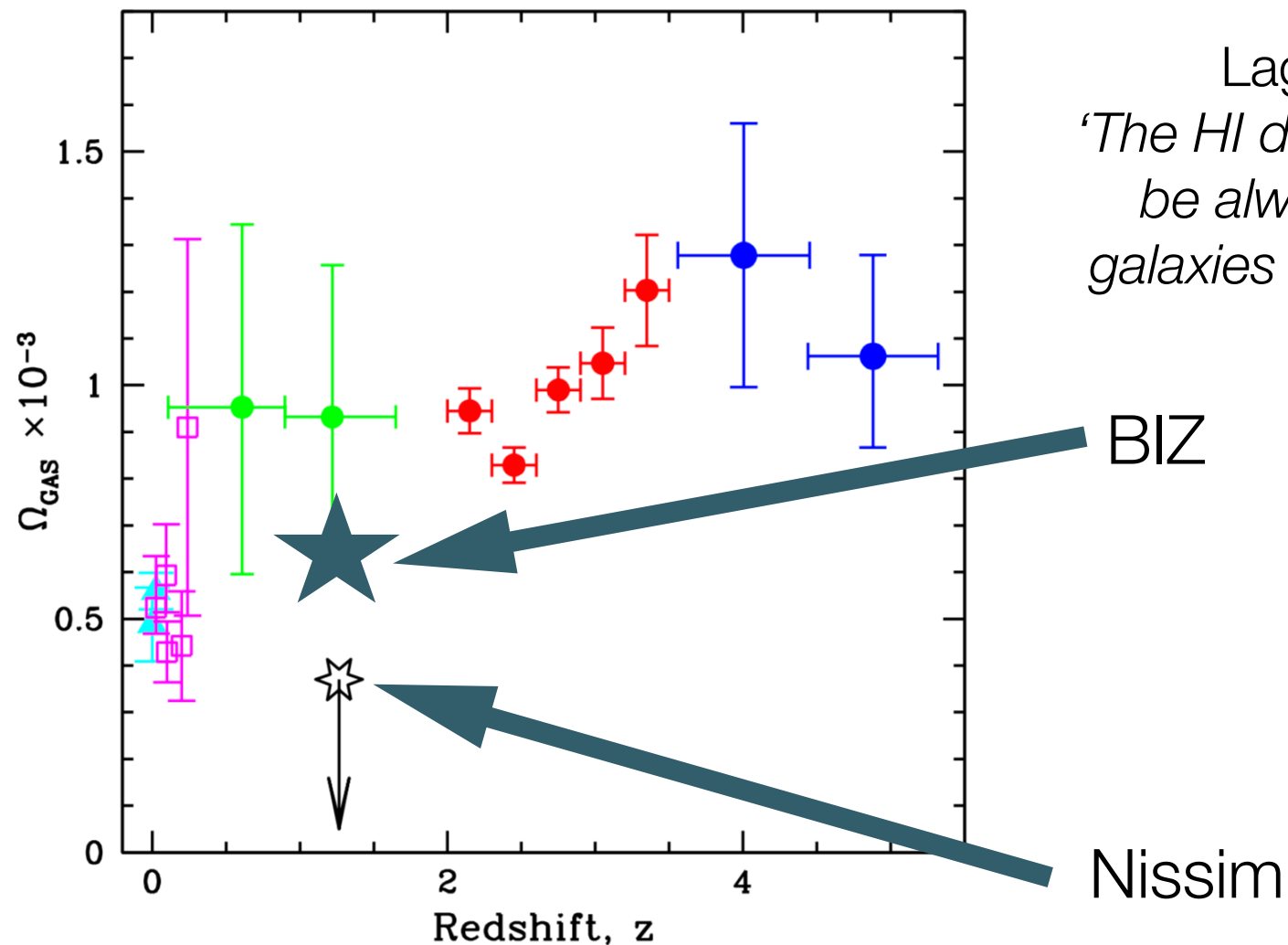


Kanekar et al



BIZ

HI mass density



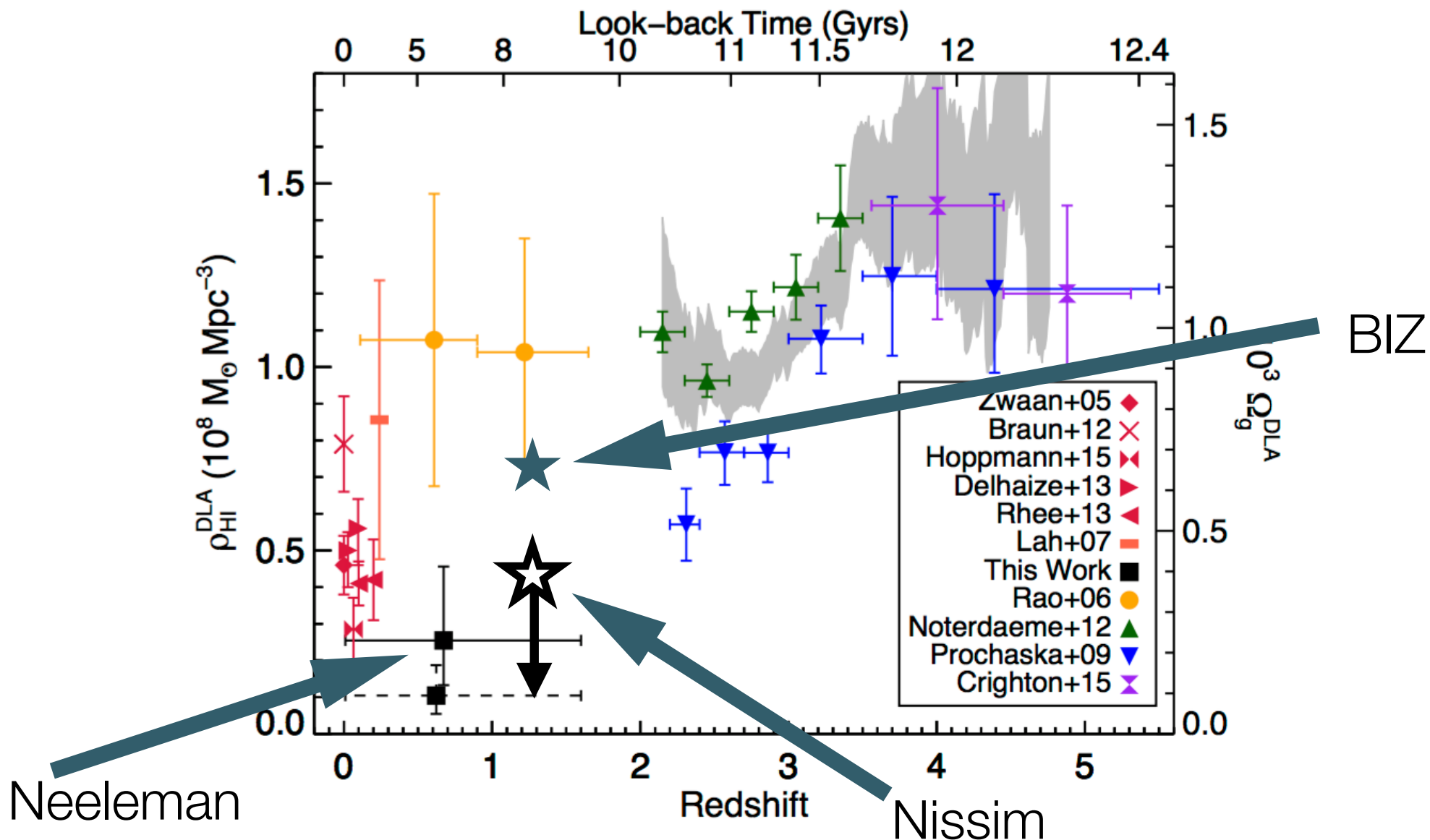
Lagos et al 2014:
*'The HI density is predicted to
be always dominated by
galaxies with $SFR < 1 M_{\odot} \text{ yr}^{-1}$ '*

BIZ

Nissim

Latest DLA results... (Neeleman et al. 2016)

Four new DLAs at $z < 1.6$



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THE PAPER MOUNTAIN

If you were to print out just the first page of every item indexed in Web of Science, the stack of paper would reach almost to the top of Mt Kilimanjaro. Only the top metre and a half of that stack would have received 1,000 citations or more, and just a centimetre and a half would have been cited more than 10,000 times. All of the top 100 are cited more than 12,000 times, besting some of the most recognizable scientific discoveries in history.

TOP-100 PAPERS

10,000+ CITATIONS
(148 papers)

TOP-10 PAPERS

Just 3 papers have received more than 100,000 citations, putting them well ahead of the rest. These runaway hits all cover biological lab techniques, which in general dominate the list of most-cited literature, including 7 of the top 10.

- 305,148 citations
Protein measurement with the folin phenol reagent (1951)
- 213,005
Cleavage of structural proteins during the assembly of head of the bacteriophage T4 (1970)
- 155,530
A rapid and sensitive method for the quantitation of microgram quantities of protein utilizing the principle of protein-dye binding (1976)
- 65,335
DNA sequencing with chain-terminating inhibitors (1977)
- 60,397
Single-step method of RNA isolation by acid guanidinium thiocyanate-phenol-chloroform extraction (1987)
- 53,349
Electrophoretic transfer of proteins from polyacrylamide gels to nitrocellulose sheets: procedure and some applications (1979)
- 46,702
Development of the Colle-Salvetti correlation-energy formula into a functional of the electron density (1988)
- 46,145
Density-functional thermochemistry. III. The role of exact exchange (1993)
- 45,131
A simple method for the isolation and purification of total lipides from animal tissues (1957)
- 40,289
CLUSTAL W: improving the sensitivity of progressive multiple sequence alignment through sequence weighting, position-specific gap penalties and weight matrix choice (1994)

Data provided by Thomson Reuters/Web of Science. Individual paper citation figures extracted 7 October 2014. Distribution of citations in database: 19 September 2014

Top-cited papers present **methods**, not scientific results

source: van Noorden et al,
Nature, 2014

Methods for measuring space densities of H I selected galaxies

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Accepted ... Received ...

ABSTRACT

Key words:

1 INTRODUCTION

2 THE METHODS

2.1 The classical $1/V_{\max}$ method

2.2 The 2DSWML method

On galaxy samples that are other than flux-limited, other parameters have to be included in the maximum likelihood calculation of the space density of objects. For a given H I source at distance D , the peak flux is directly proportional to its H I mass M_{HI} , and inversely proportional to the velocity width over which the flux is distributed. Therefore, the velocity width is the second parameter we have to include in the maximum likelihood analysis. A 2-dimensional maximum likelihood algorithm can then be employed to find the true galaxy distribution. In the following we use w_{20} , the profile width measured at 20% of peak intensity, as a measurement of the velocity width.

where

$j = 1, \dots, N_M$ and $k = 1, \dots, N_W$,

and N_M and N_W are the number of bins in M and W , respectively, and we define

$M = \log(M_{\text{HI}})$ and $W = \log(w_{20})$.

The logarithm of the likelihood of detecting all galaxies in the sample can now be expressed as

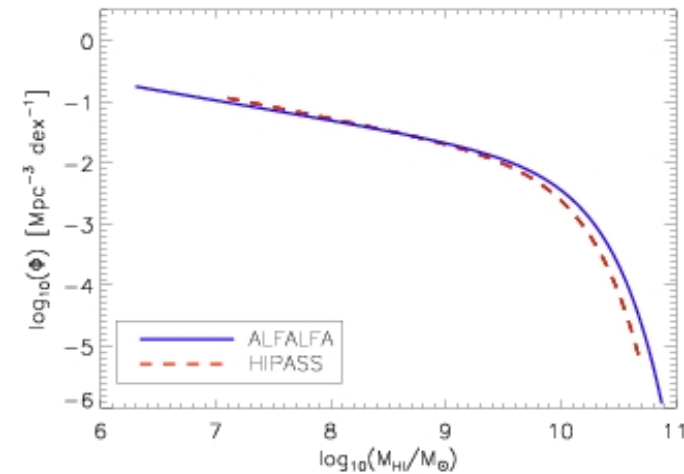
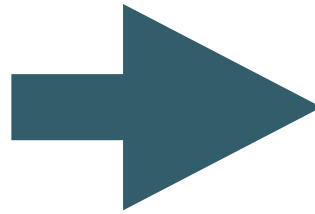
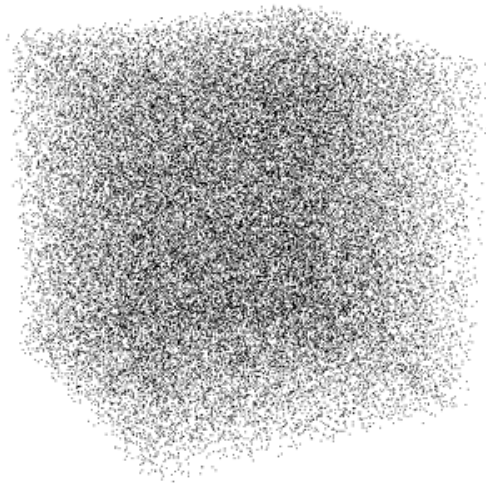
$$\ln \mathcal{L} = \sum_{i=1}^{N_g} \sum_{j=1}^{N_M} \sum_{k=1}^{N_W} V(M_i - M_j, W_i - W_k) \ln \theta_{jk} - \sum_{i=1}^{N_g} \ln \left(\sum_{j=1}^{N_M} \sum_{k=1}^{N_W} H_{ijk} \theta_{jk} \right) + c, \quad (6)$$

where c is a constant and V is a function defined by

$$V(x, y) = \begin{cases} 1, & |x| \leq \Delta M/2 \text{ and } |y| \leq \Delta W/2 \\ 0, & \text{otherwise} \end{cases}, \quad (7)$$

over galaxies in bin $\Delta M \Delta W$.

The problem...



- Complications:

- Complicated completeness limit (S_{peak} , W , profile shape, freq)



- Large scale structure

Isn't this a solved problem?

- Yes, it is for surveys with a well-defined and uni-dimensional selection function

Astron Astrophys Rev (2011)
DOI 10.1007/s00159-011-0041-9

Shedding Light on the Galaxy Luminosity Function

Russell Johnston

Accepted 26 August 2011

Abstract From as early as the 1930s, astronomers have tried to quantify the statistical nature of the evolution and large-scale structure of galaxies by studying their luminosity distribution as a function of redshift - known as the galaxy luminosity function (LF). Accurately constructing the LF remains a popular and yet tricky pursuit in modern observational cosmology where the presence of observational selection effects due to e.g. detection thresholds in magnitude, colour, surface brightness or some combination thereof can render the LF thus introduce bias into the LF.

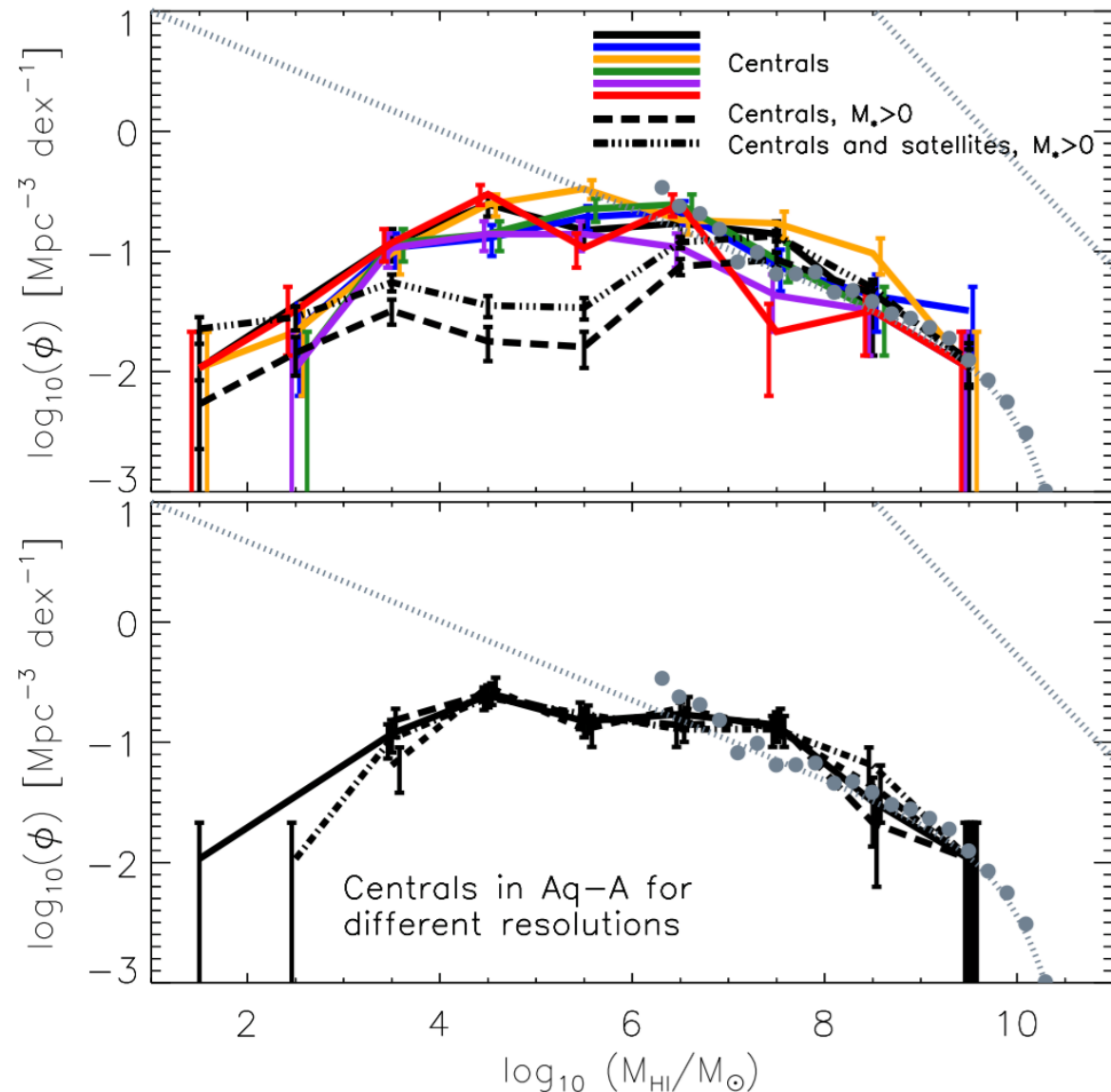
HI mass function in the late 2010s...

- **Faint end**

HI mass function in the late 2010s...

- Faint end**

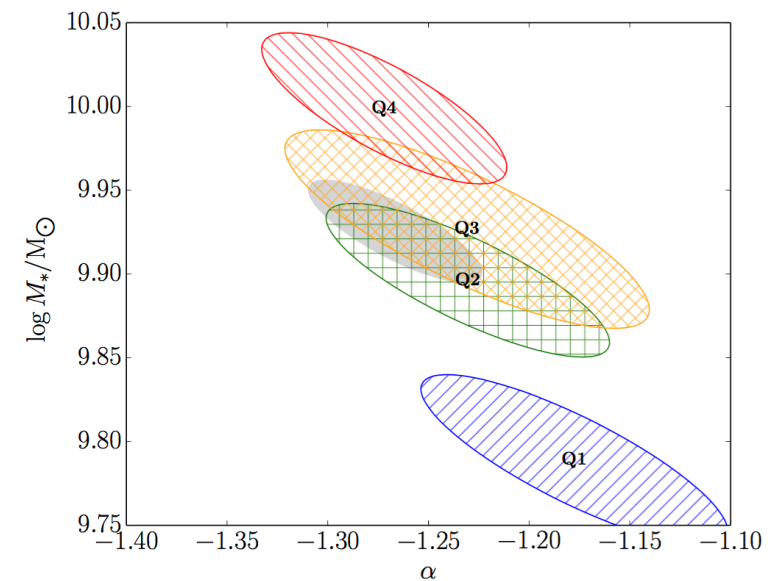
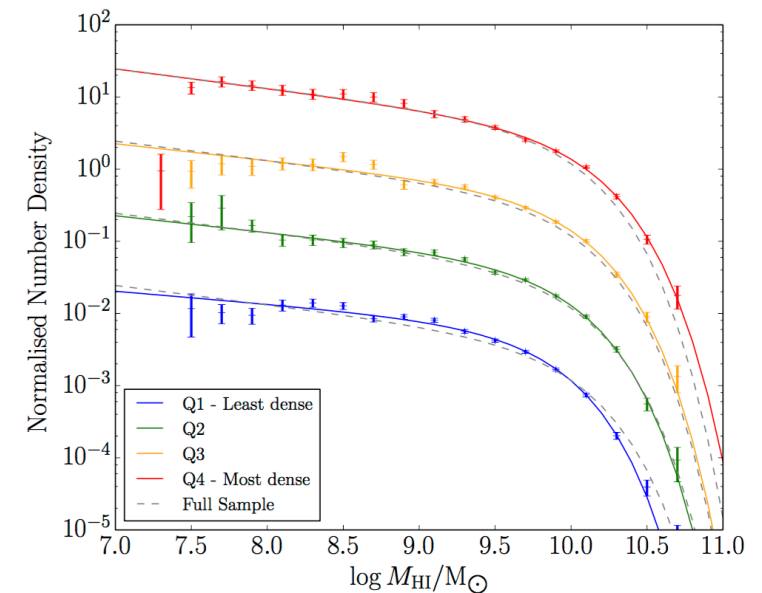
Yaryura et al. 2016:
cosmological dark matter
simulations coupled to the
semi-analytic model,
compared with measured
HI mass function (and
velocity function)



HI mass function in the late 2010s...

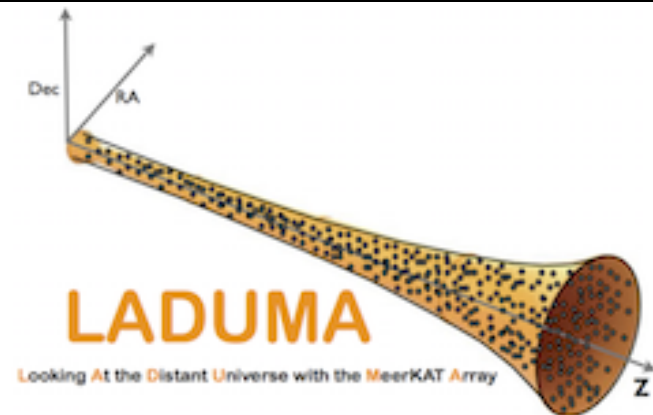
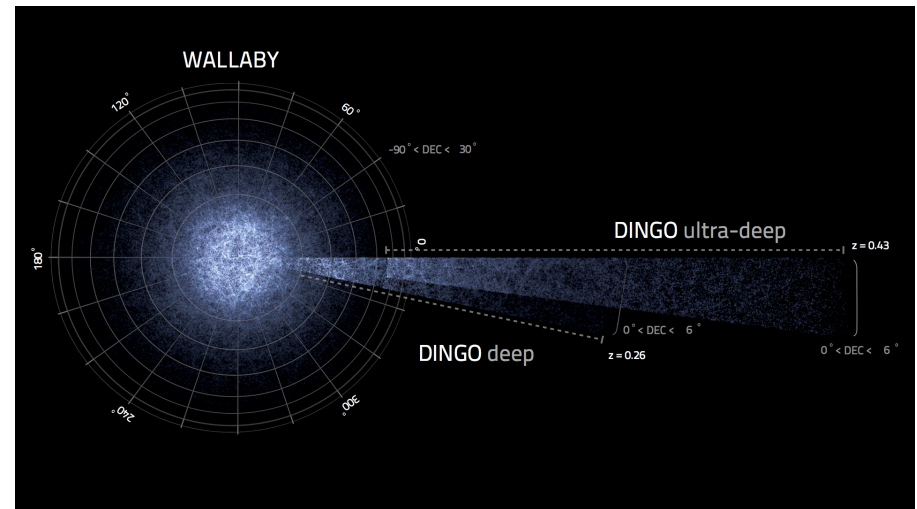
- **Faint end**
- **Environment**

Jones et al. 2016:
HI mass function 'knee' is
dependent on
environment



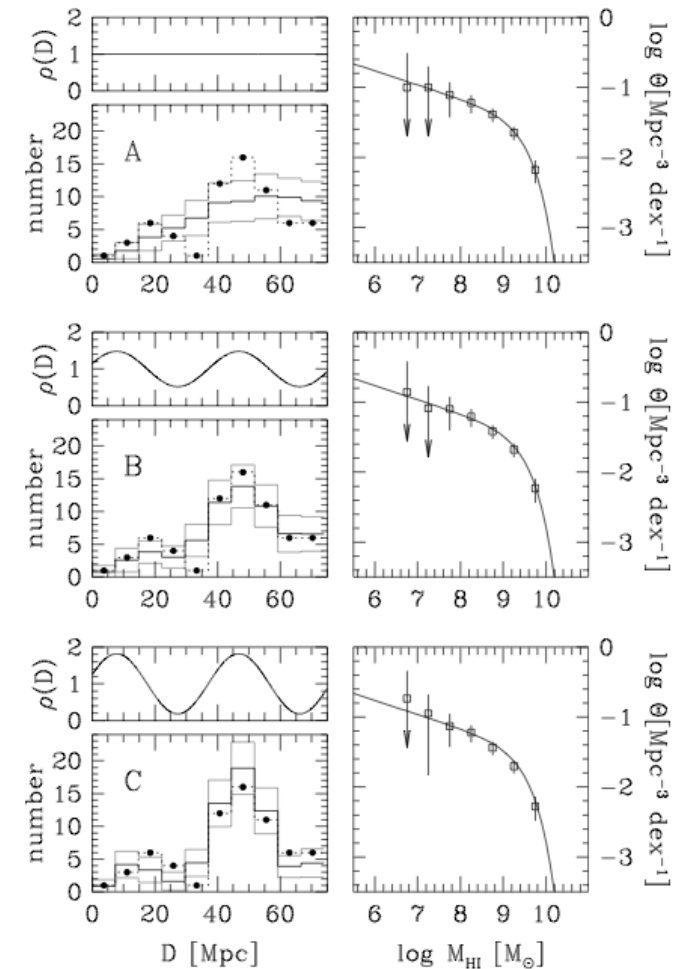
HI mass function in the late 2010s...

- **Faint end**
- **Environment**
- **Evolution**



The $1/V_{\max}$ method

- The ‘classical’ Schmidt (1968) method
- Calculate maximum distance $D_{\max}$ out to which the galaxies can be detected
- Convert $D_{\max}$ into a $V_{\max}$
- Used for early Arecibo surveys
- Advantages: simple and automatically normalised ✓
- Disadvantage: sensitive to large scale structure ✗



Maximum likelihood methods

- Defined by Efstathiou et al 1988, Sandage et al 1979
- Find θ that yields maximal joint probability of detecting all sources in sample

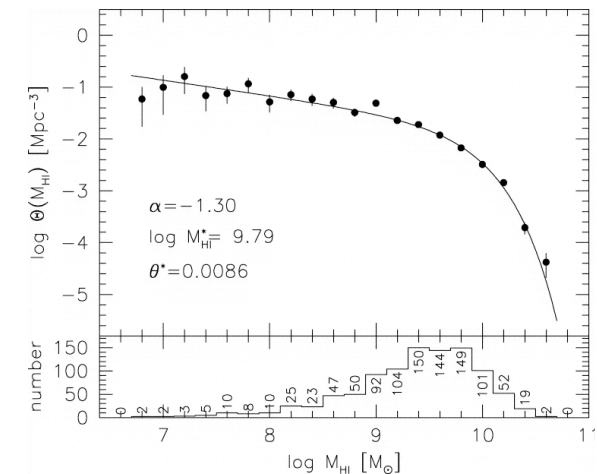
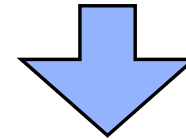
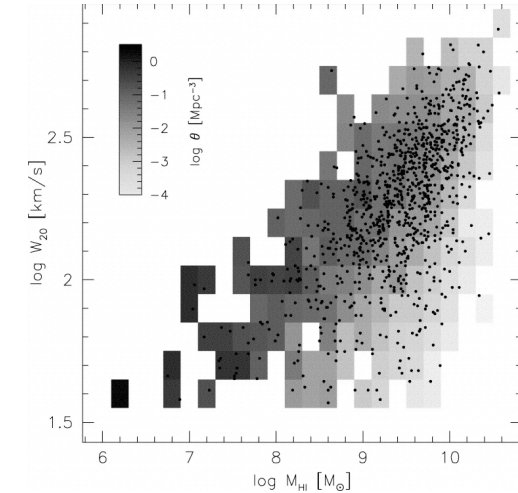
$$p(M_{\text{HI},i}|D_i) = \frac{\theta(M_{\text{HI},i})}{\int_{M_{\text{HI},\min(D_i)}}^{\infty} \theta(M_{\text{HI}}) dM_{\text{HI}}}$$

minimal detectable HI
mass at distance D_i

generally **not defined** for HI
selected samples

2D Stepwise Maximum likelihood method

- Solution: multi-dimensional stepwise maximum likelihood methods
- Find $\theta(M_{\text{HI}}, W)$
- Collapse to find HIMF
- **Used for HIPASS and ALFALFA**
- Advantage: robust against LSS ✓
- Disadvantage: slow ✗



The Turner or φ/Φ -method

- Introduced by Turner (1979) for 3C and 4C quasar catalogues
- Calculate the ratio of number of galaxies in interval dM_{HI} and number of galaxies brighter than M_{HI}

$$\begin{aligned} Y(M'_{\text{HI}})dM_{\text{HI}} &= \frac{N(dM'_{\text{HI}})}{N(\geq M'_{\text{HI}})} \\ &= \frac{\theta(M'_{\text{HI}})\rho(D)dM_{\text{HI}}dV}{\int_{M'_{\text{HI}}}^{\infty} \theta(M_{\text{HI}})\rho(D)dM_{\text{HI}}dV} \end{aligned}$$

$$Y(M'_{\text{HI}})dM_{\text{HI}} = \frac{\theta(M'_{\text{HI}})dM_{\text{HI}}}{\Theta(M'_{\text{HI}})} = \frac{d\Theta(M'_{\text{HI}})}{\Theta(M'_{\text{HI}})}$$

- Advantage: fast and robust against LSS ✓
- Disadvantage: correlated errors ✗

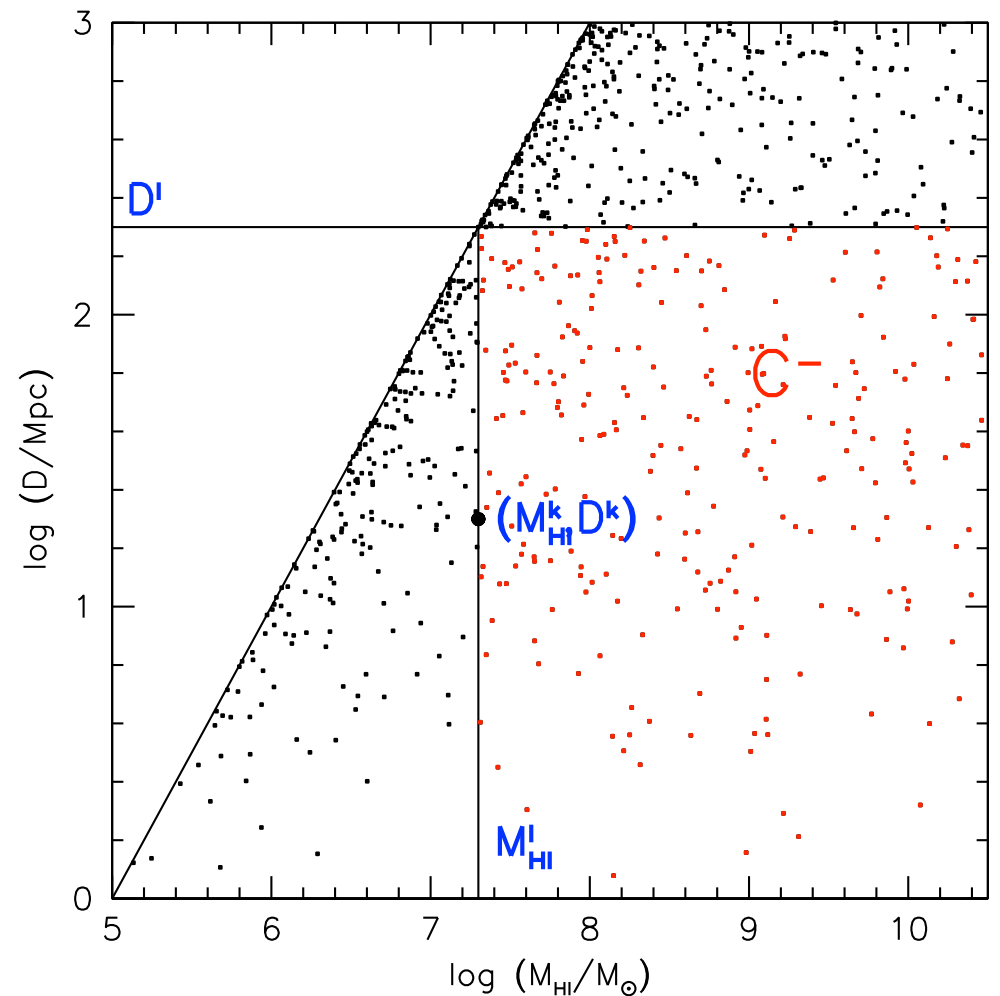
The C⁻ method

- Developed by Lynden-Bell (1971) for quasars.
- Does not require any binning.
- Does not require any assumptions about the distribution of objects within the data-set.
- Estimates the cumulative luminosity function (CLF).

The C⁻ method

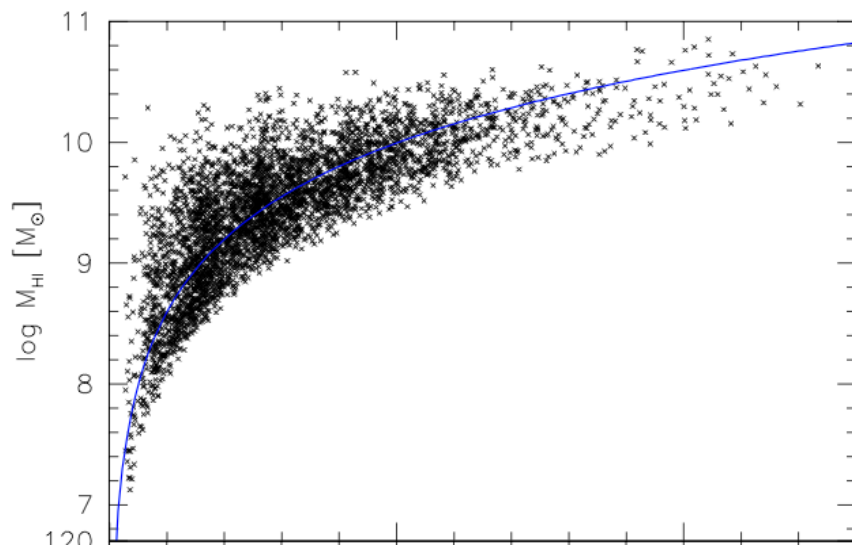
$$\begin{aligned}\Theta(M_{\text{HI}}^i) &= \int_{M_{\text{HI}}}^{\infty} \theta(M_{\text{HI}}) dM_{\text{HI}} \\ &= A \prod_k^{M_{\text{HI}}^k > M_{\text{HI}}^i} \frac{C^-(M_{\text{HI}}^k) + 1}{C^-(M_{\text{HI}}^k)}.\end{aligned}$$

- Then differentiate to obtain $\theta(M_{\text{HI}})$
- Variation of C⁻ is the C⁺ method (Zucca et al. 1998)
- Advantage: independent of clustering effects ✓ and fast ✓



Dealing with gradual drop off in completeness (as opposed to sharp flux limits...)

- All these methods are designed for optical galaxy samples with sharp magnitude limits (m_{lim})
- The 2DSWML, Turner, and C- method are easily adaptable to work for complicated completeness limits



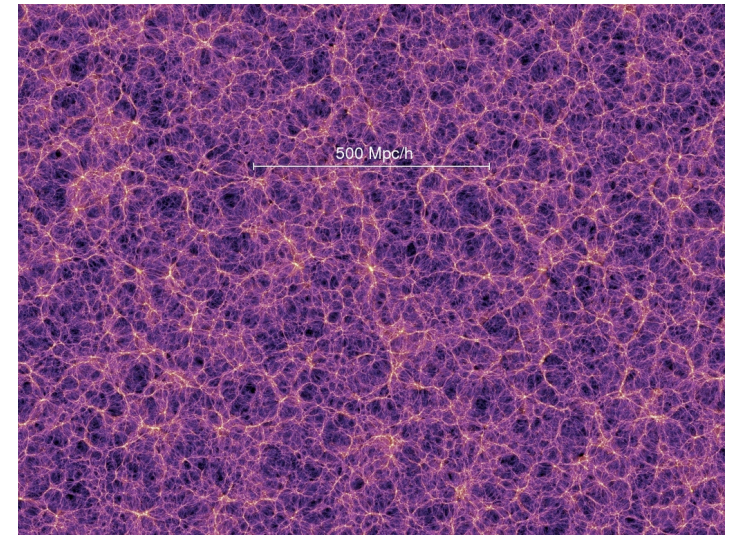
$$C_i^- = C^-(M_{\text{HI}}^i) = \sum_{M_{\text{HI}}^j > M_{\text{HI}}^i} c,$$

Effective volumes

- We always calculate V_{eff}^i per individual galaxy
- These need to be summed per HI mass bin to get the HIMF
- These values can be binned as a function of local density, morphological type, etc.

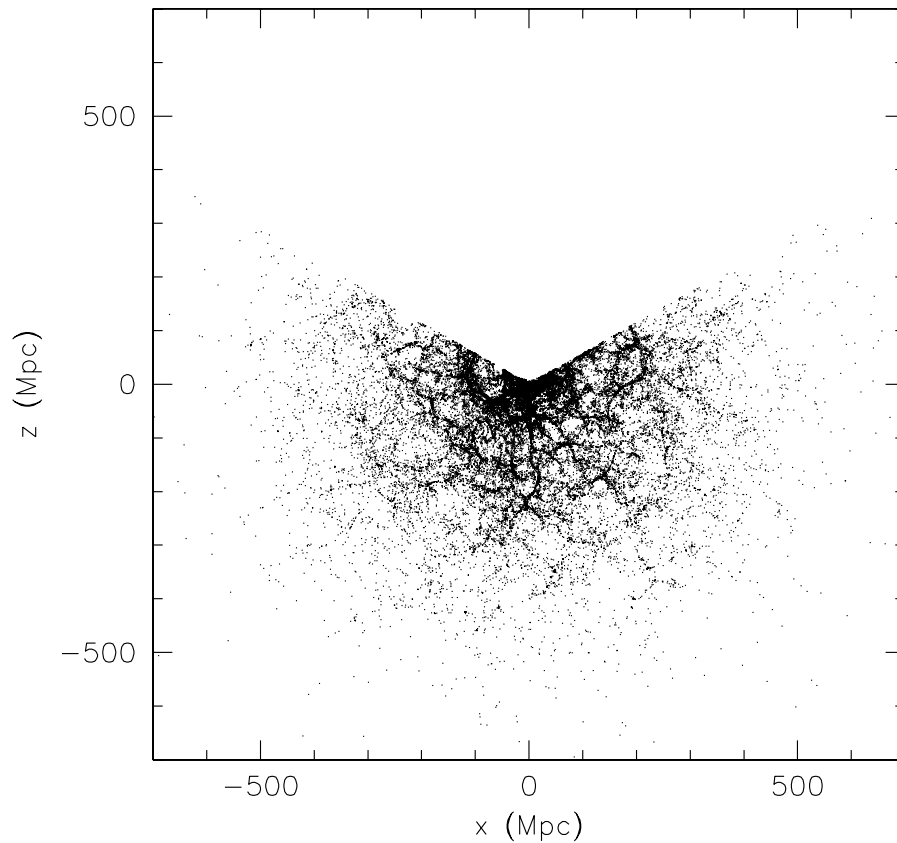
Simulations to test HIMF recovery

- Millennium Simulation (Springel et al 2005)
- Assume a HIPASS HI mass function
- Low mass ($\log M_{\text{HI}} < 8.5$) cluster around larger ones
- Realistic scatter on all parameters
- Select galaxies from simulated boxes, assuming 'optimal smoothing'

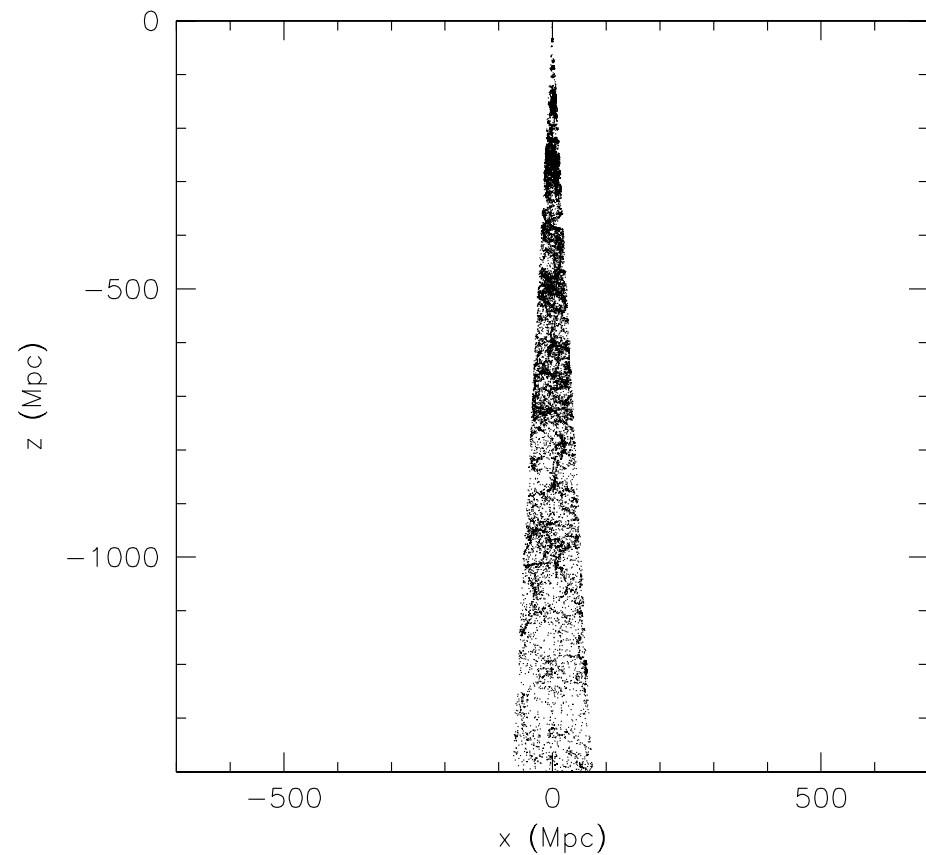


Simulated HI skies

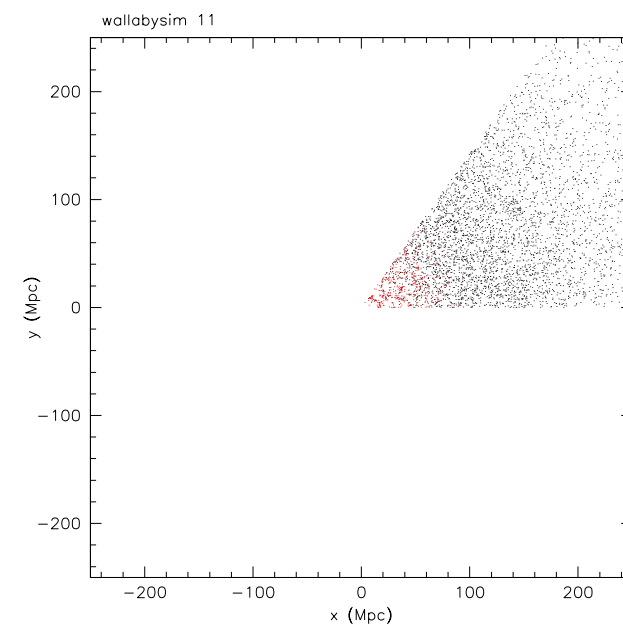
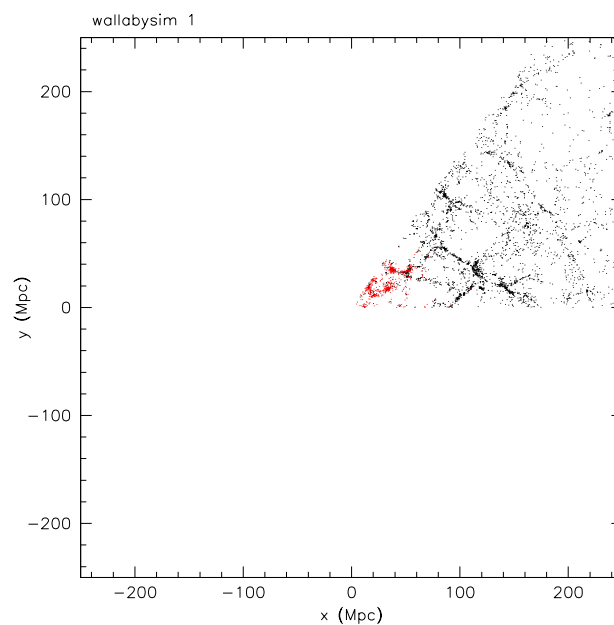
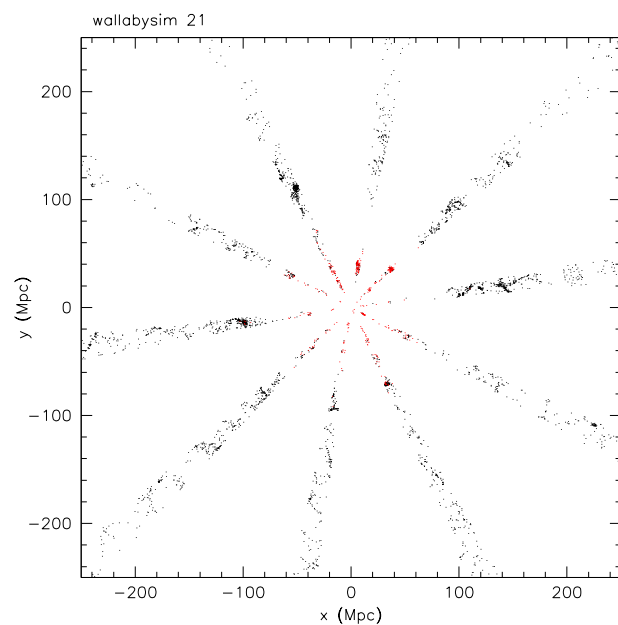
“Wallaby”
all sky shallow



“Dingo”
30deg² deep



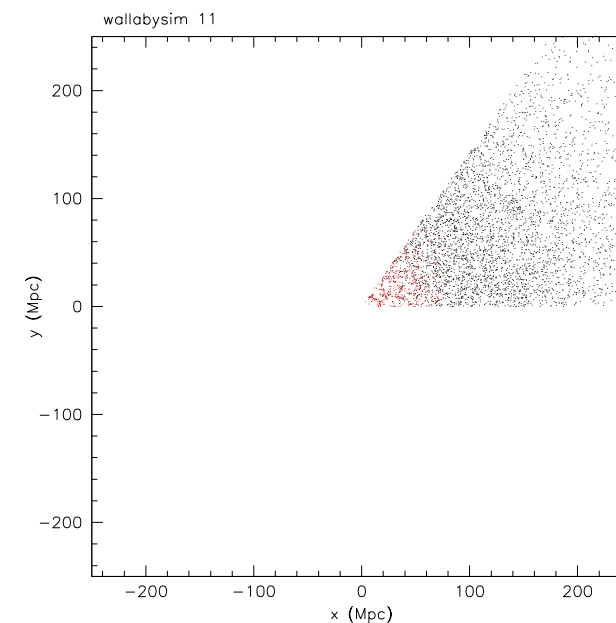
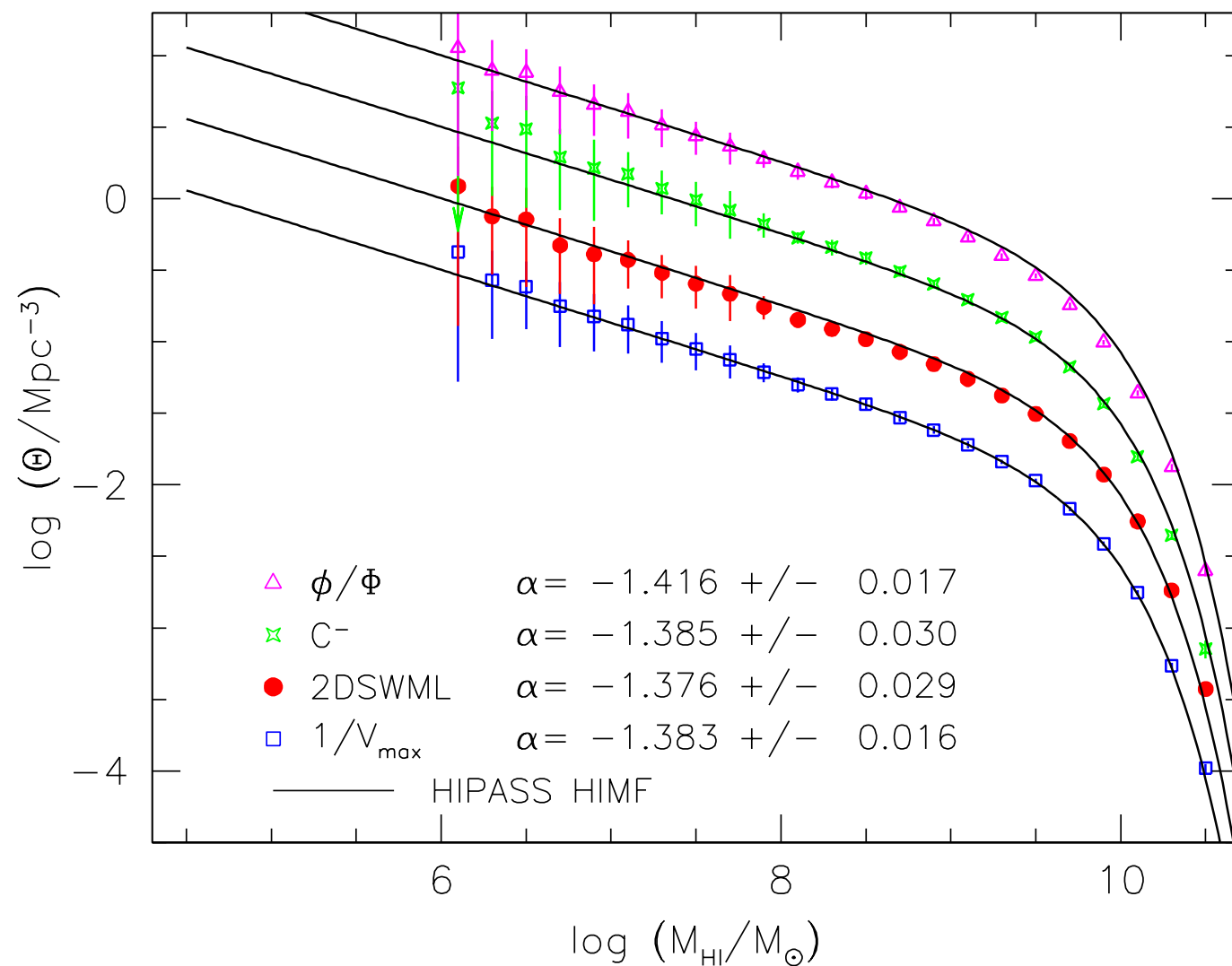
Testing the methods - 10 ASKAP pointings



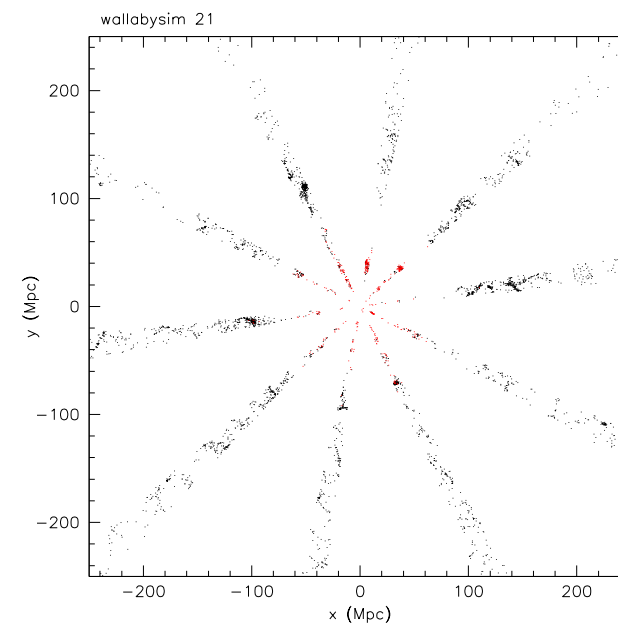
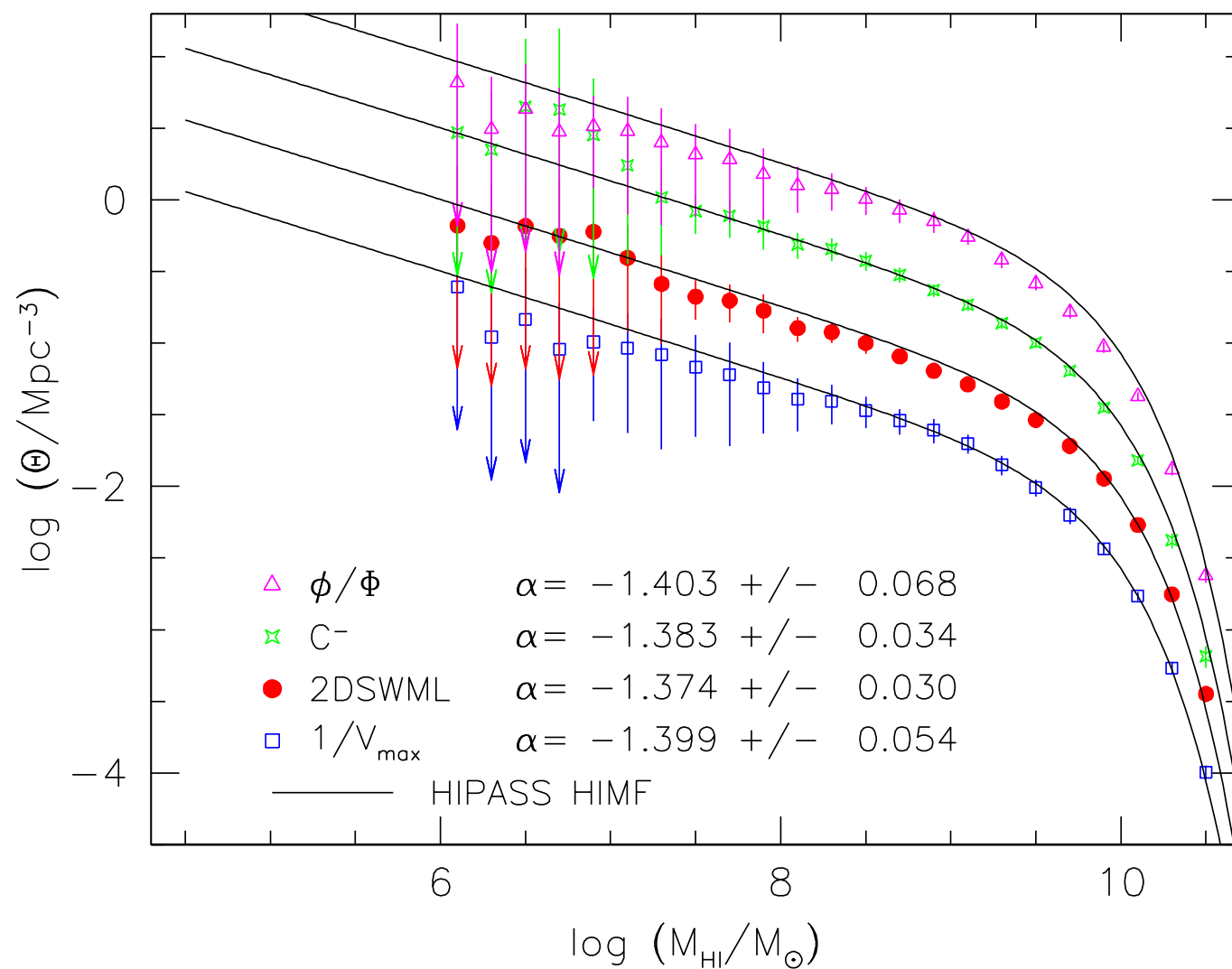
Simulate 10 times

$1/V_{\text{max}}$	ML
φ/Φ	C ⁻

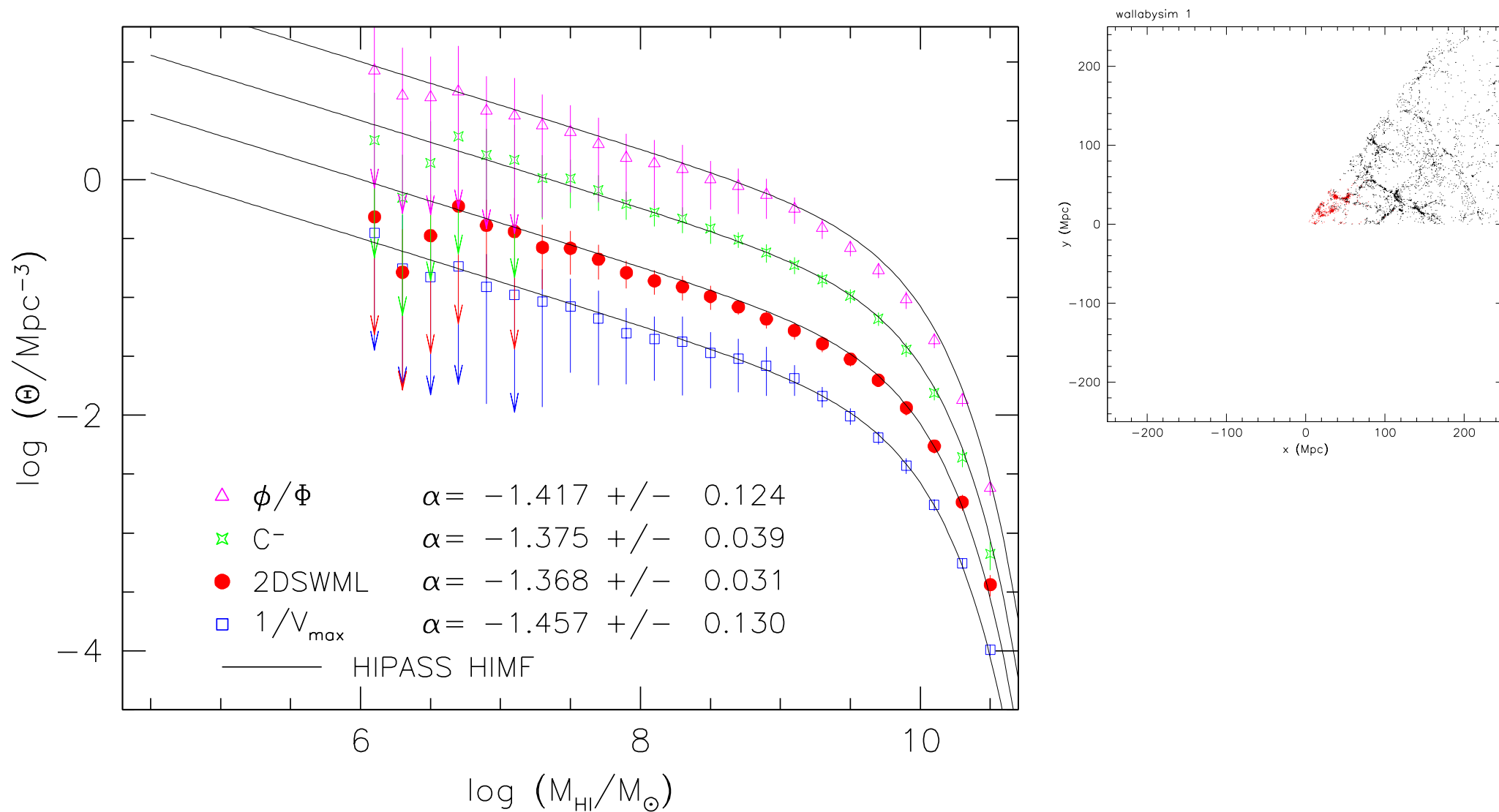
10 ASKAP pointings - No large scale structure



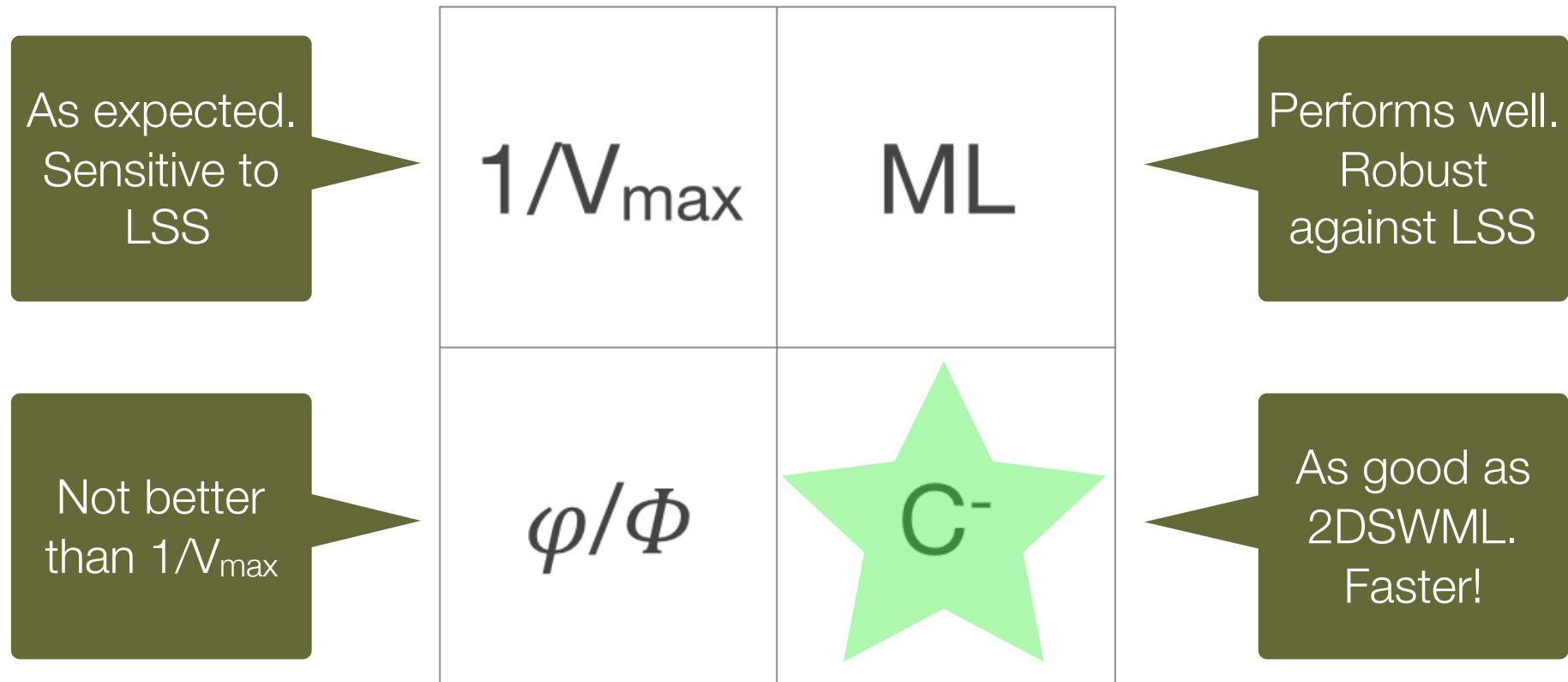
10 ASKAP pointings, with LSS - widely spaced



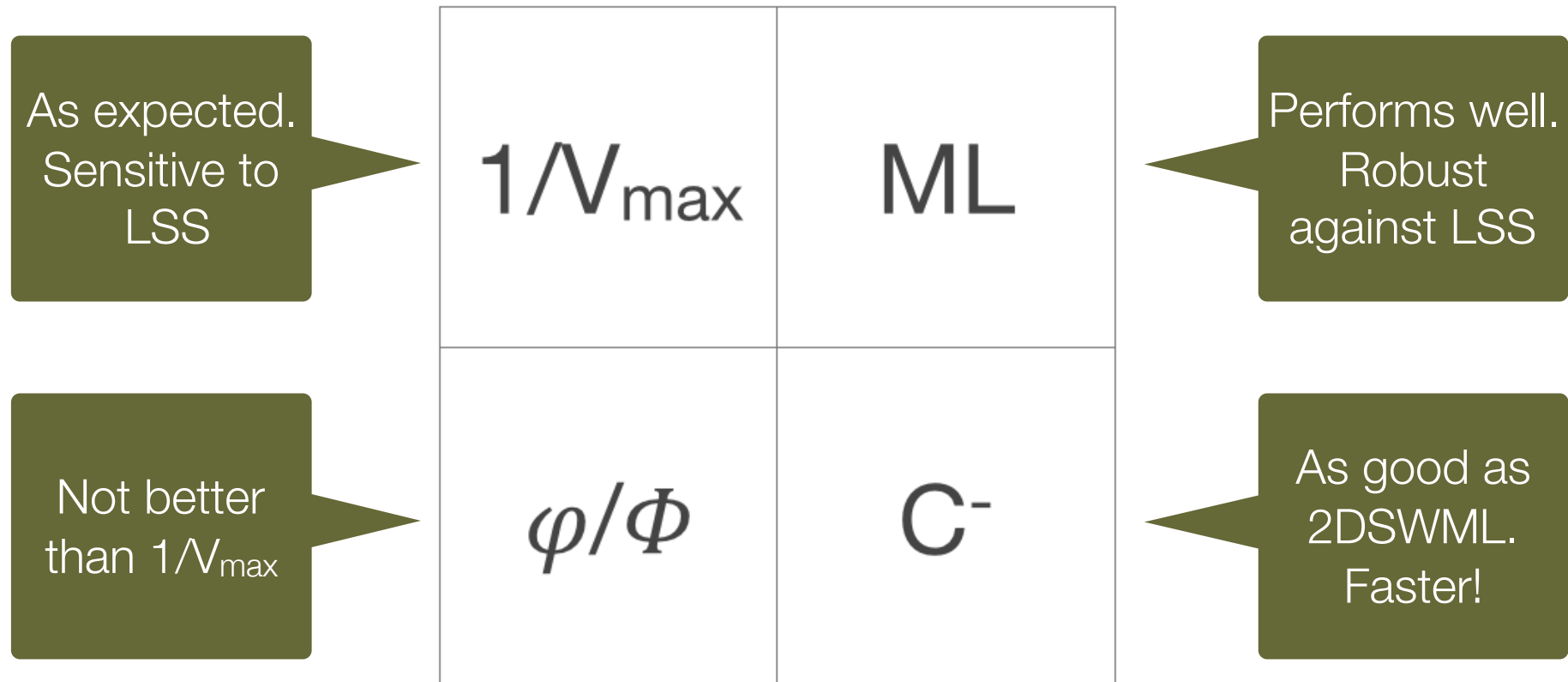
10 ASKAP pointings, with LSS- contiguous



Results of testing methods on shallow HI survey

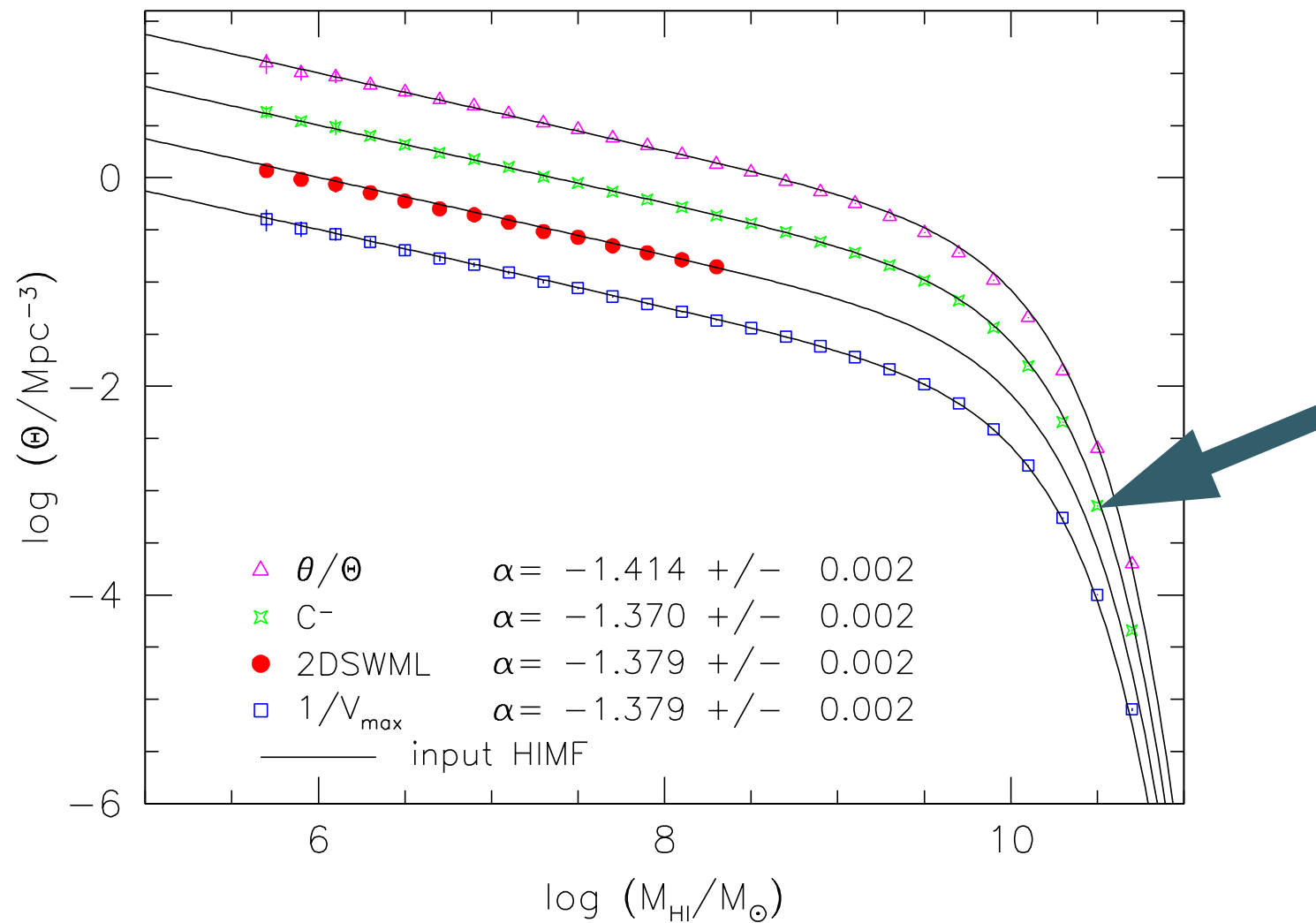


Results of testing methods on shallow HI survey

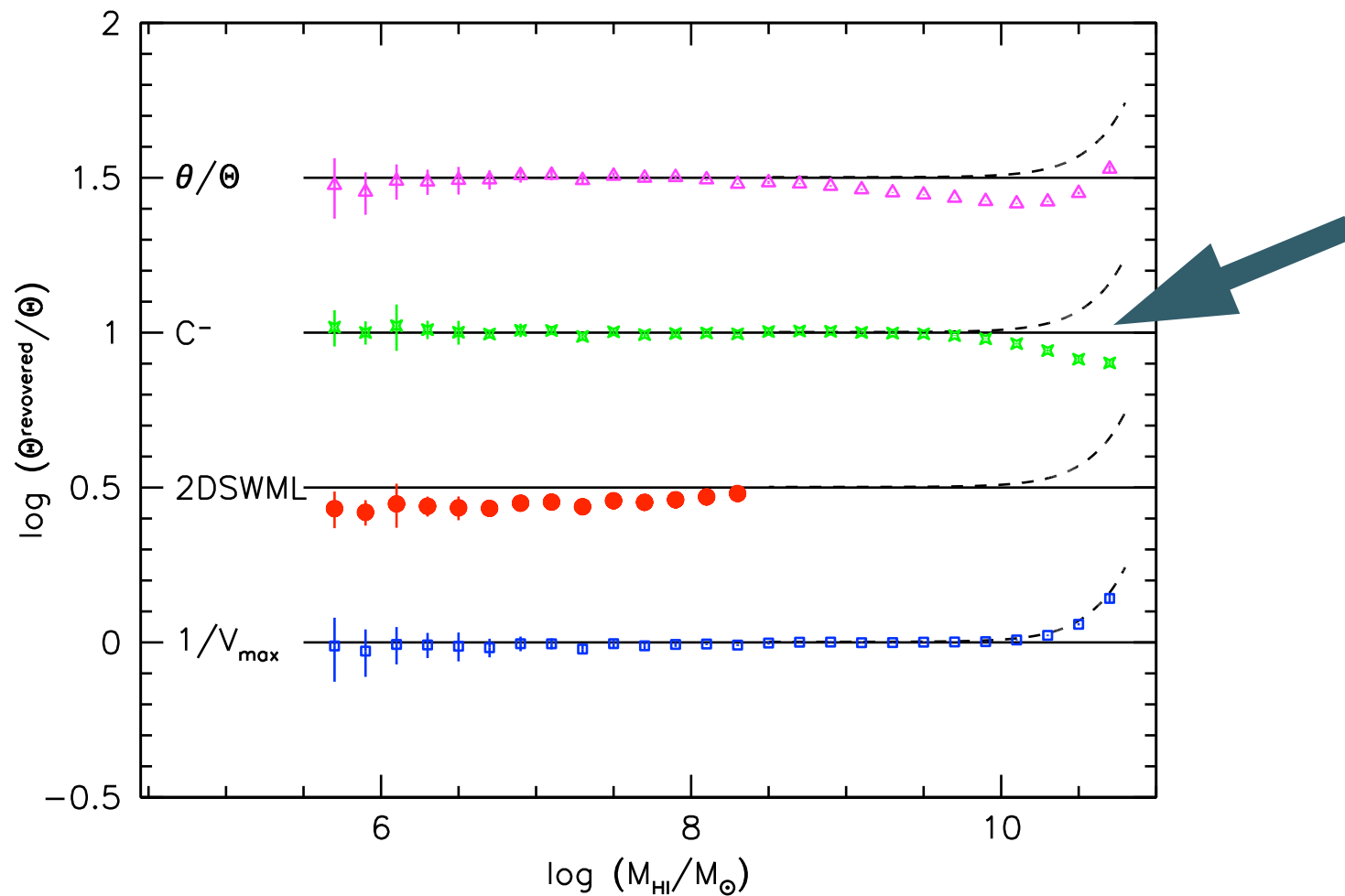


But...

A full Wallaby-type survey, no large scale structure



C⁻ method underestimating space densities at high mass end



Normalising the space densities

- Most methods (apart from $1/V_{\text{max}}$) lose the normalisation of the HIMF
- See Davis & Huchra (1982), Willmer (1997) and Johnston (2011)
- First need selection function:

$$S(D) = \frac{\int_{W_{\text{low}}}^{W_{\text{high}}} \int_{M_{\text{lim}}(D,W)}^{M_{\text{high}}} \theta(M, W) dM dW}{\int_{W_{\text{low}}}^{W_{\text{high}}} \int_{M_{\text{low}}}^{M_{\text{high}}} \theta(M, W) dM dW},$$

- Then normalise: $\int_{M_{\text{low}}}^{M_{\text{high}}} \theta(M) dM = n.$

Normalising the HIMF

- Various methods for recovering the normalisation
 - **n_3** : integral over selection function
 - **n_1** : calculating number of galaxies in redshift shells
 - **$\bar{n}$** : “minimum-variance” - weighting by selection function and second moment of correlation function
 - **counts**: compare real and expected number of galaxies

$$n_3 = \frac{N_T}{\int_0^{z_{max}} s(z) dz}$$

$$n_1 = \frac{\int_0^{z_{max}} \frac{N(z)}{s(z)} dz}{\int_0^{z_{max}} \frac{dV}{dz} dz}.$$

$$n = \frac{\sum_{i=1}^{N_g} N_i(z_i) w(z_i)}{\int_0^{z_{max}} s(z) w(z) \frac{dV}{dz} dz}$$

Normalising the HIMF

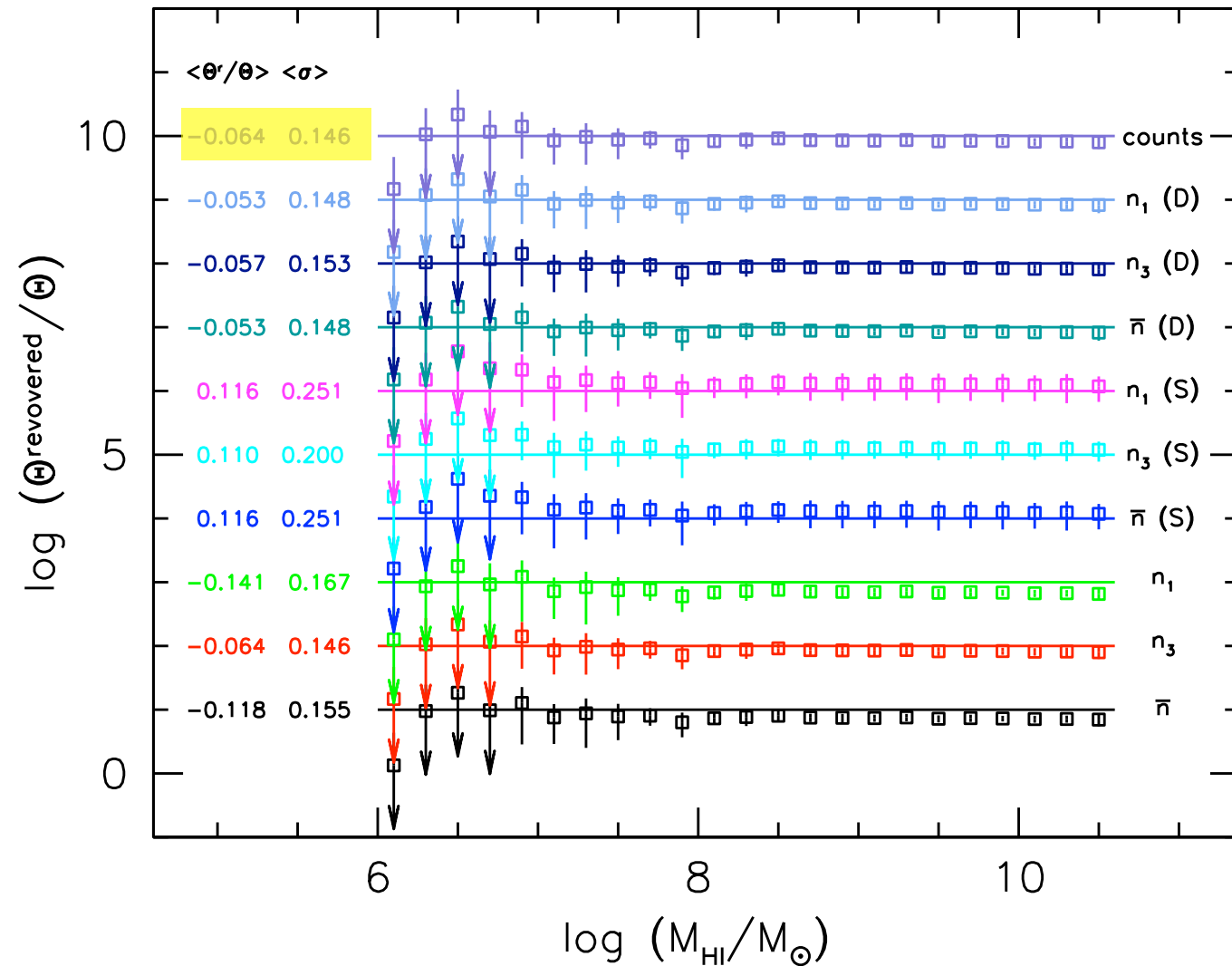
- Uncertainty of normalisation:

$$\sigma_{\bar{n}} = \left\{ \frac{\bar{n}}{\int dV S(D)w} \right\}^{1/2}$$

- For example, for HIPASS, the relative uncertainty on normalisation is at least $\sim 5\%$.

Testing normalisations

- Normalising using 'counts' is very reliable



HIMFs from next generation HI surveys

- (Probably) use the C⁻ method
 - Robust against LSS
 - Works with ‘soft’ completeness limits
 - Fast
 - Can be used for HIMF evolution, and environment
 - But it has problem at high HI masses...
- Normalisation: use counts
